

Temperature and frequency dependence on dielectric permittivity and electrical conductivity of recycled Liquid Crystals

Ana Barrera, Corinne Binet, Frédéric Dubois, Pierre-Alexandre Hébert, Philippe Supiot, Corinne Foissac, Ulrich Maschke

► To cite this version:

Ana Barrera, Corinne Binet, Frédéric Dubois, Pierre-Alexandre Hébert, Philippe Supiot, et al.. Temperature and frequency dependence on dielectric permittivity and electrical conductivity of recycled Liquid Crystals. Journal of Molecular Liquids, 2023, Journal of Molecular Liquids, 378, pp.121572. 10.1016/j.molliq.2023.121572 . hal-04310275

HAL Id: hal-04310275 https://hal.univ-lille.fr/hal-04310275

Submitted on 31 Jan2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Temperature and frequency dependence on dielectric permittivity and electrical conductivity of recycled Liquid Crystals	1 2
Ana Barrera ^a , Corinne Binet ^a , Frédéric Dubois ^b , Pierre-Alexandre Hébert ^c , Philippe Supiot ^a , Corinne Foissac ^a and Ulrich Maschke ^{a,*}	3 4
^a Unité Matériaux et Transformations (UMET), UMR 8207, CNRS, INRAE, Centrale	5
Lille, Université de Lille, 59000 Lille, France	6
^b Unité de Dynamique et Structure des Matériaux Moléculaires (UDSMM), UR 4476,	7
Université du Littoral Côte d'Opale, 62100 Calais, France	8
^c Laboratoire d'Informatique Signal et Image de la Côte d'Opale (LISIC), UR 4491,	9
Université du Littoral Côte d'Opale, 62100 Calais, France	10
*Corresponding author: E-mail address: <u>ulrich.maschke@univ-lille.fr</u>	11 12

ABSTRACT

The dielectric behavior and the electrical conductivity response of recycled nematic 14 Liquid Crystals (LCs) were studied over a wide range of frequencies (0.1 to 10⁶ Hz) 15 and temperatures (100 to -20 °C) with a broadband dielectric spectroscopy technique. 16 These LC mixtures were extracted from a total of 65700 End-Of-Life (EOL)-Liquid 17 Crystals Displays (LCD) such as computers, tablets and TVs of various sizes, brands, 18 and production years. After a stage of purification, the recycled LC blends seemed to 19 possess similar physical properties as conventional nematic LCs. In this report, two 20 types of anchoring of LC molecules, obtained with homeotropic or homogeneous 21 alignment cells, were investigated. All dielectric data were analyzed with adapted 22 23 models: Havriliak-Negami and Almond West formalisms for the frequency dependence of the data, as well as Arrhenius and Vogel-Fulcher-Tammann models for their 24 temperature dependence. In this article, it is demonstrated that despite the wide variety 25 of EOL-LCD devices used leading to a large number of LC molecules, the frequency 26 and temperature evolution of the recorded spectra are comparable to those of 27 commercial LC mixtures as reported in the literature. Therefore, their future reuse 28 29 could be envisaged for display applications.

Keywords: nematic liquid crystals; recycling; dielectric properties; modelization. 30

31

1. Introduction

The use of Electrical and Electronic Equipment (EEE) is increasingly becoming an 33 indispensable part of modern life [1,2]. A large part of the world population benefits 34 from a higher standard of living, because of their availability, widespread use and 35 simple access [3]. Nevertheless, the way we produce, consume and manage Waste 36 Electrical and Electronic Equipment (WEEE) is not long-term sustainable [4]. The 37 European Directive WEEE (2012/19/EU) has the main objective to reduce the 38 production of WEEE and requires, especially, the disassembly of all Liquid Crystals 39 Displays (LCDs) with an area over 100 cm² [5]. These devices may contain hazardous 40 chemicals, which are registered in the European RoHS Directive (2011/65/EU) 41 42 (Restriction on the use of certain Hazardous Substances) [6].

Since the last half-century, there is grown interest in technologies concerning LCs in 43 all domains of LCD devices. A presence of an average of 5.6 LCD panels per French 44 household was documented in 2018, which is constantly increasing [7]. This growing 45 demand for screens is generating more and more waste as obsolete screens are 46 replaced by newer ones. Excluding premature device replacement, the lifetime of a 47 LCD is approximatively ten years [8]. Televisions, computers and tablets based on LC 48 49 technologies includes a set of high value-added materials whose recycling represents a strategic issue [9]. 50

Effective recovery of LCD panels is an expensive process that requires sophisticated 51 technologies but revalorization of LCs, one of the main components of LCD panels, 52 could be relevant and profitable. To reuse these End-Of-Life (EOL)-LC mixtures in new 53 devices, they should present similar chemical and physical properties than the original 54 commercial LCs mixtures. For this purpose, purification processes must be established 55 in order to eliminate impurity traces. Indeed, the presence of inorganic impurities in LC 56 mixtures can induce various undesirable effects; for example, they tend to increase the 57 electrical conductivity, which may lead to a loss of electro-optical or optical properties 58 59 (screen flickering, persistence, long response time...) [10–14].

32

In the context of this report, the ability of reusing EOL-LC mixtures in new devices will 60 be investigated in terms of their chemical and physical properties. Generally, nematic 61 LCs present two modes of molecular relaxations, which can be observed using 62 dielectric spectroscopy. The first relaxation mode occurs in the frequency range 63 between 10 and 10⁶ Hz and is related to the reorientation of the molecular dipoles 64 around the short axis of the LC-molecule. The second relaxation mode, above 1 MHz, 65 is attributed to the reorientation around its long axis [15-17]. In this report, dielectric 66 permittivities and electrical conductivities of recycled LCs have been evaluated and 67 discussed over wide frequency and temperature ranges. 68

2. Experimental and methods

2.1. Materials

71 The EOL-LCs studied in this report originate from an electronic waste company located in Northern France named ENVIE²E. LC mixtures were extracted from a 72 heterogeneous deposit of EOL-LCDs including computers, tablets and TVs of different 73 sizes, types, brands and years of production. The average diagonal size of the 74 recovered EOL-LCD was 24 inch. First, EOL-LCDs were manually sorted and 75 dismantled in order to separate different components: electronic boards, batteries, 76 scrap metals, speakers, capacitors, plastic foils, lamps, LCD panels, plastics and 77 78 plasma glass.

2.2. Separation and Characterization of components

Subsequently, 65700 LCD panels were opened and LCs molecules were extracted 80 using an organic solvent bath. Additional information about the applied LCs extraction 81 process have been developed in a patent filed by U. Maschke *et al.* (2015) [18]. The 82 native extracted solutions consisted mainly of LCs molecules, organic solvent and both 83 organic and inorganic impurities. These solutions have an unusual black coloration that 84

4

69

70

79

-

is not representative of nematic LC molecules meaning that a purification process must 85 be performed. 86

For this purpose, the extracted EOL-LCs were purified to remove undesirable 87 impurities, using several distillation and chromatographic techniques. Then, the 88 purified LC samples were characterized to determine their optical, thermal and 89 dielectric properties. 90

After purification, all investigated EOL-LCs exhibited quite comparable dielectric91properties at room temperature [19]. Therefore, only one of the obtained LC mixtures92named "recycled LC mixture" will be studied as a function of temperature and93frequency in the following sections in order to avoid redundancy.94

The texture of this LC sample was observed under a Polarizing Optical Microscope 95 (POM) (Olympus BX60), coupled to a Linkam LTS 350 heating/cooling stage with a 96 TMS 94 temperature control unit. POM analysis were carried out in a descending 97 temperature range from 90 to 20 °C. Between 90 and 73 °C, the recycled LC mixture 98 remained in the isotropic phase where the LCs molecules did not show any order of 99 orientation or position (Fig. 1). The beginning of the isotropic-nematic transition was 100 observed around 73 °C, and thus, the ordered nematic state appears when the 101 temperature decreases and the sample becomes optically opaque. At room 102 temperature, the LC mixture presented a morphology similar to that of nematic LCs 103 (Schlieren texture). To complement the POM study, Differential Scanning Calorimetry 104 105 (DSC) analysis were performed using a TA Instruments DSC 10 calorimeter at heating and cooling rates of 10°C/min, under continuous nitrogen flow. In the obtained 106 thermogram (Fig. 2), a single glass transition (T_g) was detected around -90 °C. At about 107 108 70 °C, an endothermic peak corresponding to the nematic-isotropic transition (T_{NI}) was also identified, and this temperature was in good agreement with that found in the 109 optical analysis. The LC mixture presents thus a large range of temperature of the 110 nematic phase between T_q and T_{NI} . 111

112



Fig. 1. Texture of the recycled LC mixture observed under Polarizing Optical Microscope 114 (POM) between crossed polarizers in a descending temperature range from 90 to 20 °C. The scale bar represents 100 µm.



Fig. 2. DSC thermogram of the recycled LC mixture obtained during heating and cooling 119 modes in the temperature range between -150 and 100 °C. 120

2.2. Dielectric measurements

The Broadband Dielectric Spectroscopy (BDS) measurements were carried out using 122 a ModuLab-MTS impedance analyzer (Solartron Analytical, Ametek) in the frequency 123 range from 0.1 to 10⁶ Hz. The amplitude of the applied sine wave voltage was 1 V. The 124 LC samples were inserted into commercial homogeneous (HG) or homeotropic (HT) 125 alignment cells of 20 µm thickness by capillary action at a temperature slightly higher 126 than nematic-isotropic transition (73 °C). The commercial cells used in this study were 127 supplied by AWAT company (Warsaw, Poland). The active surface was 0.25 cm² and 128 the sheet resistance of the ITO coating electrode was about 10 Ω/\Box . The homogeneous 129 and homeotropic alignments of the cells employed in this study were confirmed using 130

113

115 116

117

118

121

POM, as previously reported by Barrera *et al.* 2021 [19]. The temperature of the131samples was controlled with a Linkam LTS350 stage. The temperature range chosen132was from 100 to -20 °C.133

BDS allows to determinate the complex relative dielectric permittivity ε^* , which can be 134 related to the complex conductivity σ^* by the following relation (Eq. (1)) [20]: 135

$$\sigma^*(\omega) = i\omega\varepsilon_0\varepsilon^* \tag{1} 136$$

where ω stands for the angular frequency ($\omega = 2\pi f$, *f* is the frequency) of the measuring 137 electric field, and ε_0 represents the dielectric permittivity of vacuum (8.85418 ×10⁻¹² 138 F.m⁻¹).

The complex relative dielectric permittivity and complex conductivity can be written140according to Eqs. (2) and (3):141

$$\varepsilon^*(\omega) = \varepsilon'(\omega) - i\varepsilon''(\omega) \tag{2}$$

$$\sigma^*(\omega) = \sigma'(\omega) + i\sigma''(\omega) \tag{3}$$

where ε' and ε'' describe respectively the real part (related to the stored energy) and 144 the imaginary part (related to the dissipated energy) of ε^* ; σ' and σ'' represent 145 respectively the real and imaginary parts of σ^* . 146

To obtain σ and σ from dielectric parameters, Eqs. (4) and (5) are used: 147

$$\sigma'(\omega) = \omega \varepsilon_0 \varepsilon''(\omega) \tag{4}$$

$$\sigma''(\omega) = \omega \varepsilon_0 \varepsilon'(\omega) \tag{5}$$
 149

150

The Almond-West formalism is widely applied to analyze the dependence with151frequency of the real part of the complex conductivity. This model is derived from152Jonscher's universal power law [21] and can be expressed as follows:153

$$\sigma'(\omega) = \sigma_{DC} \left(1 + \left(\frac{f}{f_c}\right)^n \right) \tag{6}$$
 154

where σ_{DC} , f_c and n represent the DC conductivity, the characteristic frequency and 155 the degree of interaction between the mobile ions and their environment, respectively. 156 There are many models allowing to fit the dependence with frequency of the complex 157 conductivity such as the Cole-Cole, Cole-Davidson, Havriliak-Negami (H-N) or Kohl- 158 Rausch-Williams-Watts equations, previously transformed into their conductivity 159 160 representation [22-24]. In order to reuse these EOL-LCs in LCD systems, it is also essential to know their 161 dielectric anisotropy ε_a . This parameter characterizes the ability of LC molecules to 162 orient themselves in response to an external electric field. ε_a is obtained from the 163 difference between the dielectric permittivities when the electric field is parallel to the 164 molecular director of the LC called "n" [25] (${\varepsilon^\prime}_{//}$ obtained with HT cells) and when the 165 field is perpendicular to "n" (${\cal E}_{\perp}$ obtained with HG cells). Therefore, dielectric 166 167 anisotropy can be expressed by Eq. (7):

$$\varepsilon_a = \varepsilon'_{//} - \varepsilon'_{\perp} \tag{7} \quad 168$$

3. Results and discussion

3.1 Temperature dependence of dielectric permittivity and electrical conductivity of the170recycled LC mixture171

3.1.1 Dielectric permittivity spectra

The evolution with temperature and frequency of the real and imaginary parts of ε^* of 173 the recycled LC mixture in HT and HG alignments is presented in Fig. 3 and 4. The 174 temperature range chosen is comprised between 100 °C and -20 °C, and the results 175 are presented in cooling mode. These temperatures were selected in relationship with 176 those used for industrial LCD applications. 177

In HT alignment, the dielectric spectra $\varepsilon'(f)$ depend strongly on both temperature and 178 frequency (Fig. 3 a). At high temperatures and a frequency of 0.1 Hz, ε' reaches its 179 highest value of 650 whereas at low temperatures ε' decreases to around 10. The 180 significant increase of ε' with decreasing frequency is due to a phenomenon of 181 polarization of charges and the contribution of the electrical conductivity [20,26]. When 182 the temperature decreases, both phenomena are shifted towards low frequencies. At 183 intermediate and high frequencies, ε' show**s** a frequency independent plateau for 184

169

temperatures from 100 to 75 °C. Below 70 °C in the frequency region about 10⁵ Hz, ε' 185 starts to decrease slightly due to the appearance of relaxation mechanisms. 186

Fig. 3 b) shows the $\varepsilon''(f)$ spectra of the recycled LC mixture in HT alignment. At low 187 frequencies, the increase of ε'' with decreasing frequency can be attributed to 188 phenomena of conduction and electrode polarization. For temperatures below 20 °C 189 and in the frequency range between 10^2 and 10^6 Hz, two overlapping dielectric 190 191 relaxation mechanisms can be observed. With decreasing sample temperature, these mechanisms shift to the intermediate frequency region. The main mechanism, 192 although of lower amplitude, is located between $\sim 10^2$ and 10^3 Hz and probably related 193 to the dielectric LC glass transition (around -90 °C). The secondary mechanism, 194 localized between $\sim 10^4$ and 10^6 Hz, is attributed to the process of rotation/reorientation 195 of LC molecules around their minor axis [17,27,28]. This phenomenon could also be 196 related to the contribution of parasitic impedances caused by connectors, cables, etc., 197 which become significant at frequencies above 100 kHz [19] [26], and the relaxation of 198 199 the measuring cell due to the resistance of the ITO layer [29-31]. However, these parasitic effects and ITO influences seem to be absent here since the empty cell 200 dielectric measurements in both alignments at different temperatures did not reveal 201 such effects in the frequency window of the experiment. 202

Following current literature, the dielectric relaxation mechanisms were analyzed by203applying known models (Debye, Cole-Cole, Cole-Davidson and Havriliak-Negami)204[22–24,32]. The study of these mechanisms will be presented in the following section.205



Fig. 3. Dielectric spectra a) $\varepsilon'(f)$ and b) $\varepsilon''(f)$ of the recycled LC mixture in homeotropic 209 alignment. The spectra were taken in a descending temperature range between 100 and 210 -20 °C. 211

The dielectric spectra $\varepsilon'(f)$ and $\varepsilon''(f)$ of the recycled LC mixture in HG alignment are 212 presented in Fig. 4 a) and 4 b), respectively. The behaviors are very similar to those 213 observed in HT alignment. However, conduction and interfacial polarization 214 phenomena appear at slightly higher frequencies for all considered temperatures. 215 Furthermore, the amplitudes of the relaxations observed at high frequency and low 216 temperature seem to be slightly less intense than in the HT alignment. 217 218 It is interesting to note that the evolutions with frequency and temperature of the spectra observed for the recycled LC mixture are comparable to those of commercial 219 LC mixtures reported in the literature [27,28]. 220

221





Fig. 4. Dielectric spectra a) $\varepsilon'(f)$ and b) $\varepsilon''(f)$ of the recycled LC mixture in homogeneous alignment. The spectra were taken in a descending temperature range between 100 and -20 °C.

3.2 Dielectric relaxation mechanisms

3.2.1 Havriliak-Negami model

The dielectric relaxation mechanisms were studied applying the H-N formalism. In 231 order to perform an analysis of the experimental data in the whole frequency range, a 232 supplementary term was added to the original formalism. The modified H-N model is 233 illustrated in Eq. (8): 234

$$\varepsilon^*(f) = \varepsilon_{\infty} + \sum_k \frac{\Delta \varepsilon_k}{\left[1 + \left(i\frac{f}{f_{ck}}\right)^{\alpha_k}\right]^{\beta_k}} + \frac{\sigma}{\varepsilon_0(i2\pi f)^n}$$
(8) 235

where *k* indexes the relaxation mechanisms, $\Delta \varepsilon$ is the strength of the relaxation 236 mechanism and represents the decrease of the real part of the complex dielectric 237 permittivity ($\Delta \varepsilon = \varepsilon_s - \varepsilon_{\infty}$ with ε_s and ε_{∞} being the low (static) and high frequency 238

permittivities, respectively). f_c corresponds to the characteristic frequency also called 239 relaxation frequency, and α and β are the shape parameters describing the width and 240 asymmetry of the dielectric spectra, respectively ($0 < \alpha, \beta \le 1$). The added term 241 corresponds to the contribution of the electrical conductivity; ε_0 represents the vacuum 242 permittivity, σ is associated with the DC conductivity, and *n* stands for the degree of 243 interaction between the mobile ions and their environment. 244

The spectra $\varepsilon'(f)$ and $\varepsilon''(f)$ were simultaneously fitted with Eq. (8) using the Levenberg-245 246 Marquardt algorithm, applying Grafity software. Fig. 5 a) shows the fit of the dielectric spectra $\varepsilon'(f)$ and $\varepsilon''(f)$ of the recycled LC mixture at -20 °C in HT alignment. The 247 experimental data and the fit are represented by symbols and red solid lines, 248 respectively. Fig. 5 b) shows the decomposition of the fit of the $\varepsilon''(f)$ spectrum into three 249 processes: the contribution of the electrical conductivity (black line at low frequencies) 250 and two H-N relaxations (the main relaxation in blue color at intermediate frequencies 251 and the secondary one in green color at high frequencies). The inset of this figure 252 shows the secondary mechanism on a linear y-scale to highlight the quality of the fit. 253





Fig. 5. a) Simultaneous fit to the H-N model of the dielectric spectra $\varepsilon'(f)$ and $\varepsilon''(f)$ of the 258 recycled LC mixture at -20 °C in HT alignment; the fit is represented by the red solid line. b) 259 Decomposition of the dielectric spectrum $\varepsilon''(f)$ into three processes: two relaxation 260 mechanisms and electrical conductivity. The inset represents the secondary mechanism using 261 a linear y-scale. c) Example of dielectric spectra $\varepsilon''(f)$ and M''(f) at -15 °C in HT alignment. 262

Table 1 reports the fit parameters obtained using the H-N model for the recycled LC264mixture in HT alignment from -20 to 15 °C and in HG alignment from -20 to 0 °C.265Outside these temperature ranges, the dielectric strength of the main mechanism is266insufficient within the studied frequency window, and the characteristic frequency of267the secondary mechanism is too high to allow fitting.268

The main relaxation mechanism of the LC mixture in both alignments exhibits Cole-269 270 Cole behavior at -20 °C and follows the Debye model for all other temperatures. The Cole-Cole equation can also be used to fit the secondary relaxation mechanism, 271 except when the LC is in HG alignment at 0 °C; in this case the H-N formalism is 272 required. As the temperature decreases, and thus approaches T_g, the characteristic 273 frequency of the mechanisms decreases and their dielectric strength increases. The 274 evolution with temperature of the characteristic frequency is detailed in the next 275 section. 276

HT cell											
Main mechanism					Secondary mechanism						
T (°C)	σ (nS.m ⁻¹)	n	f_{c^1} (kHz)	$\Delta arepsilon_1$	α	eta	<i>f</i> _c ² (kHz)	Δε2	α	β	R ²
15	3.82	0.98	39.81	0.15	1	1	414.3	4.16	0.794	1	0.995
10	2.92	0.98	23.12	0.31	1	1	269.1	4.96	0.819	1	0.995
5	2.37	1.00	14.04	0.16	1	1	174.3	5.41	0.776	1	0.996
0	1.86	1.00	7.786	0.13	1	1	108.0	5.45	0.762	1	0.997
-5	1.06	0.97	5.002	0.53	1	1	69.46	4.95	0.846	1	0.999
-10	0.739	0.97	3.135	0.55	1	1	42.88	5.05	0.842	1	0.998
-15	0.498	0.97	1.460	0.55	1	1	24.32	5.19	0.841	1	0.999
-20	0.301	0.96	0.931	0.74	0.949	1	14.50	5.32	0.867	1	0.998
					HG c	ell					
0	0.863	1.00					288.4	1.21	0.629	1	0.987
-5	0.762	0.99	3.990	0.03	1	1	92.63	1.42	0.667	1	0.989
-10	0.474	0.98	2.769	0.11	1	1	43.57	1.72	0.749	1	0.998
-15	0.411	0.99	1.441	0.17	1	1	22.55	2.44	0.773	1	0.999
-20	0.318	1.00	1.065	0.19	0.966	1	13.40	2.55	0.737	1	0.999
Uncertainty (%)	1-5	0.5-4	4-8	3-10	2-5	2-5	5-8	2-10	1-4	2-5	

Table 1. H-N model parameters for fitting dielectric spectra $\varepsilon'(f)$ and $\varepsilon''(f)$ of the recycled LC mixture in HT and HG alignments.

To complement the obtained results from the dielectric formalism, the complex 279 dielectric modulus ($M^* = 1/\varepsilon^*$) has been employed. This representation could provide 280 a more complete picture of the system and the imaginary part (M'') of M^* offers 281 information in terms of conductivity. An example of M'' as a function of frequency, 282 taken at -15 °C, is shown in Fig. 5 c). Complementary results are presented in 283 supplementary part (Fig. S1). Fig. 5 c) exhibits a Debye peak appearing at low 284 frequencies associated to the conductivity. Two other relaxations peaks, linked to main 285 and secondary modes, were also found. Analysis of all data allowed to follow the 286 frequency evolution of the peak maxima (f_{max}) as a function of temperature (see next 287 section). 288

289

290

3.2.2 Dependence of the characteristic frequency on temperature

The Arrhenius (Eq. (9)) and the Vogel-Fulcher-Tammann laws (VFT) (Eq. (10)) are 291 commonly used to analyze the Arrhenius plot of the characteristic frequency of 292 relaxation mechanisms.

$$f_c = f_{c0} \exp\left(-\frac{E_a}{kT}\right) \tag{9} 294$$

where f_{c0} is the pre-exponential factor, E_a represents the activation energy, k stands 295 for the Boltzmann constant (1.38 × 10⁻²³ J·K⁻¹) and T is the temperature in Kelvin. 296

$$f_c = f_{c0} \exp\left(-\frac{B}{T - T_0}\right) \tag{10} 297$$

where *B* represents the "pseudo activation energy" $\left(B = \frac{E_a}{k}\right)$ in Kelvin and T_0 is known 298 as the Vogel temperature. T_0 is always 30 to 70 °C lower than T_g. 299

The Arrhenius plot of the characteristic frequency of the main and secondary relaxation300mechanisms of the recycled LC mixture in HT alignment is shown in Fig. 6 as well as301obtained results from M'' analysis. It can be seen that both mechanisms exhibit302Arrhenius-like behavior. The activation energies obtained by fitting the experimental303data (Eq. (9)) are gathered in Table 2.304



305

Fig. 6. Arrhenius plot of characteristic frequencies and frequency evolution of the peak maxima306of the mechanisms for recycled LC mixture in HT alignment obtained from ε^* and M''. The307experimental data and the fits to the Arrhenius model are represented by symbols and red308solid lines, respectively. Results from HG cell are presented only for conductivity data from309M''.310

Table 2. Arrhenius activation energy of the relaxation mechanisms of the recycled LC mixture	311
in HT and HG alignments (fit for HG alignment is given in supplementary material, Fig. S2).	312

		Main mechanism	Secondary mechanism
		E_a (eV)	E_a (eV)
HT cell	<i>€</i> ′′	0.684 ± 0.003	0.602 ± 0.003
	Μ''	0.734 ± 0.002	0.645 ± 0.001
HG cell	<i>€</i> ′′	0.541 ± 0.050	0.741 ± 0.052
	Μ''	0.794 ± 0.054	0.657 ± 0.064

The main mechanism, previously associated with T_g , was expected to exhibit VFT 313 behavior. Nevertheless, the experimental data follow the Arrhenius model; this may be 314 related to the fact that the narrow temperature range investigated for the dielectric 315

study is far from T_g ~ -90 °C. The secondary mechanism also follows an Arrhenius-like 316 behavior, as reported by other authors; it is associated to the molecular rotation around 317 the minor axis [17]. The activation energies obtained are consistent with the values 318 found in the literature for commercial nematic LC mixtures [17,27,33]. Fig. 6 also 319 exhibits frequency data from the peak maxima of M''. These values, for the main and 320 secondary mechanisms, were found to be higher than the characteristic frequencies 321 from ε^* . As expected, almost parallel lines have been observed in Fig. 6 as function of 322 the reduced temperature (1000 \cdot T⁻¹). Indeed, f_c and f_{max} are proportional and 323 associated by $f_{max} = \left(\frac{\varepsilon_s}{\varepsilon_m}\right) fc$ [26]. It should be mentioned that the latter equation is 324 only valid in the case of Debye-like relaxation. The obtained activation energies (Table 325 2) are comparable since the same relaxation mechanisms are involved. Considering 326 the results of the maxima deduced from conductivity peak of M'' in both alignments, a 327 VFT-like behavior was observed. The corresponding T₀ values were 129 and 154 K for 328 HG and HT alignments, respectively, thus in good agreement with the Tg (183 K) value 329 determined by DSC analysis. 330

3.3 Dielectric anisotropy

331

Fig. 7 shows the evolution with the inverse of the temperature of the real part of the 332 complex dielectric permittivity of the recycled LC mixture, and of its dielectric 333 anisotropy ε_a determined at 1 kHz. At high temperatures, when the LC mixture is in the 334 isotropic phase, the perpendicularly and parallel permittivities are very close, 335 consequently ε_a is very small. As the temperature decreases, in the nematic phase, 336 the permittivity of the LC mixture in HG alignment changes little ($\varepsilon_{\perp}^{\prime}\sim$ 4.0), whereas 337 that in HT alignment increases steadily up to T= -5°C. The latter behavior is due to the 338 decrease in dynamic/thermal movements of the molecules as the temperature 339 decreases [34]. From -5 °C to -20°C, the permittivity of the LC mixture in HT alignment 340 varies little ($\varepsilon'_{//} \sim 8.2$). Thus, the dielectric anisotropy shows a similar evolution as that 341

20

of the longitudinal component $\varepsilon'_{//}$ of the permittivity, and reaches a value of 4.0 at T= 342 -5°C. 343



Fig. 7. Evolution with the inverse temperature of the real part of the complex dielectric345permittivity of the recycled LC mixture, and of the dielectric anisotropy (temperature (°C) as346secondary X axis).347

3.4 Electrical conductivity spectra

Figures 8 a) and b) show the dependence with temperature of the real part of the 349 complex conductivity of the recycled LC mixture in HT and HG alignments respectively, 350 351 and the fits of these spectra to the Almond-West model (Eq. (6)). For reasons of clarity, only a few temperatures were illustrated. Nevertheless, spectra over the entire 352 temperature range were presented in the supplementary material (Fig. S3). When the 353 354 recycled LC mixture is in HT alignment, two behaviors can be distinguished (Fig. 8 a): - At low frequencies, a plateau corresponding to the DC conductivity is observed. 355 Conductivity values increase with temperature. This behavior suggests that electrical 356 conductivity in LCs represents a thermally activated process [26]. At very low 357 frequencies (<1 Hz), and temperatures between 100 and 20 °C, the σ ' values of the 358 359 mixture decrease as the frequency decreases, due to the electrode polarization. This effect will not be taken into account in this report. 360

348

- At higher frequencies, conductivity depends on frequency and increases drastically 361 with it. This is due to the ionic relaxation of mobile charge carriers. This effect is known 362 as subdiffusive conductivity. Between 10 kHz and 1 MHz, a weak decrease in σ 363 appears from 20 to -20 °C due to the relaxation mechanisms already observed in the 364 dielectric spectra (Fig. 3 a and b). These observations are also valid when the LC 365 alignment is homogeneous (Fig. 8 b). However, in this configuration, at low 366 367 frequencies, the phenomenon of polarization of the electrode seems more important and visible. 368

The evolution with temperature (100 to -20°C) of the DC conductivity of the recycled369LC mixture in HT and HG alignments is analyzed in the following section using the370Almond-West formalism.371





Fig. 8. Conductivity spectra $\sigma'(f)$ of the recycled LC mixture in a) HT and b) HG alignments375from 100 to -20 °C. The experimental data and the fits to the Almond-West model are376represented by symbols and red solid lines, respectively.377

378

379

3.5 DC conductivity

The Arrhenius plot of the DC conductivity obtained with Eq. (6) are analyzed with the380Arrhenius and VFT models given by Eqs. (11) and (12), respectively [17,26].381

$$\sigma_{DC} = \sigma_{DC0} \exp\left(-\frac{E_a}{kT}\right) \tag{11} 382$$

$$\sigma_{DC} = \sigma_{DC0} \exp\left(-\frac{B}{T-T_0}\right) \tag{12}$$

where σ_{DC0} is a pre-exponential factor. The other parameters have already been 384 presented in section 3.2.2. The Arrhenius plot of the DC conductivity for the two LC 385 alignments and the fits to the VFT and Arrhenius models in the nematic phase are 386 shown in Fig. 9 for the recycled LC mixture. 387

The plot in HT alignment follows the VFT model. In contrast, the plot in HG alignment 388 cannot be correctly fitted using either Arrhenius or VFT models over the entire 389 temperature range (supplementary material, Fig. S4). In particular, the use of the VFT 390

model leads to an uncertainty on the results of more than 20% (Table 3). In fact, the391combination of the Arrhenius formalism from 80 to 10 °C and the VFT model from 10392to -20 °C allows a better fit of the activation plot of the DC conductivity of the recycled393LC mixture (Fig. 9, Table 3). VFT-like behavior was expected for the recycled LC394mixture in both alignments. If we exclude the approach, the device and the types of395cells used, it is obvious that the mixtures of purified LCs are responsible for this396behavior. Several hypotheses can explain this behavior:397

- The nature and quantity of the LC molecules and ionic species of the recycled mixture 398 are not exactly known. 399

Chemical, thermal and optical characterizations have proved that these mixtures 400 present macroscopically an eutectic behavior with very similar characteristics. 401 However, their behavior at the microscopic scale is not known. 402

Some LC molecules and ions present in the mixture may act preferentially on the HG
 or HT alignment of the cell. In addition, the proportion of molecules/ions that interact
 404
 may be different for each alignment.
 405

The Arrhenius activation energy of DC conductivity of the recycled LC mixture from 80 406 to 10 °C is equal to 0.341 eV, which is in good agreement with the data from literature 407 for commercial nematic LCs [17,27,33]. Moreover, the value of the Vogel temperature 408 of this mixture ($T_g = 166$ K) is consistent with the glass transition temperature 409 determined by DSC method ($T_g = 183$ K), as already mentioned before (cf. data from 410 M''). The DC conductivity values obtained with the Almond-West formalism and by Eq. 411 (8) are in complete agreement. 412



Fig. 9. Arrhenius plot of the DC conductivity of the recycled LC mixture in HT and HG414alignments (temperature in °C as secondary X axis). The experimental data and the fits to VFT415and Arrhenius models in nematic phase are represented by symbols and solid lines,416respectively.417

418

413

Table 3. Arrhenius and VFT models parameters for fitting Arrhenius plots of the DC419conductivity of the recycled LC mixture in HT and HG alignments (Conductivity spectra).420

					421		
HG cell							
E_a (eV)	$\sigma_{\! DC0}$ Arr (S·m ⁻¹)	<i>B</i> (K)	<i>T</i> ₀ (K)	$\sigma_{\! DC0}$ vft (S·m ⁻¹)	_		
0.341 ± 0.001	2.7E-03 ± 1.1E-05						
		2189 ± 489	59 ± 17	6.9E-05 ± 8.2E-06			
0.336 ± 0.001	2.1E-03 ± 1.2E-05	642 ± 7	166 ± 5	3.5E-07 ± 1.2E-09			
HT cell							
		779 ± 9	165 ± 3	2.3E-06 ± 2.1E-07	_		
	$\frac{E_a \text{ (eV)}}{0.341 \pm 0.001}$ 0.336 ± 0.001	HG E_a (eV) σ_{DC0} Arr (S·m ⁻¹) 0.341 ± 0.001 2.7E-03 ± 1.1E-05 0.336 ± 0.001 2.1E-03 ± 1.2E-05 HT	HG cell E_a (eV) σ_{DC0} Arr (S·m ⁻¹) B (K) 0.341 ± 0.001 2.7E-03 ± 1.1E-05 2189 ± 489 0.336 ± 0.001 2.1E-03 ± 1.2E-05 642 ± 7 HT cell 779 ± 9	HG cell E_a (eV) σ_{DC0} Arr (S·m ⁻¹) B (K) T_0 (K) 0.341 ± 0.001 2.7E-03 ± 1.1E-05 2189 ± 489 59 ± 17 0.336 ± 0.001 2.1E-03 ± 1.2E-05 642 ± 7 166 ± 5 HT cell 779 ± 9	HG cell E_a (eV) $\sigma_{DC0 \text{ Arr}}$ (S·m ⁻¹) B (K) T_0 (K) $\sigma_{DC0 \text{ VFT}}$ (S·m ⁻¹) 0.341 ± 0.001 2.7E-03 ± 1.1E-05 2189 ± 489 59 ± 17 6.9E-05 ± 8.2E-06 0.336 ± 0.001 2.1E-03 ± 1.2E-05 642 ± 7 166 ± 5 3.5E-07 ± 1.2E-09 HT cell		

422

423

4. Conclusions

In this report, the dielectric properties of a recycled LC mixture were studied using an 424 impedance analyzer in the frequency range between 0.1 and 10⁶ Hz, and in both 425 homogeneous and homeotropic alignments. Measurements were limited to the 426

temperature range from 100 to -20 °C in which an LCD device is typically used 427 428 according to the technical specifications. Furthermore, this range covered the nematic and isotropic phases of the investigated EOL-LCs. It has been illustrated that the 429 frequency and temperature evolutions of the obtained spectra are conveniently 430 comparable with those of conventional LC mixtures reported in literature. Two 431 overlapping dielectric relaxation mechanisms were observed at temperatures below 432 20 °C and in the frequency range located between 10² and 10⁶ Hz for both alignments. 433 These phenomena, named main and secondary mechanisms, were attributed to the 434 dielectric LC glass transition ($T_q = -90$ °C, determined by DSC analysis), and the 435 rotation of LC molecules around the short axis, respectively, and they follow an 436 437 Arrhenius-like behavior. In contrast, the Arrhenius plot of the DC conductivity do not follow the same model in both alignments. These results revealed that the charge 438 transport phenomena in both homogeneous and homeotropic alignments are not 439 governed by the same mechanisms. Finally, the Arrhenius activation energies of the 440 relaxation mechanisms as well as the DC conductivity of the recycled LC mixture are 441 in good agreement with those of commercial nematic LCs as documented in literature. 442 Regardless of the diverse EOL-LCDs devices collected at the company, it was possible 443 to obtain recycled LC mixtures with thermal, optical, dielectric and electric properties 444 equivalent to those of conventional LCs mixtures. Therefore, their future reuse could 445 be envisaged for display applications. 446

CRediT authorship contribution statement

Ana Barrera: Conceptualization, Methodology, Investigation, Data curation, Writing—
original draft preparation, Ulrich Maschke: Conceptualization, Writing—original draft
preparation, Corinne Binet: Methodology, Investigation, Data curation, Corinne
Foissac: Writing—review and editing, Philippe Supiot: Writing—review and editing,
Frédéric Dubois: Methodology, software and data validation, Pierre-Alexandre
Hébert: Software and data validation. All authors have read and agreed to the
published version of the manuscript.

Decl	aration of Competing Interest	455
The	authors declare that they have no known competing financial interests or personal	456
relati	ionships that could have appeared to influence the work reported in this paper.	457
		458
Ackr	nowledgments	459
The	authors acknowledge financial support from the University of Lille, the "Région	460
Haut	s-de-France (FEDER)", I-Site of Lille and ENVIE ² E du Nord.	461
		462
Fund	ding	463
This	work was supported by Région Hauts-de-France (FEDER), I-Site Lille and	464
ENV	IE ² E du Nord.	465
		466
Арр	endix A. Supplementary material	467
The	following are the Supplementary data to this article:	468
		469
Refe	rences	470
[1]	V. Goodship, A. Stevels, J. Huisman, Waste electrical and electronic equipment (WEEE) Handbook, 2 nd ed., Woodhead Publishing (Elsevier), Cambridge, UK, 2019. https://doi.org/10.1016/C2016-0-03853-6.	471 472 473
[2]	O.S. Shittu, I.D. Williams, P.J. Shaw, Global E-waste management: Can WEEE make a difference? A review of e-waste trends, legislation, contemporary issues and future challenges, Waste Manag. 120 (2021) 549–563. https://doi.org/10.1016/j.wasman.2020.10.016.	474 475 476 477
[3]	C.P. Baldé, E. D'Angelo, V. Luda, O. Deubzer, R. Kuehr, Global Transboundary E-waste Flows Monitor 2022, United Nations Institute for Training and Research, Bonn, Germany, 2022.	478 479 480
[4]	V. Forti, C.P. Baldé, R. Kuehr, G. Bel, The Global E-waste Monitor 2020: Quantities, flows and the circular economy potential., United Nations University (UNU)/United Nations Institute for Training and Research (UNITAR), Bonn- Geneva-Rotterdam, 2020.	481 482 483 484
[5]	European Union, WEEE Directive 2012/19/EU, (2012). https://eur-	485

	lex.europa.eu/legal- content/EN/TXT/?uri=uriserv%3AOJ.L2012.197.01.0038.01.ENG&toc=OJ%3 AL%3A2012%3A197%3ATOC (accessed November 17, 2022).	486 487 488
[6]	European Union, Restrictions on the use of certain hazardous substances in electrical and electronic equipment, (2017). https://eur-lex.europa.eu/summary/EN/uriserv%3A2004_4 (accessed November 17, 2022).	489 490 491
[7]	Conseil supérieur de l'audiovisuel, L'équipement audiovisuel des foyers au 1er semestre 2018, (2018). https://www.csa.fr/Informer/Collections-du- CSA/Panorama-Toutes-les-etudes-liees-a-l-ecosysteme-audiovisuel/Les- observatoires-de-l-equipement-audiovisuel/L-equipement-audiovisuel-des- foyers-au-1er-semestre-2018 (accessed November 17, 2022).	492 493 494 495 496
[8]	A. Amato, F. Beolchini, End of life liquid crystal displays recycling: A patent review, J. Environ. Manage. 225 (2018) 1–9. https://doi.org/10.1016/j.jenvman.2018.07.035.	497 498 499
[9]	A. Akcil, I. Agcasulu, B. Swain, Valorization of waste LCD and recovery of critical raw material for circular economy: A review, Resour. Conserv. Recycl. 149 (2019) 622–637. https://doi.org/10.1016/j.resconrec.2019.06.031.	500 501 502
[10]	R. Basu, Soft memory in a ferroelectric nanoparticle-doped liquid crystal, Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys. 89 (2014) 1–5. https://doi.org/10.1103/PhysRevE.89.022508.	503 504 505
[11]	Y. Garbovskiy, I. Glushchenko, Nano-objects and ions in liquid crystals: Ion trapping effect and related phenomena, Crystals. 5 (2015) 501–533. https://doi.org/10.3390/cryst5040501.	506 507 508
[12]	A. Kumar, D. Varshney, J. Prakash, Role of ionic contribution in dielectric behaviour of a nematic liquid crystal with variable cell thickness, J. Mol. Liq. 303 (2020). https://doi.org/10.1016/j.molliq.2020.112520.	509 510 511
[13]	O. V. Kovalchuk, T.M. Kovalchuk, N. Tomašovičová, M. Timko, K. Zakutanska, D. Miakota, P. Kopčanský, O.F. Shevchuk, Y. Garbovskiy, Dielectric and electrical properties of nematic liquid crystals 6CB doped with iron oxide nanoparticles. The combined effect of nanodopant concentration and cell thickness, J. Mol. Liq. 366 (2022). https://doi.org/10.1016/j.molliq.2022.120305.	512 513 514 515 516
[14]	K. Neyts, F. Beunis, Ion transport in liquid crystals, in: J.W. Goodby, C. Tschierske, P. Raynes, H. Gleeson, T. Kato, P.J. Collings (Eds.), Handb. Liq. Cryst., 2nd ed., Wiley-VCH Verlag GmbH & Co. KGaA., Weinheim, Germany, 2014: pp. 1–26.	517 518 519 520

M. Mrukiewicz, P. Perkowski, K. Garbat, R. Dąbrowski, Molecular relaxations in 521 [15] dual-frequency nematic liquid crystals, Liq. Cryst. 41 (2014) 1537-1544. 522 523 https://doi.org/10.1080/02678292.2014.932451. 524 [16] M. Bourg, M. Urbanski, Spectroscopic insight into molecular fluctuations and phase stability of nematic composites containing gold nanoparticles or carbon 525 nanotubes, Phys. Chem. Chem. Phys. 19 (2017) 23302-23308. 526 https://doi.org/10.1039/c7cp02943h. 527 [17] R. de la Fuente, D. Dunmur, Dielectric properties of liquid crystals, in: J.W. 528 Goodby, C. Tschierske, P. Raynes, H. Gleeson, T. Kato, P.J. Collings (Eds.), 529 Handb. Liq. Cryst., 2 nd ed, Wiley-VCH Verlag GmbH & Co. KGaA.: Weinheim, 530 Germany, Weinheim, Germany, 2014: 1–10. 531 pp. https://doi.org/10.1051/anphys/197803030197. 532 [18] U. Maschke, I. Moundoungou, G.J. Fossi-Tabieguia, Method for extracting the 533 liquid crystals contained in an element that comprises a first support and a 534 second support - associated device, FR3017808 (A1), EP3111276 (A1), 535 536 WO2015128582 (A1): 2017-01-04, 2015. https://patentscope.wipo.int/search/fr/detail.jsf?docId=WO2015128582 537 (accessed March 18, 2021). 538 A. Barrera, C. Binet, F. Dubois, P.A. Hébert, P. Supiot, C. Foissac, U. Maschke, 539 [19] Dielectric spectroscopy analysis of liquid crystals recovered from end-of-life 540 liquid crystal displays, Molecules. 26 (2021). 541 https://doi.org/10.3390/molecules26102873. 542 [20] W.H.H. Woodward, Broadband Dielectric Spectroscopy - A Practical Guide, in: 543 ACS Symposium Series (Ed.), Broadband Dielectr. Spectrosc. A Mod. Anal. 544 Tech., 1 st, American Chemical Society, Washington, DC, United States, 2021: 545 pp. 3-59. https://doi.org/10.1021/bk-2021-1375.ch001. 546 547 [21] A.K. Jonscher, The "universal" dielectric response, Nature. 267 (1977) 673-679. https://doi.org/https://doi.org/10.1038/267673a0. 548 D.W. Davidson, R.H. Cole, Dielectric relaxation in glycerine, J. Chem. Phys. 18 549 [22] (1950) 1417. https://doi.org/10.1063/1.1747496. 550 [23] K.S. Cole, R.H. Cole, Dispersion and absorption in dielectrics: II. Direct current 551 characteristics, J. Chem. Phys. 10 (1942)98-105. 552 https://doi.org/10.1063/1.1723677. 553 S. Havriliak, S. Negami, A complex plane representation of dielectric and 554 [24] mechanical relaxation processes in some polymers, Polymer (Guildf). 8 (1967) 555 161-210. https://doi.org/10.1016/0032-3861(67)90021-3. 556 [25] P. Pawel, H.M. Godinho, Liquid crystals. New perspectives, 1 st ed., John Wiley 557 & Sons, Inc., Hoboken, NJ, USA, 2021. 558 [26] F. Kremer, A. Schönhals, Broadband dielectric spectroscopy, 1st ed, Springer-559 Verlag Berlin Heidelberg: New York, USA, 2003. 560 http://books.google.com/books?id=sdG4ywL3qMsC. 561 M.T. Viciosa, A.M. Nunes, A. Fernandes, P.L. Almeida, M.H. Godinho, M.D. 562 [27] Dionísio, Dielectric studies of the nematic mixture E7 on a hdroxypropylcellulose 563 substrate, Liq. Cryst. 29 (2002)429-441. 564 https://doi.org/10.1080/02678290110113478. 565 G. Yadav, R. Katiyar, G. Pathak, R. Manohar, Effect of ion trapping behavior of 566 [28] TiO2 nanoparticles on different parameters of weakly polar nematic liquid crystal, 567 J. Theor. Appl. Phys. 12 (2018) 191-198. https://doi.org/10.1007/s40094-018-568 0296-x. 569 [29] P. Perkowski, D. Łada, K. Ogrodnik, J. Rutkowska, W. Piecek, Z. Raszewski, 570 Technical aspects of dielectric spectroscopy measurements of liquid crystals, 571 Opto-Electronics Rev. 16 (2008) 271-276. https://doi.org/10.2478/s11772-008-572 0008-1. 573 R. Dhar, An impedance model to improve the higher frequency limit of electrical [30] 574 measurements on the capacitor cell made from electrodes of finite resistances, 575 Indian J. Pure Appl. Phys. 42 (2004) 56-61. 576 [31] P. Perkowski, Dielectric spectroscopy of liquid crystals. Electrodes resistivity and 577 connecting wires inductance influence on dielectric measurements, Opto-578 Electronics Rev. 20 (2012) 79-86. https://doi.org/10.2478/s11772-012-0004-3. 579 [32] P. Debye, Polar Molecules, 1st ed, Dover Publications, Inc., New York, USA, 580 1929. https://doi.org/10.1002/ANGE.19290424112. 581 S. Krishna Prasad, M. Vijay Kumar, T. Shilpa, C.V. Yelamaggad, Enhancement 582 [33] 583 of electrical conductivity, dielectric anisotropy and director relaxation frequency in composites of gold nanoparticle and a weakly polar nematic liquid crystal, 584 RSC Adv. 4 (2014) 4453-4462. https://doi.org/10.1039/c3ra45761c. 585 [34] R.J. Mandle, E. Bevis, J.W. Goodby, Phase Structures of Nematic Liquid 586 587 Crystals, in: J.W. Goodby, C. Tschierske, P. Raynes, H. Gleeson, T. Kato, P.J. Collings (Eds.), Handb. Liq. Cryst., 2 nd ed, Wiley-VCH Verlag GmbH & Co. 588 KGaA.. Weinheim, Germany, 2014: 1-27. 589 pp. https://doi.org/10.1002/9783527671403.hlc037. 590 591