

# Wet foam flow: A suitable method for improving surface hygiene in the food industry

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1 Wet foam flow: a suitable method for improving surface hygiene in the food industry

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## 16 ABSTRACT

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In the food industry, the cleaning of contaminated surfaces requires new strategies to be adopted which can provide greater cleaning efficiency with minimal energy and water consumption. The use of wet foams was proposed to clean stainless steel surfaces contaminated by droplets containing Bacillus spores. Methods such as polarography, conductimetry and foam structure identification were used. Foam flow conditions with varying wall shear stresses and bubble sizes were observed to measure their impacts on surface contaminant removal kinetics. Compared to conventional cleaning-in-place method, foam flow more effectively removed both hydrophilic and hydrophobic spores. The combination of high shear stress and small bubble sizes (<0.2 mm) showed promise for improving the cleaning efficiency of the foam. Mechanisms such as fluctuation in local stresses, or in the liquid film thickness between the bubbles and the steel wall induced by bubble passage, foam imbibition and liquid film drainage were then investigated. Keywords Flow foam cleaning in place; Bacillus spores; conductimetry; polarography; bubble size; wall shear stress 

42

## 43 1. Introduction

44

In the agro-food industrial environment, avoiding surface contamination by unwanted 45 46 microorganisms is of crucial importance (Alvarez-Ordonez et al., 2019; Farag et al., 2021). Indeed, 47 contaminated surfaces can recontaminate food in contact (cross-contamination), with sometimes disastrous consequences in terms of food spoilage, but also in terms of transmission of pathogens 48 49 that can be the cause of large-scale epidemics. For example, surface contamination with 50 microorganisms would have accounted for more than 50% of the collective foodborne illnesses (CFTI) in France between 2006 and 2008 (Delmas et al., 2010). Nowadays, cleaning-in-place (CIP) is a 51 52 routine method in the food and pharmaceutical industries to ensure the equipment hygiene and the 53 quality of product. The effect of the hydrodynamics of the cleaning fluid on CIP efficiency has been 54 the subject of numerous studies (Li et al., 2019; Piepiórka-Stepuk et al., 2021). Particular interest has 55 been shown in parameters such as the mechanical action exerted by the flow in creating wall shear 56 stress (Bénézech and Faille, 2018) or the impact of Reynolds number (Fan et al., 2018). Their results 57 showed that turbulence affects the cleaning process, as an increase in the presence of turbulent 58 structures generated at the walls could improve the cleaning efficiency, but up to a certain limit. The 59 use of non-stationary flows, such as pulsed flows and intermittent jet flows, would appear to 60 increase CIP efficiency. Indeed, these kinds of flows affect both the mean and fluctuating 61 components of the wall shear stress and consequently reduce the residual contamination (Absi and 62 Azouani, 2018; Blel et al., 2013).

Gas-liquid two-phase flows (air bubbles in liquid or foam), which are also non-stationary flows, are likely to be of similar interest to pulsed flows. Indeed, liquid containing gas bubbles (non-foaming solution) had been shown to efficiently remove from surfaces both particles (Kondjoyan et al., 2009) and bacteria (Kriegel and Ducker, 2019). Recently, the efficiency of foam flow to remove *Bacillus* spores from stainless steel surfaces has been investigated by (Al Saabi et al., 2021) and the use of aqueous unstable foams has also been proposed for cleaning the sensitive surfaces of artistic and
cultural assets (Schad et al., 2021). It should also be noted that foam in static conditions, is widely
used in the food industry for cleaning open surfaces (floors, conveyor belts, worktops and
equipment) (Mierzejewska et al., 2014).

72 Aqueous foam could be considered as a non-Newtonian fluid consisting of gas dispersion in liquid 73 containing surfactants. Depending mainly on the liquid fraction, the foam could be ranked as wet 74 (small and spherical bubbles) or dry (large and polyhedral bubbles). A liquid-rich layer may be formed 75 near a solid surface due to the bubble migration, thereby giving a slip effect between the foam and 76 the wall (Tisné et al., 2003). Given this complexity, the rheological behavior of foam can follow the 77 shear-thinning Herschel-Bulkley model, which states that under strong mechanical stresses, foam 78 behaves as a viscous liquid, while under low stresses it behaves as a solid whose structure can be 79 deformed and return to its initial state (Dallagi et al., 2019). Moreover, very wet foam flow (liquid 80 fraction = 0.5) was found to clean hydrophilic spores better than dryer foam (liquid fraction between 81 0.3 to 0.4) (Al Saabi et al., 2021). However, the maximum of the shear stress that they used was 6 Pa, 82 and the work lacks the explanation of the cleaning process.

83 The aim of this study was to investigate the potential role of the foam velocity and their structure 84 (described by air/water content of the foam and size of air bubbles) in efficiently removing Bacillus 85 subtilis spores. In this respect, different flow conditions of wet foam (liquid fraction = 0.5), by varying 86 the shear stress (from 2 up to 13.2 Pa) and the foam structure, were carried out to test its efficiency. 87 All cleaning conditions were tested on two Bacillus spore strains which differ only in their hydrophilic/hydrophobic character. Another originality of this present work is the detailed physical 88 89 analysis of the dynamic behavior of foam and its structure and consequences at the wall to identify 90 the underlying mechanisms of the bacterial spores' removal and choose the foam flow condition 91 leading to optimal cleaning. The foam cleaning effectiveness was then compared to a standard CIP 92 procedure.

93

### 94 2. Materials and Methods

## 95 2.1. Foam preparation

96 The foam production prototype was built according to a previous work (Al Saabi et al., 2021). The 97 prototype is presented in Fig. 1. Briefly, foam was produced by pressurized air injection through a 98 porous medium (pore sizes ranging from 1 to 1.6 μm) into an SDS surfactant (Sodium Dodecyl Sulfate 99 (over 98.5% purity) dissolved in osmosed water (0.15% w/w) to avoid any unwanted interactions 100 with ions contained in the standard tap water.

101 The physical properties of the foaming solutions used were a 1 Bar atmospheric pressure and 102 20°C standard temperature, a 997.9 kg/m<sup>3</sup> density ( $\rho_l$ ), a 0.94 10<sup>-3</sup> Pa.s dynamic viscosity ( $\mu_l$ ) and a 103 26.2 10<sup>-3</sup> N/m surface tension ( $\sigma$ ). The foam quality ( $\beta$ ) describing the air/water content of the foam 104 was calculated as follows (Equation 1) where  $Q_g$  and  $Q_l$  are respectively the gas and liquid flow rates.

$$105 \qquad \beta = \frac{Q_g}{Q_l + Q_g} \tag{1}$$

106 A foam quality of 0.5 was chosen in this work by varying the shear stress and the capillary forces. 107 The capillary force of the bubble  $F_c = \sigma 2\pi r$  is defined as the force due to surface tension that acts 108 between the two halves of the bubble (Al-Qararah et al., 2013) with respect to the capillary pressure 109  $(P_c = \frac{\sigma}{r})$  (Schad et al., 2021), where r is the bubble radius. The balance between the viscous drag 100 forces versus surface tension forces of the foam can be represented by the non-dimensional capillary 111 number  $Ca = \frac{\mu_l v}{\sigma}$ .  $\mu_l$  and v are the dynamic viscosity of the liquid solution and the foam velocity, 112 respectively.

Therefore, the first part presents a detailed characterization of these foam flow conditions in order to understand the mechanism by which the foam detached the spores, to identify the key parameters and then to test these new foam conditions on cleaning of stainless-steel coupons contaminated by *Bacillus* spores. Table 1 (in Result section) summarizes all the FFC conditions (quality of 0.5) and CIP. The foam mean velocity ( $\bar{v}$ ) and the Reynolds number (*Re*) are determined from the following equations:

119	$Re = \frac{\rho_f.\bar{\nu}.d_h}{\mu_f}$	$d_h$ being the hydraulic diameter	(2)
120	$\bar{v} = \frac{Q_l + Q_g}{s}$	S being the duct section	(3)
121	$\rho_f = (1 - \beta). \rho_l + \beta.$	$ ho_g  ho_l$ being the liquid density	(4)
122	$\mu_f = \frac{\mu_l}{1 - \mu_l^{1/3}}$		(5)

The foam flow conditions varied from 1.5 to 13.5 cm/s ( $Q_g = Q_i$ ; are 4.5, 9, 13.5, 27, and 40.5 l/h for the cases 1, 2, 3, 4, and 5 respectively). The three first cases correspond to the conditions used in a previous work (Al Saabi et al., 2021). Cases 4 and 5 have the highest velocities, which were generated using three parallel generators. CIP has the highest velocity, while the mean shear stress was similar to that in case 4.

128

## 129 2.2. Foam characterization

### 130 2.2.1. Foam morphology

131 Foam flow was visualized in order to observe the bubble's displacement at the channel walls, for 132 a given interval. A Nikon D850camera, equipped with an AF-S Micro NIKKOR 60 mm f/2.8 G ED with a 133 resolution of 8256 x 5504 pixels, was used to record foam microtexture and its associated bubble size 134 distribution. For each foam flow condition, three images were captured then analyzed using 135 Piximètre 5.1 R1540 software. First, the chosen images were calibrated, then were subjected to 136 different image processing features (local threshold, filters...) to reduce the noise in order to better 137 observe its contours. The bubble size distribution was then calculated according to usual granulometric distributions. To evaluate the heterogeneity/uniformity of this distribution, different 138 139 parameters were identified such as F0.05, and F0.2 corresponding to the percentage of the bubbles 140 with a diameter of less than 0.05 and 0.2 mm respectively.

141 2.2.2. Thin film thickness

142 The conductimetry technique was used to measure the thickness of the thin film. This is based on 143 the measurement of the impedance of the liquid volume located between the liquid-foam interface and two outcropping electrodes at the duct wall. A generator provided an alternating sinusoidal
voltage to power these electrodes. The frequency was 50 kHz to avoid any polarization phenomena.
The calibration of the conductimetric probe was carried on using the calibrated Mylar sheets. The
estimated error was 8%.

148 2.2.3. Wall shear stress

When foam flowed, mean shear stress was calculated based on the measurement of the static pressure along the channel. The power loss measurements were carried out using a differential pressure sensor Shlumberger, type 8D. Two pressure outlets allow the connection of 2 manifold tubes placed over a scaled plate that measures on a length L of 1m the pressure drop placed after 1.5 m from the entry of the foam in the duct. Indeed, to allow steady state flow conditions, the test ducts were placed at intervals exceeding 80 times the hydraulic diameter of the vein inlet (i.e. 1.5 m). The mean wall shear stress can thus be computed as:

156 
$$\bar{\tau} = \frac{\Delta P}{\Delta x} \frac{d_h}{4}$$
 (6)

157 where  $\frac{\Delta P}{\Delta x}$  is the pressure gradient.

The local shear stress at the top wall where coupons were placed, was measured using the polarography method. This technique allows the determination of the local velocity gradient at the walls by measuring the mass transfer coefficient of the electrodes. Briefly, based on the simplified assumptions for the resolution of the convection-diffusion equation proposed by (Mitchell, 1965), the wall shear stress calculation was obtained from the velocity gradient (Lévêque solution and Sobolic correction). Details of the method used here were given previously (Tisné et al., 2003).

From a theoretical point of view, the polarographic signal is related to the slip layer thickness and foam velocities. The diffusion limit current can be represented as the Sherwood number form (*Sh*) or directly as the wall shear stress ( $\tau_w$ ). Spectral analysis of the time signals appears necessary in order to understand the underlying mechanisms and to trace back to the flow characteristics near the wall duct. The method used for spectral representation is based on the processing of the signal acquired by the Fourier transform (*FT*) with Matlab software. This was proposed by (loana et al., 2007). It expresses the frequency distribution of the amplitude, the power spectral density, as well as the fluctuation rate of the considered signals. These specters revealed peaks obtained at characteristic frequencies corresponding to the passage frequencies of bubbles respective to each foam condition and their harmonics.

174 It was found that the SDS did not support the addition of electrochemical products (Ferri-175 ferrocyanide potassium) for the use of the polarography method. In fact, a reaction other than 176 oxidation-reduction was produced, destabilizing the foam. For this reason, we replaced the SDS with 177 a non-anionic surfactant: Lauramine oxides (Ammonyx<sup>®</sup> LO (Stephan) = 0.8 ml/l). To ensure the 178 presence of the same phenomena for the same flow conditions, the thickness of the liquid film and 179 of the conductivity signals produced with that of the SDS solution was verified. The same behavior and signal variation with SDS or Ammonyx® surfactant was observed, which is in line with previous 180 181 works (Tisné et al., 2003), where another surfactant (Amonyl<sup>®</sup>) replaced the SDS for the same reason.

182

## 183 **2.3. Bacterial strains and materials**

184 Two Bacillus subtilis strains were used throughout this study. Bacillus subtilis PY79 (subsequently 185 named Bs PY79) is a laboratory strain producing hydrophilic spores surrounded by a mucous layer 186 called crust. A Bs PY79 recombinant strain deleted in spsA (Dubois et al., 2020), subsequently named 187 Bs PY79 spsA, which produces hydrophobic spores modified in their crust properties, was also used 188 throughout this study. Both Bacillus strains were tagged with green fluorescent proteins to make the 189 spores more fluorescent. Spores were produced as previously described (Faille et al., 2019) on Spo8-190 agar at 30°C. When over 95% of spores were obtained, they were harvested by scraping the surface, 191 washed five times in sterile water, and stored in sterile water at 4°C until use. In order to limit the 192 presence of spore aggregates, the spore suspensions were subjected to an ultrasonication step 193 (Bransonic 2510E-MT, Branson Ultrasonics Corporation, USA) before each experiment.

All experiments were carried out on AISI 316 stainless steel with pickled (2B) finish coupons (provided by APERAM, Isbergues, France) in the form of rectangular coupons (45 mm × 15 mm). In 196 order to have surface properties similar to those used in the dairy industry, the stainless-steel 197 coupons were subjected to a conditioning procedure consisting of 15 runs of soiling and cleaning 198 both to mimic what might occur in food processing lines and to stabilize their surface properties. The 199 soiling step was achieved by immersing the coupons in milk in semi-skimmed reconstituted milk (150 200 g/I milk powder) for 30 min at room temperature. After a quick rinse, soiled coupons were cleaned 201 by immersion in sodium hydroxide at 0.5% w/w for 30 min at 70°C. Before each experiment, coupons 202 were cleaned using an alkaline detergent (RBS T105, Traitements Chimiques des Surfaces, France). 203 Each coupon was first rubbed vigorously with pure RBS and then immersed for 10 minutes in 5% RBS 204 at 60°C. Coupons are thoroughly rinsed for 5 min with tap water, and then with softened (reverse 205 osmosis) water for another 5 min. One day before experiments, coupons were sterilized in a dry heat 206 oven at 180°C for 1h.

The hydrophilic character of the cleaned and disinfected coupons was estimated using a goniometer (Digidrop, GBX, France). A water contact angle ( $\theta_{water}$ ) of 66° was measured, which reflects the relatively hydrophilic character of this material.

210

## 211 2.4. Surface soiling and cleaning

212 For both strains, the concentration of the spore suspensions was set at  $10^8$  CFU/ml. Five 1  $\mu$ l 213 droplets of the spore suspension were placed on each coupon and evaporated for 1h at 30°C. The 214 procedure was schematically illustrated Fig. S1. Some coupons were analyzed to control the initial contamination level (5 10<sup>5</sup> spores per coupon) and to observe the droplet architecture under the 215 216 microscope (Fig. S1 in the Supporting Information). Then, sets of 4 coupons were placed inside a 217 rectangular stainless-steel duct (1.5  $10^{-2}$ m high × 1  $10^{-2}$ m width), subjected to a cleaning procedure 218 and analyzed again for residual contamination (three coupons) for enumeration and deposit 219 architecture (one coupon).

In this work, as previously (Al Saabi et al., 2021), the foam cleaning phenomenon was evaluated
at the top wall because the wall shear stress and the thin liquid film thickness would be considered to

show constant variations in this zone according to previous work (Tisné et al., 2004, 2003). This
position was selected in order to facilitate the interpretation of the role of foam flow in the cleaning
process.

225

226 The resistance of the spore deposits to a cleaning procedure was performed as follows. The test 227 ducts were connected to the pilot rig and then were subjected to the cleaning process. Experiments 228 were carried out on foam flow cleaning (FFC) under different conditions (by varying the foam shear 229 stress and the foam/bubble structure) for 15 and 35 s, 1, 3, 5, 10, and 20 min. Other experiments 230 were carried out on cleaning in place (CIP) under the same conditions to compare the relative 231 effectiveness of FFC and CIP. Details are provided in the Results section. After the cleaning process, 232 the coupons were removed from the test duct, then rinsed by immersion in sterile ultrapure water to 233 remove the rest of the foam. In order to investigate the possible role of the presence of SDS 234 surfactant on spore detachment, further experiments were performed in static conditions by dipping 235 a set of coupons in pure water or in SDS 0.15% (w/w) for 20 min (data not shown). It should be noted 236 that Bs spore incubation in SDS did not result in any significant viability loss.

237 To determine the quantity of spores on the coupons before and after cleaning, each coupon 238 surface was sampled using a dry cotton swab (Copan, Brescia, Italy). The swabs were placed in a tube 239 containing 5 ml of sterile ultrapure water and vortexed for 1 min at 2400 rpm. The suspensions 240 containing the detached spores were enumerated on TSA (Tryptone Soy Agar, Biokar, France) after 241 48h at 30°C. For microstructure examination, the observation of some deposits was performed 242 before and after cleaning, using an epifluorescence microscope (Zeiss Axioskop 2 Plus, Oberkochen, 243 Germany). The observations were first made at low magnification (x50) in order to observe the 244 whole deposit resulting from the drying of a droplet. Further observations were carried out at higher 245 magnification (x400) to determine the fine structure of the densest parts of these deposits (clusters, 246 peripheral rings...), including any 3D organization.

After cleaning, the deposits were mostly weakly fluorescent, which made them difficult to observe. It was therefore necessary to stain them before observation. For that purpose, dried coupons were stained with orange acridine (0.01%) for 15 min. In order to ensure that the deposits were not removed during this staining step, some deposits were observed before and after the staining step. No significant difference was observed.

252

253

## 254 2.5. Kinetics modeling

The removal kinetic profiles of the percentage of residual spores were fitted according to previous work (Bénézech and Faille, 2018), using GInaFIT. A two-phase model was selected (Equation 7), which presents the existence of two subpopulations characterized by different removal behaviors.

258 
$$\frac{N}{N_0} = f e^{(-k_{max_1}t)} + (1-f)e^{(-k_{max_2}t)}$$
 (7)

Where  $N_0$  is the initial bacteria count, f is the poorly adherent fraction of the population (less resistant to detachment), (1-f) is the more adherent fraction, and  $k_{max1}$  and  $k_{max2}$  [s<sup>-1</sup>] are the specific detachment rates of the two sub-populations.

262

## 263 2.6. Statistical analysis

All the experiments were carried out in triplicate. Statistical analyses were performed by general linear model procedures using SAS V8.0 software (SAS Institute, Gary, NC, USA). For each cleaning condition, variance analyses were performed to determine the possible role of the shear stress and the foam structure on the removal kinetics of spores (f,  $k_{max1}$ , and  $k_{max2}$ ) and the effectiveness of the cleaning by foam flow compared to the CIP process.

269

## 270 **3. Results**

In this study on the detachment of adherent spores by foam, it was essential to dispose of a
 procedure allowing very reproducible contamination of the surfaces, both in terms of contamination

273 rates and in terms of the structure of the deposits. For this reason, we chose to contaminate the 274 materials with droplets containing *Bacillus* spores and to let these deposits dry as previously 275 described (Deleplace et al., 2021). For the type of foam (more or less dry), we chose to use a so-276 called wet foam (foam quality of 0.5) because recently published work had shown that dryer foams 277 (liquid fraction from 0.4 to 0.3) were less effective (Al Saabi et al., 2021). We used SDS as model 278 surfactant system to produce the foam. In order to identify the key parameters affecting the 279 efficiency of foam cleaning, we first performed a series of physical measurements of foams obtained 280 under different conditions, and we measured the effect of these parameters on the cleaning 281 efficiency in terms of spores' detachment during cleaning procedure and patterns deposition after 282 cleaning.

283

## 284 **3.1. Foam flow characterization**

## 285 3.1.1. Bubbles size distributions

286 Fig. 2 shows top views of the foam at the top wall for the different flow conditions. For all foam 287 flow conditions, the bubbles had a roughly spherical shape and were separated by thick liquid films 288 (lamellae). Different foam morphologies were observed. Indeed, when the velocity increased from 289 1.5 to 9 cm/s, a clear decrease in the bubble size was observed. However, no additional decrease 290 could be observed between 9 and 13.5 cm/s. Bubble sizes were then measured for these different 291 conditions and the curves of the cumulative fraction of the bubble size distributions are presented in 292 Fig. 3. First, the distribution of the bubble size in all conditions was fitted by a lognormal function ( $R^2$ 293 ranging from 0.89 to 0.97) to better visualize the differences between the conditions tested. 294 Variation of velocity causes a change in foam structure and bubble size distribution. As the foam 295 velocity increased (Table 1), the mean diameter of bubbles decreased (from 0.34 down to 0.18 mm, 296 for foam at 1.5 and 9 cm/s respectively) and the distribution became more homogeneous, except for foam at 13.5 cm/s where the mean diameter was 0.21 mm. While the number of small bubbles (less 297 298 than 0.05 mm in diameter) was very close at 9 and 13.5 cm/s, intermediate-sized bubbles (between

299 0.2 and 0.6 mm in diameter) were more abundant at 9 cm/s than at 13.5 cm/s. Therefore, the 300 cumulative bubble size number at 0.2 mm (*F02*) was 0.68 and 0.60 respectively for 9 and 13.5 cm/s, 301 as indicated in Table 1 and were found to be statistically different (Pr < 0.0001). These differences 302 were much greater when comparing the highest velocities to the lowest ones. Indeed, considering 303 bubbles with a diameter  $\leq 0.05$  mm (see fraction *F005*, Table 1), at 13.5 cm/s the cumulative fraction 304 was more than 20 times higher than that for the lowest foam velocities, at 1.5 cm/s.

305 3.1.2. Upper liquid film thickness

306 Foam is not in direct contact with the wall surface, but bubbles are slipping over a thin liquid film 307 layer (a few micrometers width). Variations of the liquid film thickness (from 11 up to 13  $\mu$ m) 308 measured at the top wall are shown in Fig. 4 (left) for three foam mean velocities of 1.5, 3, and 4.5 309 cm/s, bearing in mind the limitations of our experimental set-up essentially due to the inability to 310 access sufficiently high foam flow rates as it was possible for the cleaning experiments. The film 311 thickness was affected by the bubble passage as shown in Fig. 4 (left) with an example of the 312 observed temporal evolution of the signal at the middle of the test section. It should be noted that 313 whatever the position of the probe the signal was observed as similar either on the sides or in the 314 middle of the duct. In the three cases, conductimetry signals presented high fluctuations reflecting 315 the non-uniformity of this liquid film. Indeed the signals showed an instantaneous variation of 72 and 316 54% relative to the mean thickness value which varied from 5 to 21  $\mu$ m, and from 8 to 22  $\mu$ m for 317 foam at 1.5 and 4.5 cm/s, respectively. In addition, the profiles obtained appeared to be very 318 different within the three cases. The higher is the velocity the higher is the frequency of the signal 319 (variation from 45 to 72.5 Hz for cases of foam at 1.5 and 4.5 cm/s, respectively), and the lower is the 320 amplitude (variation from 9.8 up to 6  $\mu$ m, respectively). These variations were directly linked to the 321 frequency of passage of the bubbles and to the size variations leading to the shape of the signal. 322 Indeed, the maximum thickness of the liquid film (peaks of the signal) corresponds to the film 323 thickness between two following bubbles passage, minimum thickness (signal hollows) 324 corresponding to the presence of the bubble, the bigger bubble the lower film thickness could be

325 stated from Fig. 2. It should be noted that conductimetry results, for both tested surfactants SDS and 326 Ammonyx, have shown the same instantaneous evolution and presented the same phenomena at 327 the walls. A variation of 61 and 42% relative to the mean thickness and an increase of frequency from 328 37.5 to 77.5 Hz for cases of foam at 1.5 and 4.5 cm/s respectively, are shown in Fig. S2 in the 329 Supporting Information.

330 3.1.3. Wall shear stress

331 To highlight the effect of the passage of a bubble on the shear rate at the wall, temporal signals 332 from the frequency response of the polarography probe located in the middle of the channel have 333 been plotted (Fig. 4, right).

334 The wall shear stresses on the upper region of the duct appeared larger and approximately 335 double the value of the average stress measured on the section (data not shown). This is in relation 336 to the low thickness of the liquid film at this location. The average shear stress increases slower than 337 the foam flow velocity but increases from 4.3 Pa at 1.5 cm/s to 7.2 and 8.1 Pa for velocities of 3 and 338 4.5 cm/s, respectively. More interesting were the temporal variations of the wall shear stress at the 339 top wall as shown in Fig. 4 (right), the signal varies periodically with a more homogeneous pattern 340 with velocity increase but a decrease in amplitude is clearly visible at 4.5 cm/s. The rate of fluctuation 341 of the wall shear (RFS) related to the signal amplitude decreases as the foam velocity increases 342 reaching 957, 688, and 491% for the three foam velocities tested, 1.5, 3, and 4.5 cm/s respectively. 343 The passage of bubbles, as seen in the film thickness variations, causes these variations in local shear 344 stress. These fluctuations are therefore characterized by a frequency that corresponds directly to the 345 frequency of bubble sliding on the upper wall.

Observing the spectral densities of the wall velocity gradient, all the curves had almost the same shape of which spectrums obtained with the probes gave a low energy level (Fig. S3 in the Supporting Information). The frequencies of bubbles are 20, 39, and 45 Hz for the first three cases respectively. For the other foam conditions, the frequency and the *RFS* were determined by the extrapolation of these experimental results and the theoretical correlations between the frequency and foam parameters (where  $f_{bi} = \frac{\beta . v}{d_{bi}}$ ,  $A_{bi} = \frac{v}{f_{bi}}$ , and  $d_{bi}$  are the frequency, amplitude, and diameter of i bubble). As was expected, foam at higher velocities (9 and 13.5 cm/s) had the highest frequency (76 Hz) with a strong decrease in the wall shear fluctuation rate (130% for foam at 9 cm/s) (Table 1). All results are reported in Table 1.

355

## 356 3.2. Spores' detachment

## 357 **3.2.1.** Resistance to cleaning procedures

358 The cleaning kinetics for the two bacterial spores are presented in Fig. 5 based on the evolution 359 of the mean values of the logarithmic reduction of adherent spores with the cleaning time. In all 360 cases, the detachment curves clearly showed two distinct phases. Both phases appeared to be exponential and therefore were quite accurately described by the biphasic model, with  $R^2$  ranging 361 362 from 0.84 to 0.97. Whatever the flow conditions, the first detachment phase lasted less than 1 min, 363 and the duration of this phase depended on the flow conditions of the foam (shorter at low velocity, 364 for the two spores), but especially on the strain. At 1 min, the resulting log reduction of the adherent 365 spores during this first phase ranged between 1.94 and 2.6, and between 0.95 and 1.33 for Bs PY79 366 and Bs PY79 spsA, respectively. The second phase was much longer and would probably continue 367 over 20 min. The detachment rate of the hydrophilic spores (Bs PY79) during this second phase was 368 quite low at 1.5, 4.5, and 13.5 cm/s resulting in an additional 0.15-0.28 log reduction of the adherent spores, whereas at 9 cm/s the reduction reached around 1.15 log. A similar trend was observed for 369 370 the hydrophobic spores with 0.2-0.3 log reduction at 1.5, 4.5, and 13.5 cm/s and 0.45-0.57 log 371 reduction at 9, and 13.5 cm/s.

The role of the flow parameters and strains on the three kinetics parameters (f,  $k_{max1}$ , and  $k_{max2}$ ) was then investigated. When comparing the two strains, the influence of the strains on the removal efficiency during the first phase represented by the two parameters f (poorly adherent fraction of the population) and  $k_{max1}$  (first kinetics constant rate) was highly significant as confirmed by the variance analysis (Pr < 0.0001 for both parameters). According to Tukey's groupings, f was significantly affected by the strain, with mean values of 0.92 and 0.998 for Bs PY79 *spsA* and Bs PY79 respectively as shown in Fig. 6. The initial detachment rate ( $k_{max1}$ ) also was largely dependent on the strain (from 1.7 up to 3.4 times higher for hydrophilic spores), while the second detachment rate ( $k_{max2}$ ) was not affected and was quite similar for the two strains (Fig. 6, no significant differences).

381 The role of the flow parameters as the shear stress (mean value, frequency, and RFS) and/or to 382 the distribution of bubbles (mean diameter and FO2) on the three kinetics parameters (f, kmax1, and  $k_{max2}$ ) was then investigated (Table 2 and Fig. 7). Whatever the parameter in relation to the flow and 383 384 the bubble size arrangement the Tukey's grouping tests gave the same ranking between the 4 385 velocities tested (Table 2) and can be visualized clearly in Fig. 7, the two groups of parameters as the 386 ones related to the wall shear stress and the ones related to the bubbles' size influencing, in the 387 same way, the removal kinetics. However, differences between the two strains could be stated. For 388 Bs PY79, the first kinetics phase rate constant  $k_{max1}$  was similar for all the studied cases (no difference 389 according to the Tukey' test) but an effect was observed on f with a significant difference between 390 the lowest and the highest velocities (1.5 and 4.5 compared to 9 and 13.5 cm/s) leading however to 391 an improvement while increasing the mean shear stress (Fig. 7, black dots). Conversely, kmax2 was 392 significantly higher in the case of foam at 9 cm/s comparatively to the other cases whatever the 393 strains. Some differences could be seen with the hydrophilic spores, according to Tukey's grouping in 394 Table 2 and Fig. 7, red dots. Indeed, the maximum foam velocity tested induced a significantly 395 higher value of f (Pr < 0.0001), a phenomenon more marked than with hydrophilic spores, as for 396  $k_{max1}$ , it was significantly higher at this velocity than at the lowest velocity tested (Pr < 0.013). 397 However, *k<sub>max2</sub>* variation trend with the velocity was similar to those of the hydrophilic strain and also 398 highly significant (Pr < 0.0007), the velocity of 9 cm/s being the best option. It should be noted that it 399 was not the highest mean shear stress tested (13.2 Pa), but it has the lowest bubble sizes (mean 400 diameter of 0.18 mm and FO2 of 0.68).

401 Cleaning with foam was then compared to a CIP procedure (pink and black curves in Fig. 5) using 402 the SDS solution at a shear stress of 9.8 Pa, shown to be the most efficient condition with foam. 403 Concerning hydrophilic spores, the first detachment's phase in the case of CIP was shorter than for 404 foam cleaning and the removal was 0.5 log CFU less efficient. In addition, very few spores were 405 detached in the second phase, resulting in about 2 log CFU reduction after 20 min cleaning compared 406 to around 4 log CFU with foam. The same trend was observed with hydrophobic spores but at a 407 lesser extent, the use of foam allowing an improvement in detachment efficiency of only 0.5 log in 20 408 min compared to CIP. The variance analysis confirmed that whatever the strains significant 409 differences could be observed between CIP and FFC considering the two kinetics parameters f and 410  $k_{max2}$  (Pr < 0.034 and less than 0.001 respectively),  $k_{max1}$  being apparently not significantly modified 411 (Fig. S4 in the Supporting Information).

## 412 **3.2.2.** Deposition patterns before and after cleaning

413 In order to observe whether the different areas of the dried drops were more or less easily 414 removed during the cleaning procedures, the dried deposits were observed before cleaning and after 415 15 seconds and 20 minutes of cleaning. Two different deposition patterns were observed in 416 epifluorescence microscopy (Fig. S4 in the Supporting Information). Whatever the spores, a regular 417 round peripheral ring, often referred to as the "coffee ring" was clearly observed. For Bs PY79 spores, 418 this external coffee ring shape was thick with relatively few adhered spores within the ring. 419 Depending on the trial, the internal deposition was more or less regular, as if the contact line 420 remained fixed or detached and then reattached, resulting in one or more small contaminated areas. 421 Conversely, with the hydrophobic spores Bs PY79 spsA, a diffuse peripheral deposit was often 422 observed, within which single spores but also spore clusters are present over the whole surface.

After cleaning, the observation of the residual patterns shown a good agreement with the enumeration results (Fig. 8). First, Bs PY79 spores seemed to be removed in a higher extent compared to Bs PY79 *spsA* spores. In the case of Bs PY79 spores after 15 s, the peripheral ring was partially removed by the cleaning procedure and this was particularly noticeable at 13.5 cm/s, but especially at 9 cm/s. After 20 min of cleaning, the ring had completely disappeared at 9 and 13.5 cm/s and to a lesser extent at 4.5 cm/s. However, spores were still visible on the inner area of the 429 deposits except for foam after 20 min at 9 cm/s. Regarding the cleaning of Bs PY79 spsA spores, 430 spores, and small clusters seemed to be greatly removed while the larger clusters seemed to have 431 better resisted the detachment, and the peripheral ring was partially removed only after 20 min. For 432 both strains, removal was visibly greater with foam flow at 9 cm/s than CIP. At 15 seconds, the 433 cleaning of the deposits of hydrophobic spores Bs PY79 spsA seemed to induce a slid of the 434 contamination along the surface in the direction of the flow showing a quite scattered pattern visible with foam flow cleaning and CIP (blue arrows on Fig 8). This sliding phenomenon was absent in the 435 436 case of hydrophilic spores Bs PY79.

437

438 **4.** Discussion

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440 Foam has been considered as a non-Newtonian fluid and the Herschel–Bulkley model can predict 441 its flow behavior as previously described (Dallagi et al., 2019). At low shear stress (2.2 Pa) the foam 442 flow behaved as a plug flow in one direction, the velocity of the foam being uniform over the entire 443 section as a 1-Dimensional (1D) flow regime. When the wall shear stress exceeded 4.2 Pa the plug 444 behavior remained only at the top side of the duct (2D flow regime), and at shear stresses greater or equal to 5.9 Pa, the foam became completely sheared (3D flow regime) (Al Saabi et al., 2021). The 445 446 rheological behavior of foams is controlled by processes that occur within the microstructures such 447 as the rearrangement of bubble size, liquid film, and Plateau borders. A slight modification on the 448 pressure or the shear stress applied could lead to deformation on the foam state. In previous studies 449 on the rheology of emulsion and foam, authors have demonstrated the deformation of the bubble as 450 a function of the capillary number (Llewellin and Manga, 2005). Their results show that as the 451 capillary number increases, bubbles start to deform and the film separating bubbles either thicken or 452 shrink trends to maintain the minimal surface energy. However, this deformation can result in developing anisotropy in foams properties if the capillary number exceeds a critical value which is 453

454 not the case here: for all studied cases *Ca* was lower than 0.1 and the bubbles remained almost455 spherical with isotropy of the foam nearly preserved.

456 It should be noted that except for the three foam conditions at 1.5, 3, and 4.5 cm/s obtained 457 using 1, 2, or 3 generators with the same surfactant, flow rate, and air pressure content, the increase 458 of the foam velocity up to 9 and 13.5 cm/s was obtained by increasing the surfactant flow rate and 459 air pressure and content in the foam generators. It was demonstrated that the pressure significantly affects the foam texture, and its stability with a potential change in bubble size distributions, liquid 460 461 drainage, and bubble coalescence (Al-Qararah et al., 2013), a reduction in the average radius of 462 bubbles was observed with the increase of the pressure. These authors reported that the bubble size 463 decreases with the increasing of the rotation speed up to a certain limit under axial mixing in a 464 Becher. At a rotation lower than 5500 RPM (revolution per minute), the foam remains stable since 465 the air is well homogenized in the surfactant solution. Conversely, above 5500 RPM, because of the 466 saturation of air content, the interactions between bubbles were more numerous which favored 467 their coalescence. Thus, their size increases, resulting in a decrease in interfacial stability. The 468 observation of the foam at the highest velocity (13.5 cm/s) allowed to suggest that the large bubbles 469 are more likely the consequence of a movement inside the foam (3D flow regime) allowing the larger 470 bubbles indeed less dense to rise towards the upper wall rather than a coalescence of the bubbles 471 explained by the high speed of the foam. This behavior was probably enhanced by the fact that the 472 foam is quite wet with a quality  $\beta$  of 0.5.

Furthermore, the conductimetry and polarography results (Fig. 4) demonstrated a significant impact of the velocity of the bubbles and their sizes on the fluctuation of the liquid film thickness between bubbles and the top wall and on the local wall shear stress increases significantly between two consecutive bubbles (Plateau borders). These variations would directly affect bacteria present at the surface influencing the removal phenomenon. The displacement of the bubbles leads to sudden increases in the Sherwood number due to shear stress peaks, appearing in a harmonic manner directly related to the bubble passage frequency. Comparing foams at 1.5, 3, and 4.5 cm/s, as the 480 velocity increased, the bubbles became smaller with an F02 of 0.29, 0.36, and 0.44. The frequency of 481 the shear stress increased in parallel from 20 up to 45 Hz while amplitude decreased in agreement 482 with previous works (Tisné et al., 2003). Hence, foam at 9 cm/s having a higher amount of small bubble sizes presented the lowest RFS (RFS of 130 %) with the high frequency (f of 70 Hz) related to 483 484 the thin liquid film height variation measurements. The peaks of the shear stress were clearly 485 observed at the two ends of the bubbles, the bubble front, and the back. As an example, the wall shear stress value between the bubbles is about 15 Pa while it remains lower than 2 Pa under the 486 487 bubbles, for foam at 4.5 cm/s. These values corresponded therefore to the maximum and minimum 488 film thickness observed, respectively. In the literature, CFD simulations, as well as experimental 489 studies in two-phase gas-liquid slug flow, e.g. in capillary membranes (Kumari et al., 2018), suggested 490 that the major contribution to the wall shear stress came from the bubbles ends and could be 491 attributed to the Laplace pressure difference caused by the surface tension. (Tisné et al., 2004) 492 revealed that it was possible to relate foam flow study to studies of a long bubble moving in 493 capillaries, and they observed that the bubbles moved as flattened bubbles in a Hele-Shaw cell 494 confirming the localization of the shear stress on the Plateau borders.

495 When compared with spore removal kinetics, parameters in relation to the shear stress and 496 bubble size variations seemed to play a significant role in the cleaning efficiency of foam flow. For 497 both Bs strains, the increase of the mean shear stress roughly leads to an improvement in the 498 cleaning efficiency. However, it was not true for the foam at 13.5 cm/s, which showed a reduction in 499 cleaning efficiency, compared to 9 cm/s. This result suggests that other parameters than the mean 500 shear stress were involved in the FFC efficiency and therefore the variation of the local wall shear 501 stress (WSS) as discussed above in relation to the size and the velocity of the bubbles could play a 502 major role. These results are consistent with previous works in literature. Silva et al. highlighted the 503 significant role of the WSS fluctuation on the cleaning efficiency using pulsed flow (Silva et al., 2021). 504 Results showed that shorter periods promoted higher removal of raw milk cells due to the 505 occurrence of the annular effect and its important frequency. A higher frequency of peaks WSS was

506 needed to make turbulent pulse flow effective to remove tomato deposits (Absi and Azouani, 2018).
507 The authors explained obtained results by the fact that at higher frequency the surfaces are
508 submitted more often to the phenomenon caused by peaks. Hence, as the foam has a higher amount
509 of small bubbles, the contamination will be more subjected to the pick of WSS and perhaps also to
510 the viscous dissipation in the narrow Plateau borders between consecutive bubbles since smaller
511 bubbles possess a higher Laplace pressure.

512 Whatever the FFC conditions, the removal kinetics showed a strong detachment rate at the 513 beginning of the cleaning process (less than 1 min), followed by a second kinetic phase described by 514 a slow detachment rate. Previous works on biofilm removal under CIP (Bénézech and Faille, 2018) 515 and bacteria spores under different foam qualities (Al Saabi et al., 2021) have reported similar 516 observations. Neither the shear stress nor the bubble size has affected the first kinetic phase rate of 517 removal with Bs PY79 spores, while high shear stress induced low  $k_{max1}$  with Bs PY79 spsA. For both 518 strains, as the shear stress increased, the f parameter increased (the spore population detached 519 more easily). The synergy between high shear stress and small bubbles (foam at 9 cm/s) would 520 explain the highest  $k_{max2}$  and a better cleaning efficiency. To prove or disprove this suggestion, a new 521 foam flow condition was produced at 9 cm/s, using a single generator. The air pressure in the 522 generator in this condition was 3 times more than that for the 9 cm/s foam, produced by three 523 parallel generators. As was expected, the mean shear stress was 9.8 Pa, while the bubble size 524 increased (mean diameter of 0.19 mm with a presence of large bubbles (Fig. S5 in the Supporting 525 Information). When comparing the second detachment phases, foam with the smallest bubbles was 526 more efficient for PY79 spores, by 1 log CFU over 20 min.  $k_{max2}$  was 1.5 times higher (with 9.8 Pa foam 527 and small bubble size) than that in the alternative condition tested. For spsA spores,  $k_{max2}$  seems to 528 be similar in these two conditions.

According to (Kondjoyan et al., 2009), the role of capillary forces is not negligible and depends on the size of the bubbles and their speed. In (Al Saabi et al., 2021), the detachment of the hydrophilic spores of *Bacillus amyloliquefaciens* 98/7 occurred at bubble velocities of around 3 cm/s and 532 dramatically decreased at 5 cm/s in line with the work of (Kondjoyan et al., 2009), where the greatest 533 detachment rate was observed, as illustrated by high  $k_{max1}$  constant rate values. In this work, no 534 visible changes could be seen for the first kinetic phase with the increase in the foam velocity from 535 9.5 up to 13.5 cm/s, but the second kinetic phase showed a significant decrease in the removal 536 efficiency. For hydrophobic particles, according to (Kondjoyan et al., 2009), the detachment rate 537 would start at higher bubble velocities. In this work, the increase in the foam velocity indeed was 538 accompanied by an increase in the detachment efficiency during the first minutes of the foam 539 cleaning, which corresponds to an increase in the f parameter, resulting in a more rapid and efficient 540 spore removal.

541 Other phenomena would occur under the combination of the drainage and the imbibition effects 542 (Fig. 9). It is known that under gravity and capillarity, the liquid flows through drainage networks 543 called Plateau borders. This phenomenon has been demonstrated in previous works (Maestro et al., 544 2013). Thus, as discussed before, the spores present in the thin liquid film between the bubble and 545 the solid surface are prone to be absorbed within the Plateau borders under the imbibition effect. 546 This mechanism is driven by the capillary forces. Indeed, these Plateau borders have a curvature that 547 creates a capillary under-pressure in the liquid phase. Owing to this pressure difference, the foam 548 was able to absorb the spores (which were in small size,  $1 \mu m$ ), just as a sponge would. It has been 549 proven in the literature that the imbibition mechanism was highly dependent on the bubble size, 550 which is the strongest for foam with a smaller bubble size (Maestro et al., 2013; Mensire et al., 2015). 551 On the other hand, such a mechanism could help in understanding the difference in the results 552 obtained by the two strains of spores. Unlike hydrophilic spores which can be easily trapped by the 553 liquid film, absorbed into the Plateau border and then flushed out of the duct, the hydrophobic 554 character of Bs PY79 spsA gives spores a strength to oppose the imbibition effect. They mostly 555 remain in the liquid film and then due to bubble fluctuations, some of them are able to re-adhere to 556 the solid surface. To highlight this phenomenon, the ability of hydrophobic spores to adhere to 557 sterile coupons was tested at a distance of around 10 cm downstream of the contaminated area. 558 Whatever the flow conditions, results showed that 1.3% of the detached hydrophobic spores after 20 559 min of cleaning were able to re-adhere to the surface. Comparatively, the re-adhesion of hydrophilic 560 spores was found to be negligible, at less than 0.1% of the detached spores after 20 min. Otherwise, 561 due to the continuous fluctuations of the shear stress and the liquid film thickness induced by the 562 bubble passage, some spores that were strongly attached to the surface were pushed to other 563 locations, thereby causing a sliding phenomenon. This phenomenon was more visible in the case of hydrophobic spores than for hydrophilic ones. This could be explained by the strongly hydrophobic 564 565 character of Bs PY79 spsA. Recently, Schad et al. have shown that foams with small bubble sizes are 566 more efficient in cleaning glass surfaces contaminated with fluorescent oil (Schad et al., 2021). The 567 authors explained the cleaning by an interplay of imbibition, drainage, and wiping mechanisms. 568 Imbibition and drainage of oil into the Plateau borders of foams were found to be the strongest with 569 small bubbles, which is in concordance with our results. However, the wiping mechanism, i.e., 570 shifting of the contact line between foam, oil and glass was found to be more efficient with unstable foam than with a stable one. This is not surprising since they used static foam. Given similar 571 572 conditions, unstable foams clearly clean better than stable foams, being continuously in motion and 573 accelerating the re-arrangements of their bubbles with time which leads, as they explained, to 574 violent wiping motions between the foams and the interface. Consequently, these are not 575 completely irrelevant with our work since in our study foams were used in a dynamic form and the 576 movement/flow was imposed, leading to forceful wiping. However, the instability of foam can reduce 577 the cleaning efficiency through the generation of large bubbles (as observed in the case of foam at 9 578 cm/s).

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580 5. Conclusion
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In this study, we analysed the ability of wet foam flow to remove *Bacillus subtilis* spores from dried droplets on stainless steel surfaces. Two strains, *Bs* PY79 and *Bs* PY79 *spsA*, of different 584 hydrophobicity were used. Regardless of the strain, the foam flow showed a better removal capacity 585 than conventional cleaning in place at the same average shear stress. Optimisation and improvement 586 of cleaning efficiency using foam flow requires a quantification of the key parameters, such as bubble 587 size, liquid film thickness and wall shear stress. Analyses obtained from polarography, conductimetry 588 and bubble size quantification have shown that the magnitude and frequency of fluctuations in wall 589 shear stress and liquid film thickness are related to bubble size and flow velocity. These phenomena 590 are most likely to be the cause of spore removal, as described in the literature, although to a lesser 591 extent, for cleaning with liquid solutions. In contrast, with foam, the particularity of fluctuations at 592 the wall (liquid film thickness and shear stress) in relation to the bubble sizes and their velocities, 593 combined with other probable phenomena such as imbibition and drainage of the liquid film within 594 the foam, favours not only the release of contaminants, but also their entrapment. This leads to the 595 conclusion that a combination of high WSS, and the presence of small bubbles (<0.2 mm) indeed 596 improves surface biocontamination cleaning efficiency. Further studies would be interesting to test 597 the ability of this process to clean actual industrial equipment, which is frequently of complex design. 598 Further research could also be carried out into the removal of other types of microbial fouling by 599 these flowing foams, such as into biofilm contaminations developed on the surfaces of materials 600 other than steel, such as polymers, which are sometimes rougher and often more hydrophobic. The 601 implementation of such a foam cleaning technology for CIP cleaning would request an effective foam 602 removal during the rinsing step. Future work could be envisaged to solve this problem using stimuli-603 responsive surfactants allowing at request a complete disappearance of the foam.

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608

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Fig. 1. Experimental set-up for studying the foam cleaning in place process.



**Fig. 2.** Top view of foam at the wall of the transparent pipe (the entire width, 1 cm) for the different foam flow conditions.



**Fig. 3.** Cumulative fraction of the bubble size distribution for all the foam flow conditions. The markers present the brut data and the lines represent the Lognormal model fitted to each condition.



**Fig. 4.** Example of the evolution of wall film thickness (left) and wall shear stress (right) given by the temporal signals of conductimetry and polarography over the top wall for the three cases, respectively. The values indicated in the images show the mean thickness and the mean wall shear stress during 40 seconds of acquisition.



**Fig. 5.** Removal kinetics of *B. subtilis* PY79 (left) and *B. subtilis* PY79 *spsA* spores under the different flow conditions. Markers showed the data of three repetitions, and the line represents the biphasic fitted model fit to each condition: 1.5 cm/s (sky square), 4.5 cm/s (purple circle), 9 cm/s (pink triangle), 13.5 cm/s (red diamond), and CIP (black triangle).



**Fig. 6.** kinetics parameters f, K<sub>max1</sub>, and K<sub>max2</sub> for Bs PY79 (black color) and Bs PY79 *spsA* (red color) spores under different flow conditions.



**Fig. 7:** Variations induced on the kinetics parameters f,  $K_{max1}$  and  $K_{max2}$  for Bs PY79 (black color) and Bs PY79 *spsA* (red color) spores under different flow conditions: 1.5 cm/s (square), 4.5 cm/s (circle), 9 cm/s (triangle), 13.5 cm/s (diamond), and CIP.



**Fig. 8:** Residual deposition patterns of droplet contaminated with Bs spores after complete evaporation and further subjected to 5 seconds and 20 minutes of FFC and CIP procedures, examined by epifluorescence with a Zeiss Axioskop 2 plus microscope (x50). Blue arrows indicated the sliding of the spores observed at 15 s for Bs PY79 *spsA*.



Fig. 9: Schematic of the suggested cleaning mechanisms using foam flow

**Table 1:** Description of the different flow conditions carried out. The FFC and CIP represent the foam flow conditions and the cleaning-in-place that has been processed for removal of Bacillus spores from stainless steel surfaces, respectively.

Flow	u (cm.s <sup>-1</sup> )	Ca [10 <sup>-3</sup> ]	$\tau$ (Pa)	Dmean	F005	F02	RFS	f (II )
conditions	1.5	~ ~		(mm)	0.000	0.00	(%)	(HZ)
FFC	1.5	0.5	2.2	0.34	0.002	0.29	957	20
	3	1.1	4.2	0.30	0.006	0.36	688	39
	4.5	1.6	5.9	0.28	0.013	0.44	491	45
	9	3.3	9.8	0.18	0.022	0.68	130	70
	13.5	4.8	13.2	0.21	0.043	0.60	193	76
CIP	180		10					

uping for PY79 and *spsA* spores. According to the Tukey grouping, letters were indicated with potentially three classes A, AB, B, C and D depending on the parameter: common letters meaning no significant differences.

Mean Velocity of FFC conditions	f		Kmax1		Kmax2	
	PY79	spsA	PY79	spsA	PY79	spsA
1.5 cm/s	С	С	А	А	В	В
4.5 cm/s	В	D	А	AB	В	В
9 cm/s	А	В	А	AB	А	А
13.5 cm/s	А	А	А	В	В	В