

# Experimental determination of the rate constants of the reactions of HO2 + DO and DO2+ DO2

Mohamed Assali, Jozef Rakovsky, Ondrej Votava, Christa Fittschen

#### ▶ To cite this version:

Mohamed Assali, Jozef Rakovsky, Ondrej Votava, Christa Fittschen. Experimental determination of the rate constants of the reactions of HO2 + DO and DO2+ DO2. International Journal of Chemical Kinetics, 2020, International Journal of Chemical Kinetics, 52 (3), pp.197-206. hal-02992711

# HAL Id: hal-02992711 https://hal.univ-lille.fr/hal-02992711

Submitted on 6 Nov 2020

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



# Experimental Determination of the Rate Constants of the Reactions of HO2 + DO2 and DO2 + DO2

Journal:	International Journal of Chemical Kinetics
Manuscript ID	KIN-19-0148.R1
Wiley - Manuscript type:	Article
Date Submitted by the Author:	n/a
Complete List of Authors:	Assali, Mohamed; University of Lille Rakovsky, Jozef; J Heyrovsky Institute of Physical Chemistry Czech Academy of Sciences Votava, Ondrej; J Heyrovsky Institute of Physical Chemistry Czech Academy of Sciences Fittschen, Christa; University of Lille
Keywords:	kinetics, HO2 radicals, DO2 radicals

SCHOLARONE™ Manuscripts

# Experimental Determination of the Rate Constants of the Reactions of $HO_2 + DO_2$ and $DO_2 + DO_2$

Mohamed Assali<sup>1</sup>, Jozef Rakovsky<sup>2</sup>, Ondrej Votava<sup>2</sup>, Christa Fittschen<sup>1,#</sup>

- <sup>1</sup> Université Lille, CNRS, UMR 8522 PC2A Physicochimie des Processus de Combustion et de l'Atmosphère, F-59000 Lille, France
- <sup>2</sup> J. Heyrovský Institute of Physical Chemistry v.v.i., Academy of Sciences of the Czech Republic, Dolejškova 3, 18223 Prague, Czech Republic

#Corresponding author: Christa Fittschen (<u>christa.fittschen@univ-lille.fr</u>)
Tel: ++ 33 3 20 33 72 66

Submitted to

International Journal of Chemical Kinetics

Revised version

# **Abstract**

The rate constants of the reactions of  $DO_2 + HO_2$  (R1) and  $DO_2 + DO_2$  (R2) have been determined by the simultaneous, selective and quantitative measurement of HO<sub>2</sub> and DO<sub>2</sub> by cw-CRDS in the near infrared, coupled to a radical generation by laser photolysis. HO<sub>2</sub> was generated by photolysing Cl<sub>2</sub> in the presence of CH<sub>3</sub>OH and O<sub>2</sub>. Low concentrations of DO<sub>2</sub> were generated simultaneously by adding low concentrations of D<sub>2</sub>O to the reaction mixture, leading through isotopic exchange on tubing and reactor walls to formation of low concentrations of CH<sub>3</sub>OD and thus formation of DO<sub>2</sub>. Excess DO<sub>2</sub> was generated by photolysing Cl<sub>2</sub> in the presence of CD<sub>3</sub>OD and O<sub>2</sub>, small concentrations of HO<sub>2</sub> were always generated simultaneously by isotopic exchange between CD<sub>3</sub>OD and residual H<sub>2</sub>O. The rate constant  $k_1$  at 295 K was found to be pressure independent in the range 25 – 200 Torr helium, but increased with increasing D<sub>2</sub>O concentration  $k_1 = (1.67 \pm 0.03) \times 10^{-12} \times (1 + (8.2 \pm 1.6) \times 10^{-18})$ cm<sup>3</sup>× [D<sub>2</sub>O] cm<sup>-3</sup>) cm<sup>3</sup>s<sup>-1</sup>. The rate constant for the DO<sub>2</sub> self reaction  $k_2$  has been measured under excess DO<sub>2</sub> concentration, and the DO<sub>2</sub> concentration has been determined by fitting the  $HO_2$  decays, now governed by their reaction with  $DO_2$ , to the rate constant  $k_1$ . A rate constant with insignificant pressure dependence was found:  $k_2 = (4.1\pm0.6)\times10^{-13}$  (1 +  $(2\pm2)\times10^{-20}$  cm<sup>3</sup>× [He] cm<sup>-3</sup>) cm<sup>3</sup>s<sup>-1</sup> as well as an increase of  $k_2$  with increasing D<sub>2</sub>O concentration was observed:  $k_2 = (4.14\pm0.02)\times10^{-13}\times(1+(6.5\pm1.3)\times10^{-18}\,\mathrm{cm}^3\times\,\mathrm{[D_2O]}\,\mathrm{cm}^{-3})$ cm $^3$ s $^{-1}$ . The result for  $k_2$  is in excellent agreement with literature values, while this is the first determination of  $k_1$ .

### Introduction

The hydroperoxy radical, HO<sub>2</sub>, is a major radical in oxidation chemistry. In the atmosphere its concentration is closely linked to the OH radical by recycling it through reaction with NO. The self reaction of HO<sub>2</sub> radicals (R3) is a major sink for odd hydrogen in the atmosphere, and presents also the major source of H<sub>2</sub>O<sub>2</sub> in the stratosphere. The reaction has attracted a large interest in the last decades and its rate constant presents interesting features such as negative temperature dependence and pressure dependence. Small amounts of water vapor or methanol enhance the rate constant, while addition of NH<sub>3</sub> initially enhances the rate constant but then slows down the reaction at higher concentrations [1,[2,[3,[4,[5,[6,[7,[8,[9,[10,[11]. An extended review on the current knowledge and on the interpretation of this behavior is given by Stone and Rowley [7].

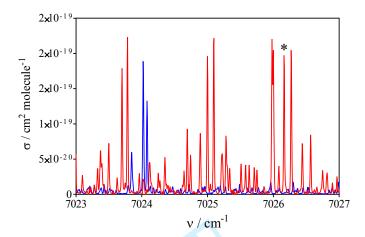
The self reaction of  $DO_2$ ,  $k_2$ , shows a similar behavior, but with a rate constant around four times slower than  $k_3$  and the pressure dependence roughly a factor of 2 weaker. The less pronounced pressure dependence for  $DO_2$  compared to  $HO_2$  is attributed to the faster approach to the high pressure limit for  $DO_2$  while the strong kinetic isotope effect has been explained in an early work by Mozurkewich and Benson [4] by using a cyclic transition state instead of the tetraoxide transition state favored by Patrick, Barker and Golden [12]. The six-membered ring complex as transition state has also been confirmed by more recent high level calculations [13,[14].

The rate constant of the cross reaction between the two isotopes HO<sub>2</sub> and DO<sub>2</sub>

$$DO_2 + HO_2 (+M) \rightarrow HDO_2 + O_2 (+M)$$
 (R1)

has never been measured to our knowledge. While this reaction is of no importance in the atmosphere, it might play a role in laboratory experiments, when deuterated compounds are used to elucidate reaction mechanisms [15]. The difficulty in measuring this rate constant is that both isomeric radicals need to be detected in a selective way. In most of the earlier works, HO<sub>2</sub> or DO<sub>2</sub> radicals have been detected by UV-absorption spectroscopy which does not allow a separation between both isomers, because their absorption spectra in the UV range are broad and unstructured. Thrush and Tyndall have in 1982 detected the HO<sub>2</sub> radical selectively in the  $v_3$  band at 1117 cm<sup>-1</sup> and have measured the rate constant of the self reaction [16]. Six years later, Martin and Thrush have in 1988 selectively detected the DO<sub>2</sub> radical in the v<sub>2</sub> band at 1020 cm<sup>-1</sup> and measured the rate constant of the DO<sub>2</sub> self reaction [17]. However, to our knowledge, no paper has reported so far experiments on the simultaneous, time resolved and selective detection of both isotopes. In the current work, we have used the selective detection of both radicals in the near IR range, where two highly structured absorption ranges can be found: the 2v<sub>1</sub> transition for HO<sub>2</sub> lies at around 6600 cm<sup>-1</sup> and has been extensively studied and used for its detection [10,[18,[19,[20,[21,[22,[23,[24,[25,[26]]. This transition is also used in this work for the quantification of HO2. The corresponding DO2 transition is expected to be around 5000 cm<sup>-1</sup> [27], but has not been studied to our knowledge, probably due to increased experimental difficulties occurring in this wavelength range compared to 6600 cm<sup>-1</sup>. Instead, the low-lying electronic transition of the  $\tilde{A}^{2}A' \leftarrow \tilde{X}^{2}A''$  000-000 band at around 7000 cm<sup>-1</sup> has been studied and used for the detection of DO<sub>2</sub> [15,[28,[29]]. Being an electronic transition, the HO<sub>2</sub> radical can be detected in the same wavelength range [29,[30,[31]. Figure 1 presents a portion of the  $\tilde{A}^2A' \leftarrow \tilde{X}^2A''$  000-000 spectrum for HO<sub>2</sub> (blue) and DO<sub>2</sub> (red). The transition used in this work for quantifying the DO<sub>2</sub> concentration

has been chosen to be free of interference from HO<sub>2</sub> transitions and is marked with an asterisk.



**Figure 1**: Portion of the HO<sub>2</sub> (blue) and DO<sub>2</sub> (red) spectrum at 25 Torr, adapted from [29]. The line marked with an asterisk indicates the line used in this work for DO<sub>2</sub> quantification.

In this work we present the experimental determination of  $k_1$ , by simultaneous selective detection of both isotopes. The rate constant of the self reaction of DO<sub>2</sub> radicals:

$$DO_2 + DO_2 (+M) \rightarrow D_2O_2 + O_2 (+M)$$
 (R2)

has also been measured. Both rate constant determinations are based on the determination of the absolute HO<sub>2</sub> concentration by fitting HO<sub>2</sub> decays to the well-known rate constant of the HO<sub>2</sub> self reaction, (R3):

$$HO_2 + HO_2 (+M) \rightarrow H_2O_2 + O_2 (+M)$$
 (R3)

In other words, the rate constants  $k_1$  and  $k_2$  have been measured relative to  $k_3$ .

# **Experimental Section**

Experiments have been carried out using pulsed laser photolysis coupled to cw-CRDS. The set-up has been described in detail elsewhere [29,[32,[33,[34]]] and only a brief description will be given here. The experimental set-up contains two cw-CRDS paths, which cross the photolysis beam symmetrically in a small angle. This allows the simultaneous, time resolved and absolute detection of two species. An overlap between photolysis and detection beam of 37.7cm is obtained. To control if the photolysis laser is well aligned, *i.e.* both cw-CRDS paths sample the same path lengths, HO<sub>2</sub> radicals are detected in an initial experiment on both paths: the retrieved concentrations from both paths agreed to better than 5%.

Experiments have been carried out under either excess  $HO_2$  over  $DO_2$ , or excess  $DO_2$  over  $HO_2$ .  $HO_2$  was detected on one path for all experiments within the  $2v_1$  vibrational overtone

band at either the strongest transition at 6638.205 cm<sup>-1</sup> or, when HO<sub>2</sub> concentrations were high and absorption became too strong, on a smaller transition at 6638.58 cm<sup>-1</sup>. DO<sub>2</sub> has been detected on the other path in the  $\tilde{A}$   $^2A$ '  $\leftarrow \tilde{X}$   $^2A$ '' 000-000 band at 7026.16 cm<sup>-1</sup>.

The reaction is initiated by 351 nm excimer laser photolysis (Lambda Physik, LPX 202i) of Cl<sub>2</sub>

$$Cl_2 + hv_{351nm} \rightarrow 2 Cl$$
 (R4)

The laser fluence was typically 30-40 mJ cm<sup>-2</sup>, leading to a photolysis yield of around 1% for Cl<sub>2</sub>. HO<sub>2</sub> radicals were generated by the reaction of the Cl-atoms with CH<sub>3</sub>OH in the presence of O<sub>2</sub>:

$$Cl + CH_3OH \rightarrow HCl + CH_2OH$$
 (R5)

$$CH2OH + O2 \rightarrow CH2O + HO2$$
 (R6)

For experiments with excess HO<sub>2</sub> over DO<sub>2</sub>, low concentrations of DO<sub>2</sub> radicals in the presence of excess HO<sub>2</sub> were generated by adding low concentrations of D<sub>2</sub>O to the gas flow. This way, rapid H/D isotope exchange between the labile –OH in CH<sub>3</sub>OH and D<sub>2</sub>O lead to formation of some CH<sub>3</sub>OD which then reacts with Cl [29]:

$$Cl + CH_3OD \rightarrow HCl + CH_2OD$$
 (R7)

with subsequent formation of DO<sub>2</sub> through

$$CH_2OD + O_2 \rightarrow CH_2O + DO_2 \tag{R8}$$

Excess  $DO_2$  over  $HO_2$  was generated through the reaction of Cl atoms with  $CD_3OD$  in the presence of  $O_2$ :

$$Cl + CD_3OD \rightarrow DCl + CD_2OD$$
 (R9)

$$CD_2OD + O_2 \rightarrow CD_2O + DO_2 \tag{R10}$$

H<sub>2</sub>O, always present in the system, led through H/D exchange always to the formation of some CD<sub>3</sub>OH and subsequently to formation of HO<sub>2</sub> through:

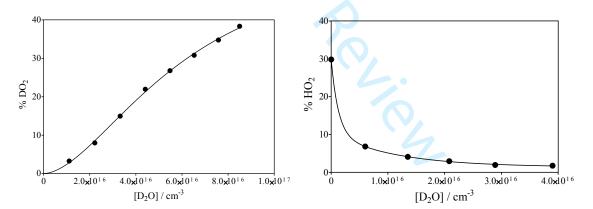
$$Cl + CD_3OH \rightarrow DCl + CD_2OH$$
 (R11)

$$CD_2OH + O_2 \rightarrow CD_2O + HO_2 \tag{R12}$$

In order to decrease the fraction of  $CD_3OH$  and thus  $HO_2$ , the gas flow was partially saturated with  $D_2O$  by flowing a varying fraction of the bath gas Helium through a trap containing liquid  $D_2O$ . Concentrations of  $CH_3OH$ ,  $CD_3OD$  and  $D_2O$  were estimated from partial pressures and flow rates, no information on the  $CD_3OH$  concentration was available, but can be estimated from the  $HO_2$  concentration formed in (R11) and (R12).

(R5) and (R6) are fast ( $k_5 = 5.5 \times 10^{-11} \text{ cm}^3 \text{s}^{-1}$  and  $k_6 = 9.6 \times 10^{-12} \text{ cm}^3 \text{s}^{-1}$  [35]), leading with typical concentrations of [CH<sub>3</sub>OH]  $\approx 8 \times 10^{14} \text{ cm}^{-3}$  and [O<sub>2</sub>] =  $1.1 \times 10^{17} \text{ cm}^{-3}$  to pseudo-first

order rates of  $k_5$ ' = 4.4×10<sup>4</sup> s<sup>-1</sup>, and  $k_6$ ' = 1.1×10<sup>6</sup> s<sup>-1</sup>, and are thus completed within few 10 μs. No rate constants for the corresponding reactions with the deuterated species are available in the literature, and even though it can be expected that these reactions will be slower due to the kinetic isotope effect, DO<sub>2</sub> will be equally formed on the µs time scale and will be in any case much faster than the peroxy self- and cross-reactions, which take place on the ms time scale. The ratio of [HO<sub>2</sub>] / [DO<sub>2</sub>] is given by the ratio of  $k_5 \times [CH_3OH] / k_7 \times [CH_3OD]$  (or  $k_9 \times [\text{CD}_3\text{OD}] / k_{11} \times [\text{CD}_3\text{OH}]$  respectively), whereby the last 2 terms are not known. However, a nearly linear increase of [DO<sub>2</sub>] with increasing [D<sub>2</sub>O] to the mixture Cl/CH<sub>3</sub>OH/O<sub>2</sub> was observed, while the fraction of HO<sub>2</sub> radicals in a mixture Cl/CD<sub>3</sub>OD/O<sub>2</sub> decreased very rapidly upon addition of a small flow of D<sub>2</sub>O. Examples for both experimental systems are shown in Figure 2: the left graph shows the percentage of Cl-atoms being converted to DO<sub>2</sub> in a mixture Cl/CH<sub>3</sub>OH/D<sub>2</sub>O/O<sub>2</sub> as a function of added [D<sub>2</sub>O], whereby the DO<sub>2</sub> concentration is obtained as the difference in HO<sub>2</sub> concentration before and after addition of D<sub>2</sub>O. The right graph shows the percentage of Cl-atoms being converted to HO<sub>2</sub> in a mixture Cl/CD<sub>3</sub>OD/D<sub>2</sub>O/O<sub>2</sub> as a function of [D<sub>2</sub>O]. Here, HO<sub>2</sub> concentrations have been measured directly using the absorption cross sections, DO<sub>2</sub> concentrations have been determined as explained further down.



**Figure 2**: Left graph: Fraction of Cl-atoms converted to  $DO_2$  as a function of added  $D_2O$  to a mixture  $Cl/CH_3OH/D_2O/O_2$ . Total pressure was 50 Torr,  $[Cl]_0 = 1.1 \times 10^{14}$  cm<sup>-3</sup>. Right graph: Remaining  $HO_2$  after adding different concentrations of  $D_2O$  to a mixture  $Cl/CD_3OD/D_2O/O_2$  in order to decrease concomitantly generated  $HO_2$ . Total pressure was 50 Torr,  $[Cl]_0 = 1.0 \times 10^{14}$  cm<sup>-3</sup>.

Experiments have been carried out in the pressure range 25 - 200 Torr, and initial [CI] concentration varied typically between  $5 - 15 \times 10^{13}$  cm<sup>-3</sup>. The concentration of the excess isotope (HO<sub>2</sub> or DO<sub>2</sub>) was typically in 5-10 times higher than the concentration of the other isotope (DO<sub>2</sub> or HO<sub>2</sub>).

The gas flow into the photolysis reactor was controlled using calibrated flowmeters (Tylan FC-260). The main flows consisted of Helium and O<sub>2</sub> and were directly taken from the cylinder (both Alphagaz 2). The precursor Cl<sub>2</sub> was also directly taken from a commercial cylinder (5% Cl<sub>2</sub> in Helium, Alpha Gaz), CH<sub>3</sub>OH (CD<sub>3</sub>OD) was added to the gas mixture by flowing helium through a bottle containing liquid CH<sub>3</sub>OH (CD<sub>3</sub>OD) at room temperature, the concentration was estimated from the CH<sub>3</sub>OH vapor pressure, the total pressure in the bottle and the measured flow rate. D<sub>2</sub>O was either prepared as a mixture (2%) in a 20 l glass balloon and small concentrations were added through flowmeter in case of HO<sub>2</sub>-excess experiments or high concentrations were added by flowing a fraction of the main Helium through a bottle containing liquid D<sub>2</sub>O for DO<sub>2</sub>-excess experiments.

# Results and Discussion

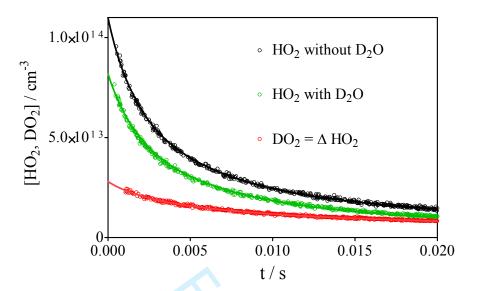
#### Measurement of the rate constant of $HO_2 + DO_2$

Experiments have been carried out under large excess of HO<sub>2</sub> over DO<sub>2</sub>. Under these conditions, the decay of the excess HO<sub>2</sub> is nearly exclusively governed by the rate constant for its self-reaction  $k_3$ , while for the DO<sub>2</sub> radical the self-reaction plays only a minor role due to lower concentrations and to a lower rate constant of the self-reaction  $k_2$ . Therefore, the DO<sub>2</sub> decay is governed nearly exclusively by their reaction with excess HO<sub>2</sub>; i.e. by the rate constant  $k_1$  and the absolute concentration of the HO<sub>2</sub> radical. These experimental conditions therefore result in a reliable determination of the rate constant of the cross reaction  $k_l$ , if the absolute HO<sub>2</sub> concentration can be determined. This is possible not only by converting the absorption-time profiles into concentration-time profiles using the well-known absorption cross sections [10,[21,[22,[23,[25,[26]], but also by fixing  $k_3$  to the recommended literature value and subsequently adjusting the initial HO<sub>2</sub> concentration to best fit the HO<sub>2</sub> decay: both methods have been applied independently and led always to initial HO<sub>2</sub> concentration agreeing to better than +/-10% (see below). In other words: by choosing excess HO<sub>2</sub> concentration conditions, the rate constant  $k_1$  is measured relative to the well-known rate constant of the  $HO_2$  self-reaction  $k_3$ . Any uncertainties in the absorption cross sections of  $HO_2$ and DO<sub>2</sub> (the absorption lines are very narrow, especially at low pressure, and a small drift in the wavelength emitted by the DFB laser will result in a change of the effective absorption cross section) will not impact the retrieved rate constants  $k_1$ : the HO<sub>2</sub> absorption cross section is refined for each experiment by fitting the  $HO_2$  decays to the fixed  $k_3$ , while the absolute

value of the absorption cross section of  $DO_2$  has only a very minor impact for retrieving  $k_I$ . These conditions minimize the uncertainty on  $k_I$ , because even if cross sections and pressure broadening factors of several strong absorption lines in the  $2v_1$  transition of  $HO_2$  have been well studied [10,[21,[22,[23,[25,[26], the rather small line at 6638.58 cm<sup>-1</sup>, that was used in this work in order to avoid saturation of the absorption due to the high initial radical concentrations, is less well known.

Figure 3 shows a typical experiment at 50 Torr helium: the black trace shows the HO<sub>2</sub> profile of an experiment in the absence of D<sub>2</sub>O, the green trace shows the HO<sub>2</sub> profile measured under the same conditions, but after addition of a small flow of D<sub>2</sub>O to the gas flow; the red trace shows the corresponding DO<sub>2</sub> profile measured simultaneously on the second cw-CRDS path. In a first step, the HO<sub>2</sub> absorption profile in absence of D<sub>2</sub>O is fitted to the literature value for the rate constant of the self-reaction,  $k_3$ , using a custom-designed Labview based program. This way the absorption cross section of HO<sub>2</sub> can be retrieved at the given pressure. This step (measuring HO<sub>2</sub> profiles in absence of D<sub>2</sub>O) is carried out at each pressure only for some  $Cl_2$  concentrations. In this work, we have used the rate constant for  $k_3$  in helium from Sander et al [1]. They have measured the pressure dependence of  $k_3$  for different bath gases (He, Ar, N<sub>2</sub>, O<sub>2</sub> and SF<sub>6</sub>) and their value for air (ie.  $0.2 \times k_{3,O2} + 0.8 \times k_{3,N2}$ ) is in excellent agreement with recommended values from IUPAC [36]. The reliability of  $k_3$  is estimated by the IUPAC committee to be  $\pm 40\%$ . Because all results in this work are determined relative to  $k_3$ , this uncertainty of  $\pm 40\%$  needs to be considered also for our  $k_1$  and  $k_2$  measurements. However, in what follows, we will only consider uncertainties occurring from our measurements and not add the uncertainty of  $k_3$ .

Once the  $HO_2$  concentration (and with this the absorption cross section) in absence of  $DO_2$  has been determined, the  $HO_2$  absorption time profile in presence of  $D_2O$  at the same  $Cl_2$  concentration is converted to a concentration-time profile using that absorption cross section. Now, the initial  $DO_2$  concentration can be calculated as the difference between both  $HO_2$  concentrations and hence the  $DO_2$  absorption cross section at the given pressure can be determined from the absorption at t=0 s. Finally, the rate constant  $k_I$  is adjusted to best fit the  $DO_2$  profile.



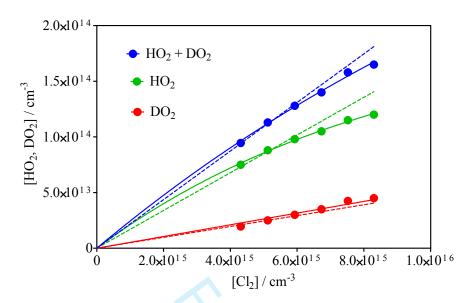
**Figure 3**: HO<sub>2</sub> concentration time profiles at 50 Torr in absence (black) and presence (green) of D<sub>2</sub>O and corresponding DO<sub>2</sub> profile (red) measured simultaneously (*i.e.* the red trace is not the difference between black and green). Initial HO<sub>2</sub> concentrations have been retrieved by fitting decay kinetics using  $k_3 = 1.58 \times 10^{-12}$  cm<sup>3</sup>s<sup>-1</sup> [1], initial DO<sub>2</sub> concentration has been normalized as the difference in HO<sub>2</sub> concentration, the DO<sub>2</sub> decay has been fitted by adjusting  $k_1$ .

In practice, for retrieving the rate constant  $k_1$ , the HO<sub>2</sub> and DO<sub>2</sub> absorption profiles for a series of experiments at different Cl<sub>2</sub> concentrations (typically 6 at each pressure, see below) have initially all been converted to concentration-time profiles using the same absorption cross section obtained such as described above. However, HO<sub>2</sub> absorption lines are narrow, and a very minor deviation of the near IR-laser wavelength between different experiments can lead to a variation in the absorption cross section, especially for the low pressure experiments. Therefore, all 6 HO<sub>2</sub> and DO<sub>2</sub> concentration-time profiles have been fitted simultaneously using a very small mechanism consisting of (R1), (R2) and (R3) plus small losses due to diffusion of the radicals out of the photolysis volume. Unique rate constants for all reactions have been used for all different Cl-concentrations at each pressure. The individual HO<sub>2</sub> decays are very sensitive to the initial concentration (see below) and have been finely adjusted to best fit the unique rate constant  $k_3$ . The correction for the individual absorption cross sections was less than 10% for all experiments. The values for the absorption cross sections for HO<sub>2</sub>, averaged for each pressure over all experiments, and DO<sub>2</sub> obtained such as described above, are summarized in Table 1. Assuming a Voigt profile of the pressure broadened absorption lines, the integrated line strength and pressure broadening parameter have been determined from the data in Table 1 with values leading for DO<sub>2</sub> (HO<sub>2</sub>) to a line strength of  $(6.8\pm0.7)\times10^{-21}$  ((7.9±0.8)×10<sup>-22</sup>) cm and a broadening coefficient of 0.165 (0.06) cm<sup>-1</sup>/atm.

**Table 1**: Absorption cross sections of  $HO_2$  and  $DO_2$  for all pressures. Error bars contain only the uncertainty from our measurements, an additional uncertainty of  $\pm 40\%$  needs to be considered due to the uncertainty of  $k_3$  [36]

P / Torr	$\sigma_{HO2}$ / $10^{-20} cm^2$ at 6638.58 cm <sup>-1</sup>	$\sigma_{DO2}$ / $10^{\text{-}20}$ cm² at 7026.16 cm <sup>-1</sup>	
25	$3.3 \pm 0.3$	20 ± 2	
50	$2.8 \pm 0.3$	13 ± 1	
75	$2.4 \pm 0.2$	12 ± 1	
100	$2.1 \pm 0.2$	$9.3 \pm 0.9$	
150	$1.7 \pm 0.2$	$6.3 \pm 0.6$	
200	$1.4 \pm 0.1$	$4.5 \pm 0.4$	

This type of experiment has been carried out at each pressure for a series of 6 different initial  $Cl_2$  concentrations. The  $HO_2$  and  $DO_2$  concentrations as well as the sum of both radicals should increase linearly with increasing  $Cl_2$  concentration. **Table 2** shows these data for a series of experiments at 50 Torr, and the values are plotted in **Figure 4**. The increase in radical concentration with  $Cl_2$  is not perfectly linear (dashed line, forced through origin), but a slight saturation can be observed at the highest  $Cl_2$  concentrations. This might be due to a fast reaction such as  $Cl + CH_2OH$  ( $k = 6.6 \times 10^{-10}$  cm<sup>3</sup>s<sup>-1</sup>) [37], that starts to compete to (R6) or (R8) with increasing Cl concentration. Only one determination is available for the rate constant of the reaction  $Cl + CH_2OH$ , and such experiments could possibly be designed in a way (varying  $O_2$  concentration at high Cl-concentration) to re-determine this rate constant, but is out of scope of this paper.

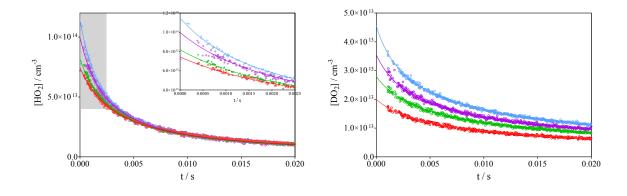


**Figure 4**: Data from **Table 2**, showing the initial HO<sub>2</sub> (in presence of D<sub>2</sub>O), DO<sub>2</sub> and sum of both as a function of Cl<sub>2</sub> concentration

**Table 2**: Evolution of initial HO<sub>2</sub> (in presence of D<sub>2</sub>O) and DO<sub>2</sub> concentration with increasing Cl<sub>2</sub> concentration, the example at 50 Torr

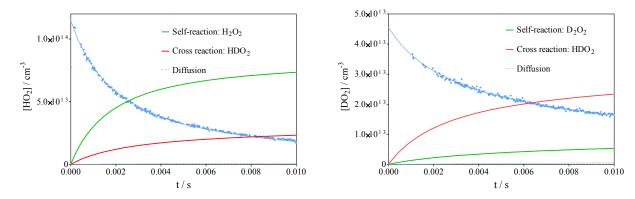
er concentration, the example at 50 Toll								
$\mathrm{Cl_2}/\ 10^{15}\ \mathrm{cm^{-3}}$	$HO_2/10^{14} \text{ cm}^{-3}$	$DO_2/10^{14} \text{ cm}^{-3}$	$HO_2 + DO_2 / 10^{14} \text{ cm}^{-3}$					
8.30	1.20	0.45	1.65					
7.52	1.15	0.43	1.58					
6.73	1.05	0.35	1.40					
5.93	1.0	0.30	1.28					
5.12	0.9	0.25	1.15					
4.31	0.8	0.20	1.0					

Figure 5 shows a series of HO<sub>2</sub> and DO<sub>2</sub> profiles with different Cl<sub>2</sub> concentrations. For better visibility, the profiles of two intermediate Cl<sub>2</sub> concentrations have been omitted. The full lines show the fit to the simple model containing (R1), (R2) and (R3) as well a loss through diffusion for each radical (very minor impact), using the rate constants from **Table 3**, the inset in the left graph shows a zoom on the HO<sub>2</sub> profiles at short reaction times.



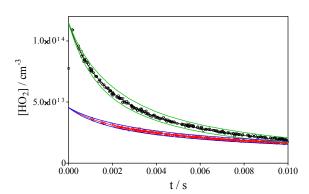
**Figure 5**: Typical series of experiments with increasing Cl<sub>2</sub> concentration. Left graph: HO<sub>2</sub> profiles, right graph DO<sub>2</sub> profiles. The inset in the left graph shows a zoom of the HO<sub>2</sub> profiles on the first 2.5 ms. No ring-down events occur during the first 1 ms on the DO<sub>2</sub> absorption path, possibly due to perturbation of the DFB laser from scattered 351 nm photons. This phenomena is not observed when using 248 nm photolysis (see [29])

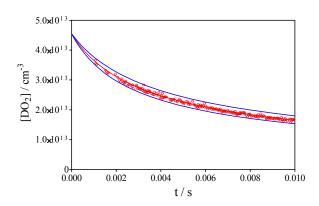
**Figure 6** shows the distribution of products for both reaction partners ( $HO_2$  on the left graph,  $DO_2$  on the right graph, green for self-reaction, red for cross reaction) for the experiment with the highest initial radical concentration from **Figure 5**: it can be seen that for  $HO_2$  the self-reaction is the major loss pathway, while for  $DO_2$  self-reaction is very minor and the decay is nearly exclusively governed by its reaction with  $HO_2$ .



**Figure 6**: Signal with highest Cl<sub>2</sub> concentration from **Figure 5** with the loss through self-reaction (green line) and cross reaction (red line) for HO<sub>2</sub> (left) and DO<sub>2</sub> (right).

Figure 7 shows the sensitivity of the signals to the rate constant  $k_I$ : the left graph shows both, HO<sub>2</sub> (black) and DO<sub>2</sub> (red), while DO<sub>2</sub> is shown again on the right graph on a zoomed y-scale. The full lines show the best fit using the rate constants from **Table 3** and are barely visible, while for the upper and lower green lines the initial HO<sub>2</sub> concentration in the model has been varied by +/- 20% (*i.e.* simulating an uncertainty in the HO<sub>2</sub> absorption cross section of +/- 20%). For better visibility, the simulated green curves have been multiplied by 1.2 / 0.8 to match the experimental data, rather than plotting three different experimental curves in which the absorption coefficient α=f(t) would have been converted to [HO<sub>2</sub>]=f(t) using three different absorption cross sections σ. It can be seen that the HO<sub>2</sub> concentration-time profile is not well reproduced anymore with a 20% change in initial concentration, ie. at a fixed rate constant for the HO<sub>2</sub> self reaction, the HO<sub>2</sub> concentration can be determined to better than 20%. At the same time, the 20% change in initial HO<sub>2</sub> concentration leads to a change of the DO<sub>2</sub> concentration-time profile well outside the experimental uncertainty (blue lines).





**Figure 7**: Simulation showing the sensitivity to the initial  $HO_2$  concentration for an experiment at 50 Torr. Black dots: experimental  $HO_2$  concentration time profile, red dots: experimental  $DO_2$  concentration time profile (zoomed in the right graph). Full lines (barely visible): fit to the model of **Table 3**, with the  $HO_2$  concentration adjusted to best fit the decay, lower and upper green lines show variation of  $HO_2$  profile with a +/- 20% change in initial concentration. For visual demonstration of the effect, the simulated green curves have been adapted to match the initial experimental  $HO_2$  concentration, *ie.* have been multiplied by 1.2 / 0.8.

Final results for the rate constants  $k_1$  (and  $k_2$ , see below) for all different pressures are presented in **Figure 9** and **Table 3**. The rate constant for the cross reaction of HO<sub>2</sub> and DO<sub>2</sub> radicals can be given as pressure independent in the range 25 - 200 Torr helium as

$$k_I = (1.6 \pm 0.3) \times 10^{-12} \text{ cm}^3 \text{s}^{-1}$$

The error bar includes an uncertainty of 10% from the fitting of the decays and an uncertainty of 10% for the rate constant of the HO<sub>2</sub> self reaction.

**Table 3**: Rate constants used to fit all experiments. Values  $k_2$  in **bold** are measured in absence of added  $D_2O$ , values in *italic* are interpolated. Values for  $k_3$  are calculated from Sander et al. [1] with M = He

P / Torr	$k_1$ $HO_2 + DO_2$ $/ 10^{-12} \text{ cm}^3 \text{s}^{-1}$	$k_2$ $DO_2 + DO_2$ $/ 10^{-13} \text{ cm}^3 \text{s}^{-1}$	$k_3$ $HO_2 + HO_2$ $/ 10^{-12} \text{ cm}^3 \text{s}^{-1}$	Diffusion / s <sup>-1</sup>
25	1.7±0.3	4.0±0.8	1.56	5
50	1.6±0.3	<b>4.1</b> ±0.8	1.58	3
75	1.6±0.3	4.2±0.8	1.60	2
100	1.6±0.3	<b>4.3</b> ±0.8	1.61	1
150	1.6±0.3	4.4±0.8	1.65	1
200	1.6±0.3	<b>4.5</b> ±0.8	1.69	0

#### Measurement of the rate constant $k_2$ for the reaction $DO_2 + DO_2$

Determining the rate constant for the  $DO_2$  self-reaction,  $k_3$ , requires knowledge of the absolute  $DO_2$  concentration. However, only very little is known about absolute absorption cross sections and pressure broadening in the near infrared [29]. Therefore, two different methods can be used for determining the  $DO_2$  concentration. (a) the Cl-atom concentration is initially determined by adding  $CH_3OH$  and thus transforming them quantitatively into  $HO_2$  radicals. After switching to  $CD_3OD/D_2O$  mixtures, the  $DO_2$  and  $HO_2$  absorption time profiles are determined, and the initial  $DO_2$  concentration (and thus the absorption cross section) can be determined under the hypothesis that the Cl-atom concentration has not changed. This is equivalent to the above described method. (b) Fitting the  $HO_2$  decays in presence of excess  $DO_2$  to the reaction system (R1) – (R3) and fixing the rate constant  $k_1$  to the above obtained value. Now the  $HO_2$  decay will depend on the initial  $DO_2$  concentration. Both methods (a) and (b) have been applied independently and have returned values for the initial  $DO_2$  concentrations that agreed to better than 15% to each other. The so-obtained absorption cross section for  $DO_2$  was then used to convert absorption time profiles into absolute  $DO_2$  concentration time profiles, and to determine the rate constant for the  $DO_2$  self-reaction.

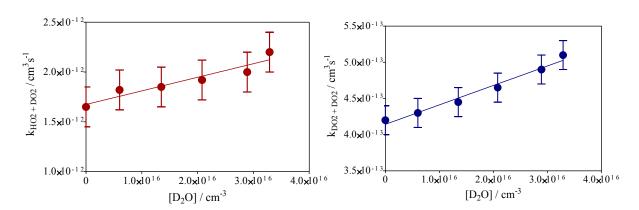
In order to minimize the  $HO_2$  concentration that is always formed through (R11) and (R12),  $D_2O$  has been added to the gas flow by flowing a fraction of the bath gas Helium through a bubbler containing  $D_2O$ . The fraction of Cl-atoms being converted to  $HO_2$  decreased rapidly from around 30% in absence of any added  $D_2O$  over 8% in presence of  $[D_2O] \approx 6 \times 10^{15}$  cm<sup>-3</sup> to less than 2% with  $[D_2O] \approx 3.5 \times 10^{16}$  cm<sup>-3</sup> (see **Figure 2**).

It was observed that both rate constants,  $k_1$  and  $k_2$ , increased with increasing D<sub>2</sub>O concentration, as shown in **Figure 8**. This is in agreement with observations for an increase in  $k_2$  with increasing H<sub>2</sub>O concentration.

$$k_1 = (1.67 \pm 0.03) \times 10^{-12} \times (1 + (8.2 \pm 1.6) \times 10^{-18} \text{ cm}^3 \times [D_2\text{O}] \text{ cm}^{-3}) \text{ cm}^3 \text{s}^{-1}$$
  
 $k_2 = (4.14 \pm 0.02) \times 10^{-13} \times (1 + (6.5 \pm 1.3) \times 10^{-18} \text{ cm}^3 \times [D_2\text{O}] \text{ cm}^{-3}) \text{ cm}^3 \text{s}^{-1}$ 

Error bars represent statistical error only for intercept and an addition error of 10% for uncertainty in the  $D_2O$  concentration. The effect observed in this work with  $D_2O$  is 2 to 3 times larger than the effect with  $H_2O$ , which has been found by Kircher and Sander [2] to be

$$k_3 = 1.6 \times 10^{-12} \times (1 + 2.25 \times 10^{-18} \text{ cm}^3 \times [\text{H}_2\text{O}] \text{ cm}^{-3}) \text{ cm}^3 \text{s}^{-1}$$

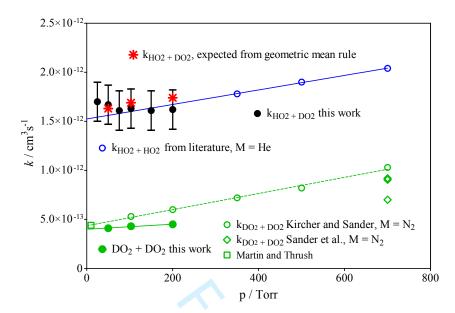


**Figure 8**: Rate constants  $k_1$  (left) and  $k_2$  (right) as a function of  $D_2O$  concentration.

Measurements have been carried out at 50, 100 and 200 Torr (values are given in **bold** in **Table 3**) and insignificant pressure dependence is observed:

$$k_2 = (4.1\pm0.6)\times10^{-13} (1 + (2\pm2)\times10^{-20} \text{ cm}^3\times \text{[He] cm}^{-3}) \text{ cm}^3\text{s}^{-1}$$

The error bars represent statistical error only, the systematic error due to the uncertainty in  $k_3$  is not taken into account. The data are plotted in **Figure 9**, together with the values for  $k_1$  as obtained in this work, and  $k_2$  from literature. No data for  $k_2$  in helium could be found in the literature, but extrapolation to p = 0 shows a very good agreement of the current results with literature data. Interestingly, in the work of Sander et al,  $k_2$  has been measured at 760 Torr using 2 different precursors: either using  $Cl_2$  photolysis in the presence of  $CD_3OD/O_2$  or in the presence of  $D_2/O_2$ . The rate constant they obtained using the first system, was significantly higher than when they used the second system, a difference which was considered by the authors beyond experimental uncertainty and could not be unexplained. A possible explanation would be that in their experiments residual  $H_2O$  was also present, which led to some  $HO_2$  next to the desired  $DO_2$ . However, their detection method (UV-absorption spectroscopy) did not allow distinguishing  $DO_2$  from  $HO_2$ , and thus the signal represented the sum of  $HO_2$  and  $DO_2$  and the decay was partially due to the faster reaction of  $DO_2$  with  $HO_2$ . When using  $D_2$  as precursor, no formation of  $HO_2$  is expected, and the observed decay is only due to the slower  $DO_2$  self-reaction.



**Figure 9**: Rate constants  $k_1$  (HO<sub>2</sub> + DO<sub>2</sub>, black dots) and  $k_2$  (DO<sub>2</sub> + DO<sub>2</sub>, green dots) from this work. Literature data for  $k_2$ : open circles from Kircher and Sander [2], open diamands from Sander *et al* [1] and open square from Martin and Thrush [17]. Red symbols show  $k_1$  such as expected from geometrical rule, taking  $k_2$  from this work and  $k_3$  from Kircher and Sander [2] (blue symbols).

The geometric mean rule is an empirical approach that allows for the estimation of cross-reaction rate coefficients from the self-recombination rate constants of the reacting partners [38]

$$k_{A+B} = 2 \times \sqrt{k_{A+A} \times k_{B+B}}$$

It has shown to work to better than 20% in the prediction of radical-radicals rate coefficients for a series of hydrocarbon radicals [39]. In absence of any literature data for  $k_1$ , we have tentatively applied this rule to estimate  $k_1$  by using the literature values for  $k_3$  and the values for  $k_2$  such as obtained in this work. In **Figure 9** are shown as red stars the values for  $k_1$  obtained this way: excellent agreement is found with our experimental determinations.

## Conclusion

The rate constant of the cross reaction between  $HO_2$  and  $DO_2$  radicals has been measured for the first time thanks to a simultaneous and selective measurement of absolute concentrations of  $HO_2$  and  $DO_2$  radicals by time resolved cw-CRDS in the near infrared. The rate constant was determined under conditions where the  $HO_2$  radical was in excess over the  $DO_2$  radical. Under these conditions, the decay of  $DO_2$  is essentially given by the absolute concentration of  $HO_2$  radicals and the rate constant  $k_1$ . The  $HO_2$  concentration in turn can be determined by

fitting its decay, which is essentially governed by its self-reaction. The rate constant  $k_I$  has been found independent on pressure in the range 25 – 200 Torr helium, but increasing with added  $D_2O$ :

$$k_1 = (1.67 \pm 0.03) \times 10^{-12} \times (1 + (8.2 \pm 1.6) \times 10^{-18} \text{ cm}^3 \times [D_2 \text{O}] \text{ cm}^{-3}) \text{ cm}^3 \text{s}^{-1}$$

The rate constant for the self reaction of  $DO_2$  has been determined by quantifying the  $DO_2$  concentration in two different ways: either Cl-atoms were initially converted to  $HO_2$  radicals (which can be quantified by cw-CRDS) by addition of  $CH_3OH$  which can be quantified. Then  $CH_3OH$  was replaced by  $CD_3OD$ , and it was considered that the Cl-atom concentration had not changed. In an independent method, the decay of the remaining small concentration of  $HO_2$  radicals next to the excess  $DO_2$  radicals was fitted using the above determined rate constant  $k_I$ . Both methods gave consistent concentrations and led to a rate constant with insignificant pressure dependence for the self-reaction of  $DO_2$ , in good agreement with the available literature data:

$$k_2 = (4.1\pm0.6)\times10^{-13} (1 + (2\pm2)\times10^{-20} \text{ cm}^3\times \text{[He] cm}^{-3}) \text{ cm}^3\text{s}^{-1}$$

For this reaction also, an increase of  $k_2$  with increasing concentration of  $D_2O$  was found:

$$k_2 = (4.14 \pm 0.02) \times 10^{-13} \times (1 + (6.5 \pm 1.3) \times 10^{-18} \text{ cm}^3 \times [D_2 \text{O}] \text{ cm}^{-3}) \text{ cm}^3 \text{s}^{-1}$$

# Acknowledgements

This project was supported by the French ANR agency under contract No. ANR-11-Labx-0005-01 CaPPA (Chemical and Physical Properties of the Atmosphere), the Région Hauts-de-France, the Ministère de l'Enseignement Supérieur et de la Recherche (CPER Climibio) and the European Fund for Regional Economic Development. O. Votava and J. Rakovsky thank for financial support through the PHC Barrande project no. 38203PM.

# References

- 1. Sander, S. P.; Peterson, M.; Watson, R. T.; Patrick, R. J Phys Chem 1982, 86, 1236-1240.
- 2. Kircher, C. C.; Sander, S. P. The Journal of Physical Chemistry 1984, 88, 2082-2091.
- 3. Christensen, L. E.; Okumura, M.; Sander, S. P.; Salawitch, R. J.; Toon, G. C.; Sen, B.; Blavier, J. F.; Jucks, K. W. Geophys. Res. Lett. 2002, 29, 1299.
- 4. Mozurkewich, M.; Benson, S. W. Int J Chem Kinet 1985, 17, 787-807.
- 5. Kurylo, M. J.; Ouellette, P. A.; Laufer, A. H. J Phys Chem 1986, 90, 437-440.
- Kanno, N.; Tonokura, K.; Tezaki, A.; Koshi, M. J Phys Chem A 2005, 109, 3153-3158.
- 7. Stone, D.; Rowley, D. M. PCCP 2005, 7, 2156 2163.
- 8. Hamilton, E. J.; Lii, R.-R. Int J Chem Kinet 1977, 9, 875-885.
- 9. Lii, R.-R.; Gorse, R. A.; Sauer, M. C.; Gordon, S. J Phys Chem 1980, 84, 819-821.
- 10. Tang, Y.; Tyndall, G. S.; Orlando, J. J. J Phys Chem A 2010, 114, 369-378.
- 11. Kanno, N.; Tonokura, K.; Koshi, M. J. Geophys. Res. 2006, 111, D20312.
- 12. Patrick, R.; Barker, J. R.; Golden, D. M. The Journal of Physical Chemistry 1984, 88, 128-136.
- 13. Zhu, R. S.; Lin, M. C. PhysChemComm 2001, 4, 106-111.
- 14. Donaldson, D. J.; Francisco, J. S. PCCP 2003, 5, 3183-3187.
- 15. Estupiñán, E. G.; Smith, J. D.; Tezaki, A.; Klippenstein, S. J.; Taatjes, C. A. J Phys Chem A 2007, 111, 4015-4030.
- 16. Thrush, B. A.; Tyndall, G. S. J Chem Soc, Faraday Trans 1982, 78, 1469-1475.
- 17. Martin, N. A.; Thrush, B. A. Chem Phys Lett 1988, 153, 200-202.
- 18. Taatjes, C. A.; Oh, D. B. Appl Opt 1997, 36, 5817-5821.
- 19. DeSain, J. D.; Ho, A. D.; Taatjes, C. A. J Mol Spectrosc 2003, 219, 163-169.
- 20. Noell, A. C.; Alconcel, L. S.; Robichaud, D. J.; Okumura, M.; Sander, S. P. J Phys Chem A 2010, 114, 6983-6995.
- 21. Assaf, E.; Liu, L.; Schoemaecker, C.; Fittschen, C. Journal of Quantitative Spectroscopy & Radiative Transfer 2018, 211, 107-114.
- 22. Thiebaud, J.; Crunaire, S.; Fittschen, C. J Phys Chem A 2007, 111, 6959-6966.
- 23. Ibrahim, N.; Thiebaud, J.; Orphal, J.; Fittschen, C. J Mol Spectrosc 2007, 242, 64-69.
- 24. Thiebaud, J.; Aluculesei, A.; Fittschen, C. J Chem Phys 2007, 126, 186101.
- 25. Thiebaud, J.; Fittschen, C. Appl Phys B 2006, 85, 383-389.
- 26. Onel, L.; Brennan, A.; Gianella, M.; Ronnie, G.; Lawry Aguila, A.; Hancock, G.; Whalley, L.; Seakins, P. W.; Ritchie, G. A. D.; Heard, D. E. Atmos. Meas. Tech. Discuss. 2017, 10, 4877-4894.
- 27. Rothman, L. S.; Gordon, I. E.; Babikov, Y.; Barbe, A.; Chris Benner, D.; Bernath, P. F.; Birk, M.; Bizzocchi, L.; Boudon, V.; Brown, L. R.; Campargue, A.; Chance, K.; Cohen, E. A.; Coudert, L. H.; Devi, V. M.; Drouin, B. J.; Fayt, A.; Flaud, J. M.; Gamache, R. R.; Harrison, J. J.; Hartmann, J. M.; Hill, C.; Hodges, J. T.; Jacquemart, D.; Jolly, A.; Lamouroux, J.; Le Roy, R. J.; Li, G.; Long, D. A.; Lyulin, O. M.; Mackie, C. J.; Massie, S. T.; Mikhailenko, S.; Müller, H. S. P.; Naumenko, O. V.; Nikitin, A. V.; Orphal, J.; Perevalov, V.; Perrin, A.; Polovtseva, E. R.; Richard, C.; Smith, M. A. H.; Starikova, E.; Sung, K.; Tashkun, S.; Tennyson, J.; Toon, G. C.; Tyuterev, V. G.; Wagner, G. J Quant Spectrosc Radiat Transfer 2013, 130, 4-50.
- 28. Fink, E. H.; Ramsay, D. A. J Mol Spectrosc 2002, 216, 322-334.
- 29. Assaf, E.; Asvany, O.; Votava, O.; Batut, S.; Schoemaecker, C.; Fittschen, C. J Quant Spectrosc Radiat Transfer 2017, 201, 161-170.
- 30. Fink, E. H.; Ramsay, D. A. J Mol Spectrosc 1997, 185, 304-324.
- 31. Clifford, E. P.; Farrell, J. T.; DeSain, J. D.; Taatjes, C. A. J Phys Chem A 2000, 104, 11549-11560.
- 32. Parker, A.; Jain, C.; Schoemaecker, C.; Szriftgiser, P.; Votava, O.; Fittschen, C. Appl Phys B 2011, 103, 725-733.
- 33. Votava, O.; Masat, M.; Parker, A. E.; Jain, C.; Fittschen, C. Rev Sci Instrum 2012, 83, 043110.
- 34. Assaf, E.; Fittschen, C. J Phys Chem A 2016, 120, 7051-7059.

- 35. Atkinson, R. B., D.L.; Cox, R.A.; Crowley, J.N.; Hampson, R.F, Jr.; Kerr, J.A.; Rossi, M.J.; Troe, J. IUPAC Subcommittee on Gas Kinetic Data Evaluation for Atmospheric ChemistryWeb Version December 2001 2001, 1 56.
- 36. Atkinson, R.; Baulch, D. L.; Cox, R. A.; Crowley, J. N.; Hampson, R. F.; Hynes, R. G.; Jenkin, M. E.; Rossi, M. J.; Troe, J. ACP 2004, 4, 1461-1738.
- 37. Pagsberg, P.; Munk, J.; Sillesen, A.; Anastasi, C. Chem Phys Lett 1988, 146, 375-381.
- 38. Jasper, A. W.; Klippenstein, S. J.; Harding, L. B. J Phys Chem A 2007, 111, 8699-8707.
- 39. Klippenstein, S. J.; Georgievskii, Y.; Harding, L. B. PCCP 2006, 8, 1133-1147.

