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Magnetization transfer from protons to quadrupolar nuclei in solid-state NMR using PRESTO or dipolar-mediated refocused INEPT methods

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Abstract. In solid-state NMR spectroscopy, the through-space transfer of magnetization from protons to quadrupolar nuclei is employed to probe proximities between those isotopes. Furthermore, such transfer, in conjunction with Dynamic Nuclear Polarization (DNP), can enhance the NMR sensitivity of quadrupolar nuclei, as it allows the transfer of DNP-enhanced ¹H polarization to surrounding nuclei. We compare here the performances of two approaches to achieve such transfer: PRESTO (Phase-shifted Recoupling Effects a Smooth Transfer of Order), which is currently the method of choice to achieve the magnetization transfer from protons to quadrupolar nuclei and which has been shown to supersede Cross-Polarization under Magic-Angle Spinning (MAS) for quadrupolar nuclei and *D*-RINEPT (Dipolar-mediated Refocused Insensitive Nuclei Enhanced by Polarization Transfer) using symmetry-based SR4₁² recoupling, which has already been employed to transfer the magnetization in the reverse way from half-integer quadrupolar spin to protons.

We also test the PRESTO sequence with R16₇⁶ recoupling using 270₀90₁₈₀ composite π -pulses as inversion elements. This recoupling scheme, which has previously been proposed to reintroduce ¹H Chemical Shift Anisotropy (CSA) at high MAS frequencies with high robustness to rf-field inhomogeneity, has not so far been employed to reintroduce dipolar couplings with protons. These various techniques to transfer the magnetization from protons to quadrupolar nuclei are analyzed using (i) an average Hamiltonian theory, (ii) numerical simulations of spin dynamics, and (iii) experimental ¹H \rightarrow ²⁷Al and ¹H \rightarrow ¹⁷O transfers in as-synthesized AlPO₄-14 and ¹⁷O-labelled fumed silica, respectively. The experiments and simulations are done at two magnetic fields (9.4 and 18.8 T) and several spinning speeds (15, 18-24 and 60 kHz). This analysis indicates that owing to its γ -encoded character, PRESTO yields the highest transfer efficiency at low magnetic fields and MAS frequencies, whereas owing to its higher robustness to rf-field inhomogeneity and chemical shifts, *D*-RINEPT is more sensitive at high fields and MAS frequencies, notably for protons exhibiting large offset or CSA, such as those involved in hydrogen bonds.

Key words. PRESTO; R-INEPT; HETCOR; quadrupolar nuclei; proton; dipolar coupling; composite π -pulse.

I. Introduction

Quadrupolar nuclei, with nuclear spin quantum number $S > \frac{1}{2}$, represent 75 % of stable NMR-active nuclei [1]. Numerous solids, such as organic compounds, biomolecules, hybrid or porous materials, nanoparticles, hydrates or heterogeneous catalysts, contain both quadrupolar nuclei and protons. For these materials, two-dimensional (2D) *D*-HETCOR (Dipolar-mediated HETeronuclear CORrelation) NMR experiments between quadrupolar nuclei and protons allow the unambiguous identification of

proximities between sites occupied by these isotopes. Hence, these experiments facilitate the assignment of NMR spectra and provide precious information on the atomic-level structure of these materials. For instance, *D*-HETCOR experiments between ^1H and ^{27}Al isotopes have been employed to investigate the aluminum incorporation in aluminosilicate mesoporous material [2], the aluminum local environment in phyllosilicate [3], the dissolution mechanism of aluminosilicate glasses in water [4], the structure of aluminophosphates [5,6], aluminum-based metal-organic frameworks [7], alumina surfaces [8,9], olefin metathesis catalysts supported on chlorinated alumina support [10], methylaluminumoxane-modified silica [11], the nature of Brønsted acid sites at the surface of amorphous silica alumina [12], and the location of Al atoms in zeolites [13,14]. Similarly, ^1H - ^{11}B *D*-HETCOR experiments have been employed to probe the changes in the local environment of boron atoms in borosilicate zeolites in the course of hydration and dehydration [15]. ^1H - ^{43}Ca *D*-HETCOR experiments have also been applied to observe the proximities between Ca atoms and hydroxyl groups in hydroxyapatite materials [16,17]. ^1H - ^{17}O *D*-HETCOR experiments have been employed to examine the structure of silica surfaces and silica-supported catalysts [18,19], the hydrogen bonds in crystalline and amorphous forms of pharmaceutical compounds [20], and crystalline peptides [21]. Furthermore, it has been demonstrated that the sensitivity gain provided by Dynamic Nuclear Polarization (DNP) enables the acquisition of ^1H - ^{17}O *D*-HETCOR 2D spectra for isotopically unmodified solids, despite the low natural abundance of ^{17}O isotope [19,22,23]. Recently, ^1H - ^{35}Cl *D*-HETCOR experiments have been introduced to characterize the molecular-level structure of active pharmaceutical ingredients [24,25]. Besides half-integer spin quadrupolar nuclei, *D*-HETCOR experiments have been used for the indirect detection of ^{14}N isotope which has a spin $S = 1$, via protons [26,27]. Such ^1H - ^{14}N experiments have been used to study the self-assembly of guanosine derivatives [28–30], the intermolecular hydrogen bonds and the nitrogen protonation in pharmaceuticals [30–35], the structure of layered aluminophosphate materials containing amine structure directing agents [5], and the host-guest interactions in metal-organic frameworks functionalized by amine groups [7]. *D*-HETCOR experiments can be achieved using either direct or indirect detection [36].

In direct detection, the magnetization of the excited nucleus is transferred to the detected one. Such transfer between spin-1/2 and quadrupolar isotopes under Magic-Angle Spinning (MAS) has first been performed using Cross-Polarization (CP) [37]. However, CP experiments that involve quadrupolar nuclei present numerous limitations when they are performed under MAS (CPMAS) [38,39]. First, the transfer efficiency is reduced since it is difficult to spin-lock the magnetization of quadrupolar nuclei for all crystallites simultaneously in a rotating powder [40]. Second, for half-integer spin quadrupolar nuclei, the most efficient CPMAS transfers are usually achieved for selective spin-lock of the central transition (CT) with a low radio-frequency (rf) field [39]. As a result, such transfers are then highly sensitive to resonance offset and Chemical Shift Anisotropy (CSA). Third, the optimization is difficult because the efficiency of the spin-lock for a half-integer spin quadrupolar isotope drops at the Rotary Resonance Recoupling (R^3) conditions; i.e. when the nutation frequency of the CT is a multiple of the MAS frequency, ν_R [41]. Fourth, CPMAS transfers are also sensitive to the strength of the quadrupole coupling constant, C_Q , and hence they may not be efficient for two sites exhibiting distinct C_Q values [42].

Alternative *D*-HETCOR methods with direct detection have been introduced in order to circumvent the shortcomings of CPMAS transfers involving quadrupolar nuclei. These approaches include the *D*-RINEPT (Dipolar-mediated Refocused Insensitive Nuclei Enhanced by Polarization Transfer) [43–46] and PRESTO (Phase-shifted Recoupling Effects a Smooth Transfer of Order) [23,47–49] polarization transfers. The first introduced *D*-RINEPT experiment is TEDOR (Transferred-Echo DOuble Resonance) [43,50] using the REDOR (Rotational-Echo DOuble Resonance) scheme [51] as hetero-nuclear dipolar recoupling. However, REDOR does not eliminate the homo-nuclear dipolar couplings and hence is not suitable for *D*-HETCOR experiments with protons. More recently, *D*-RINEPT experiments, in which the

hetero-nuclear dipolar couplings are reintroduced using the R^3 scheme, have been reported [44,45]. In particular, the R^3 scheme using an rf-field $\nu_1 = 2\nu_R$ has been employed to acquire D -RINEPT 2D spectra between protons and half-integer spin quadrupolar isotopes, such as ^{27}Al or ^{17}O [21,44]. Very recently, D -RINEPT experiments, in which heteronuclear couplings with protons are reintroduced using the symmetry-based $\text{SR}4_1^2$ recoupling scheme [52], have also been proposed to correlate protons with spin-1/2 nuclei with low gyromagnetic ratio, such as ^{83}Y , ^{103}Rh or ^{183}W [53], or half-integer spin quadrupolar nuclei, such as ^{35}Cl [46]. Nevertheless, to the best of our knowledge, such sequence has not yet been applied to transfer the magnetization from protons to half-integer spin quadrupolar nuclei.

In the PRESTO sequence, the hetero-nuclear dipolar couplings are reintroduced using symmetry-based single-quantum (1Q) hetero-nuclear γ -encoded dipolar recoupling schemes, such as $\text{R}18_1^7$ or $\text{R}18_2^5$ [23,47,48]. These symmetry-based sequences suppress the homo-nuclear dipolar interactions in the first-order average Hamiltonian. The sensitivity gain afforded by PRESTO has notably been used to transfer the DNP-enhanced proton polarization to ^{17}O , without any sample labelling [23]. So far, the RN_n^γ schemes which have been used in PRESTO experiments employed single π -pulses as inversion element. Recently, schemes based on symmetries, such as $\text{R}20_9^8$, $\text{R}18_8^7$, $\text{R}14_6^5$, $\text{R}16_7^6$ and $\text{R}12_5^4$ and using 270_090_{180} composite π -pulses as inversion element, have been introduced to measure the ^1H CSA at MAS frequencies of 60 and 70 kHz [54]. Here, the standard notation for composite pulses is used: ξ_φ indicates a rectangular, resonant rf-pulse with flip angle ξ and phase φ , and the angles are written in degrees. These schemes reintroduce the same components of the spin interactions as the RN_n^γ schemes employed in PRESTO, but they benefit from higher robustness to rf-field inhomogeneity. However, to the best of our knowledge, these symmetry-based sequences have not yet been employed to reintroduce the heteronuclear dipolar interactions.

D -HETCOR experiments can also be used to increase the sensitivity for the NMR detection of half-integer spin quadrupolar isotopes. Such sensitivity gain has notably been reported when PRESTO scheme is used to transfer the DNP-enhanced proton polarization to quadrupolar nuclei, such as ^{17}O and ^{27}Al [23,48].

Indirect detection is an alternative to direct detection for D -HETCOR experiments; i.e. the excited isotope is also the detected one and the coherences are transferred back and forth between the isotopes. These indirectly detected D -HETCOR experiments particularly include the D -HMQC (Dipolar-mediated Hetero-nuclear Multiple-Quantum Correlation) schemes [44,55,56]. In these experiments, the hetero-nuclear dipolar couplings are reintroduced using various schemes, such as REDOR, R^3 , SFAM (Simultaneous Frequency and Amplitude Modulation) and symmetry-based sequences [36,57–59]. In the case of D -HMQC experiments correlating protons and quadrupolar isotopes, the symmetry-based $\text{SR}4_1^2$ recoupling [52] is often employed. Indeed, this scheme: (i) exhibits high efficiency and robustness, (ii) is compatible with high MAS frequency, and (iii) can easily be optimized. We have also recently introduced another indirectly detected D -HETCOR experiment, called D -HUQC (Dipolar-mediated Hetero-nuclear Universal-Quantum Correlation), which employs γ -encoded symmetry-based recoupling schemes on the detected channel, and exhibits lower t_1 -noise in the case of nuclei subject to large CSA [60].

The relative sensitivities of direct and indirect detections depend notably on the gyromagnetic ratios, the longitudinal relaxation times and the spectral widths of the correlated isotopes [46,61]. Furthermore, contrary to the directly detected D -HETCOR experiments, those using indirect detection cannot be used to acquire directly 1D spectra of quadrupolar isotope by transferring the DNP-enhanced polarization of protons to the nearby quadrupolar nuclei.

We focus here on the directly detected D -HETCOR experiments with proton excitation and detection of quadrupolar isotope. We compare the efficiency and the robustness of two techniques: D -RINEPT using $\text{SR}4_1^2$ recoupling and PRESTO using simple 180_0 pulses or 270_090_{180} composite ones.

The two techniques are first described using an average Hamiltonian theory. Their performances are then compared using numerical simulations of spin dynamics and ^1H - ^{27}Al and ^1H - ^{17}O experiments on $\text{AlPO}_4\text{-14}$ and fumed silica, respectively.

II. Pulse sequences and theory

II.1. PRESTO

In the PRESTO sequence, the hetero-nuclear dipolar couplings between the protons and the quadrupolar nuclei are reintroduced under MAS by the application on the ^1H channel of symmetry-based γ -encoded recoupling schemes, such as R18_1^7 , R18_2^5 , R16_3^2 , R18_4^1 and R16_7^6 . Those schemes recouple the $|m| = 2$ space components and the one-quantum (1Q) terms of hetero-nuclear dipolar coupling (^1H -S) and ^1H CSA (CSA_H) [47], while they suppress the ^1H isotropic chemical shifts, the heteronuclear J -couplings with protons and the ^1H - ^1H dipolar couplings in the first-order average Hamiltonian. The rf-field requirements of R18_1^7 , R18_2^5 , R16_3^2 , R18_4^1 and R16_7^6 recouplings with simple π -pulses are $v_1/v_R = 9, 4.5, 2.66, 2.25$ and 1.14 , respectively, whereas that of R16_7^6 scheme using 270_090_{180} composite π -pulses, denoted $\text{R16}_7^6\text{-C}$ hereafter is 2.28 . Schemes with high rf-field requirements, such as R18_1^7 , may not be compatible with fast MAS. However, it must be noted that other γ -encoded recoupling schemes, with much lower rf-field requirements can be used [23].

During these recoupling schemes, the contribution of the dipolar coupling between $I = ^1\text{H}$ and S quadrupolar nuclei to the first-order average Hamiltonian is equal to [47]:

$$\bar{H}_{D,IS}^{(1)} = \omega_{D,IS} S_z [I^+ \exp(i2\varphi) + I^- \exp(-i2\varphi)] \quad (1)$$

where $I^\pm = I_x \pm iI_y$ are the shift operators. In Eq.1, the magnitude and phase of the recoupled I -S dipolar coupling are given by

$$\omega_{D,IS} = -\kappa \frac{\sqrt{3}}{2} b_{IS} \sin^2(\beta_{PR}^{D,IS}) \quad (2)$$

$$\varphi = \gamma_{PR}^{D,IS} + \alpha_{RL}^0 - \omega_R t^0 \quad (3)$$

In Eq.2, (i) b_{IS} is the dipolar coupling constant in $\text{rad}\cdot\text{s}^{-1}$, (ii) the dipolar scaling factor $\kappa = 0.182, 0.175, 0.161$ and 0.152 for $\text{R18}_1^7, \text{R18}_2^5, \text{R16}_3^2$ and R18_4^1 schemes, respectively, with simple π -pulse as basic element and $\kappa = 0.15$ for R16_3^2 scheme using 270_090_{180} composite π -pulse, which is denoted $\text{R16}_3^2\text{-C}$ hereafter, and (iii) the Euler angles $\{0, \beta_{PR}^{D,IS}, \gamma_{PR}^{D,IS}\}$ relate the inter-nuclear I -S vector to the MAS rotor frame. In Eq.3, $\omega_R = 2\pi\nu_R$ and t^0 refers to the starting time of the symmetry-based scheme. The norm of $\bar{H}_{D,IS}^{(1)}$ does not depend on the $\gamma_{PR}^{D,IS}$ angle and hence, these recoupling schemes are called γ -encoded [45,62]. The recoupled Hamiltonian described in Eq.1 does not commute among different spin-pairs and the PRESTO experiment is hence affected by dipolar truncation, which may limit the observation of long I -S inter-nuclear distances. However, it must be noted that such dipolar truncation has been used to selectively correlate the signals of covalently bonded ^{13}C and ^1H nuclei [63,64].

In PRESTO, these heteronuclear dipolar recoupling schemes also reintroduce CSA_H , with the same scaling factor, and its first-order average Hamiltonian is equal to [47]:

$$\bar{H}_{CSA,I}^{(1)} = \omega_{CSA,I}^* I^+ + \omega_{CSA,I} I^- \quad (4)$$

where $\omega_{CSA,I}$ is the frequency of the recoupled CSA_H and $\omega_{CSA,I}^*$ is its complex conjugate. This frequency is given by

$$\omega_{CSA,I} = -\frac{\kappa}{\sqrt{2}} [A_{22}^{CSA,I}]^R \exp\{-2i(\alpha_{RL}^0 - \omega_R t^0)\} \quad (5)$$

where $[A_{22}^{CSA,I}]^R$ is given by Eq.5 in ref.[36]. Eqs.1 and 4 show that the recoupled CSA_H and S - ^1H dipolar coupling terms do not commute. Therefore, the spin dynamics during PRESTO simultaneously depends on both CSA_H and S - ^1H hetero-nuclear dipolar coupling.

In the present article, we mainly employed the PRESTO-III variant, which is depicted in Fig.1a [47]. A π -pulse is applied at the centers of the two defocusing and refocusing periods, denoted τ and τ' ,

respectively, and simultaneously the phase of the ^1H channel irradiation is shifted by 180° . Such procedure limits the interference of CSA_H , because this interaction is only fully refocused when its tensor is axially symmetric and collinear with the S - H vector. During the τ delay, the longitudinal ^1H magnetization is converted into ^1H 1Q coherences, which are antiphase with respect to the S spin. The $\pi/2$ pulse on the S channel transforms the antiphase ^1H 1Q coherences into S 1Q coherence antiphase with respect to ^1H . A π -pulse is also applied at the center of the refocusing period, τ' , in order to refocus the evolution under the isotropic shifts of the S nuclei, whereas the phase of the ^1H channel irradiation is shifted by 180° . Furthermore, the $\tau/2$ and $\tau'/2$ delays must be integer multiples of the rotor period so that the evolution under CSA_H and the second-order quadrupole interaction of the S nucleus is better refocused. Herein, we employed $\tau = \tau'$.

For a S spin coupled to a single proton with vanishing CSA_H , the PRESTO signal with $\tau = \tau'$ is proportional to

$$S(\tau) \propto \langle \sin^2(\omega_{I,IS}\tau) \rangle = \frac{1}{2} - \frac{1}{3^{1/4}} \sqrt{\frac{\pi}{8\kappa b_{IS}\tau}} \left\{ F_c \left(3^{1/4} \sqrt{\frac{2\kappa b_{IS}\tau}{\pi}} \right) \cos(\sqrt{3}\kappa b_{IS}\tau) + F_s \left(3^{1/4} \sqrt{\frac{2\kappa b_{IS}\tau}{\pi}} \right) \sin(\sqrt{3}\kappa b_{IS}\tau) \right\} \quad (6)$$

where the angular bracket $\langle \dots \rangle$ denotes the powder average. Eq.6 was derived using a closed analytical form for γ -encoded $|m| = 1$ recoupling sequences, and $F_c(x)$ and $F_s(x)$ are the Fresnel cosine and sine integrals, respectively [62]. This equation can be used for distance determination, in place of the spin dynamics simulations that have been employed in Ref [23]. In the absence of losses and CSA_H , the shorter τ value producing the maximal signal intensity is given by:

$$\text{PRESTO (CSA}_\text{H} = 0) \rightarrow \tau^{\text{opt}} = 2.18/(\kappa b_{IS}) \quad (7)$$

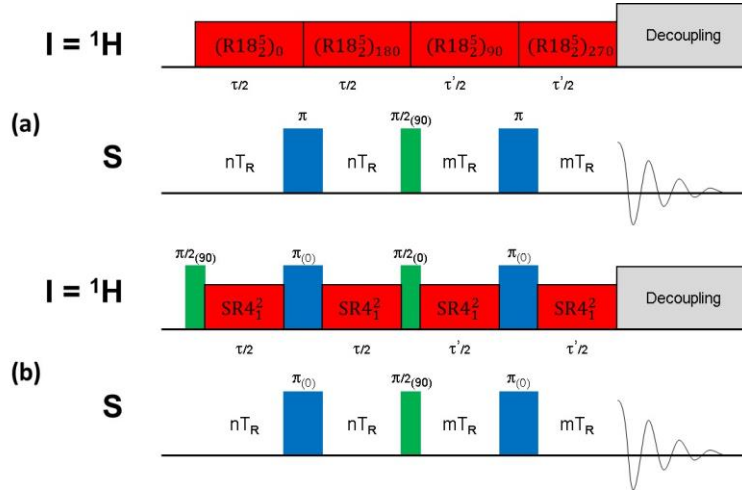


Fig.1. Pulse sequences for $I = ^1\text{H} \rightarrow S$ transfers: (a) PRESTO-III- R18_2^5 and (b) D -RINEPT- SR4_1^2 . The S isotope is quadrupolar with half-integer spin value and the S pulses are CT-selective. For the acquisition of D -HETCOR 2D spectra, the period t_1 is inserted (a) before the first R18_2^5 block and bracketed by two $\pi/2$ -pulses, (b) between the first $\pi/2$ -pulse and the first SR4_1^2 block. The quadrature detection along the indirect dimension was achieved using the States-TPPI procedure [65] by incrementing the phase of the first pulse prior to t_1 period. The phase cycling and pulse programs of the various sequences are given at the end of the SI.

II.2. D -RINEPT

In D -RINEPT, described in Fig.1b, the S - ^1H dipolar couplings are reintroduced using the SR4_1^2 recoupling [52]. This scheme is a 3-step multiple-quantum super-cycled version of $\text{R4}_1^2\text{R4}_1^{-2}$, each block lasting one rotor period: $T_R = 1/\nu_R$, i.e. $\text{SR4}_1^2 = [\text{R4}_1^2\text{R4}_1^{-2}]_0[\text{R4}_1^2\text{R4}_1^{-2}]_{120}[\text{R4}_1^2\text{R4}_1^{-2}]_{240}$, with $\text{R4}_1^2 = \pi_{90}\pi_{-90}\pi_{-90}\pi_{90}$ and $\text{R4}_1^{-2} = \pi_{-90}\pi_{90}\pi_{90}\pi_{-90}$, where π_{90} and π_{-90} denote resonant, rectangular π -pulses

on ^1H channel with phase γ and $-\gamma$, respectively. This sequence, during which the protons are irradiated with $\nu_1 = 2\nu_R$, reintroduces the CSA_H and the $|m| = 1$ space component of the I-S dipolar coupling, whereas it suppresses the ^1H isotropic chemical shifts, the J_{IS} -couplings, and the ^1H - ^1H dipolar couplings to the first-order. The SR4_1^2 scheme achieves zero-quantum hetero-nuclear dipolar recoupling and the contribution of the I-S dipolar coupling to the first-order average Hamiltonian is equal to [52,58]:

$$\bar{H}_{D,IS}^{(1)} = 2\omega_{D,IS}I_zS_z \quad (8)$$

$$\omega_{D,IS} = \frac{1}{4}b_{IS}\sin^2(\beta_{PR}^{D,IS})\cos(2\varphi). \quad (9)$$

The norm of $\bar{H}_{D,IS}^{(1)}$ depends on the φ phase, given by Eq.3, and hence on the $\gamma_{PR}^{D,IS}$ angle. Therefore, the SR4_1^2 scheme is non- γ -encoded. The recoupled Hamiltonian of Eq.8 commutes among different spin-pairs, hence allowing the observation of long I-S inter-nuclear distances. Furthermore, the CSA_H term recoupled by SR4_1^2 is proportional to I_z and thus commutes with the I-S dipolar interaction of Eq.8. Hence, the CSA_H does not interfere with the evolution under I-S dipolar interaction during D -RINEPT.

In the D -RINEPT sequence, the first $\pi/2$ -pulse creates a transverse ^1H magnetization in-phase with respect to the S spin. During the defocusing delay, τ , this magnetization evolves into transverse ^1H magnetization antiphase with respect to the S spin. The simultaneous π -pulses on S and ^1H channels at the center of the τ delay refocus the evolution under the CSA_H , while allowing that under the I-S dipolar interaction. Simultaneous $\pi/2$ -pulses on the S and ^1H channels convert the antiphase ^1H magnetization into antiphase S one. During the refocusing delay τ' , this antiphase S magnetization is transformed into transverse in-phase S magnetization, which is detected during the acquisition period. The simultaneous π -pulses on S and ^1H channels at the center of the τ' delay refocus the evolution under the S isotropic chemical shifts, while allowing that under I-S dipolar couplings. For a proton coupled to a single S spin, the NMR signal of D -RINEPT- SR4_1^2 experiment with $\tau = \tau'$ is proportional to

$$S(\tau) \propto \frac{1}{2} \left\{ 1 - \frac{\pi\sqrt{2}}{4} J_{1/4} \left(\frac{b_{IS}}{4} \tau \right) J_{-1/4} \left(\frac{b_{IS}}{4} \tau \right) \right\} \quad (10)$$

where $J_{\pm 1/4}(x)$ denotes the Bessel functions of the first kind and $\pm 1/4$ -order. In the absence of losses and CSA_H , the shorter τ value producing the maximal signal intensity is given by:

$$D\text{-RINEPT} (\text{CSA}_\text{H} = 0) \rightarrow \tau^{\text{opt}} = 9.44/b_{IS} \quad (11)$$

III. Experimental section

III.1. Simulation parameters

All numerical simulations of spin dynamics were performed with the SIMPSON software (version 4.1.1) [66]. The powder average was calculated using 2304 $\{\alpha_{\text{MR}}, \beta_{\text{MR}}, \gamma_{\text{MR}}\}$ Euler angles. The 256 $\{\alpha_{\text{MR}}, \beta_{\text{MR}}\}$ angles, which relate the molecular and rotor frames, were selected according to the REPULSION algorithm [67], while the 9 γ_{MR} angles were equally stepped from 0 to 360° . The dsyev method, with the corresponding Linear Algebra PACKage (LAPACK) [68], was used to accelerate the simulations [69]. During the PRESTO and D -RINEPT sequences, only CT-selective pulses are applied to the quadrupolar isotope. Therefore, its satellite transitions weakly contribute to the detected signal. This statement is supported by Fig. S1, showing that the simulated signal of D -RINEPT sequences for ^{13}C - $^1\text{H}_4$ and ^{27}Al - $^1\text{H}_4$ spin-systems are very similar, whereas the CPU time was 30-fold shorter for the former spin system than for the latter one. The ratio of the CPU times required for simulations on ^{13}C - $^1\text{H}_4$ and ^{27}Al - $^1\text{H}_4$ spin systems is consistent with the fact that the duration of the SIMPSON simulations is limited by the matrix-matrix multiplications and matrix diagonalizations. As the number of arithmetic operations for these processes scales with the cube of the dimension of the density matrix, a simulation for ^{13}C - $^1\text{H}_4$ spin system should be ca. 27-fold faster than for ^{27}Al - $^1\text{H}_4$ one [70]. Therefore, all simulations (except those of Fig. S1b) were performed for two spin-1/2 isotopes, $I = ^1\text{H}$ and $S = ^{13}\text{C}$, in order to accelerate the simulations.

The simulations were carried out for one isolated ^{13}C - ^1H spin-pair, except those in Figs. 6-8, S1a, S6 and S9, which were carried for a ^{13}C - $^1\text{H}_4$ spin-system in order to compare the robustness of the sequences to ^1H - ^1H dipolar couplings, and those of Fig. S1b, which were carried for a ^{27}Al - $^1\text{H}_4$ spin-system. In these five-spins systems, the four protons were located on the vertices of a tetrahedron and the ^{13}C or ^{27}Al nucleus was located on a symmetry axis of this tetrahedron. All ^1H - ^1H dipolar coupling constants were identical and the ^{13}C or ^{27}Al nucleus was dipolar coupled with its closest proton with $|b_{\text{IS}}|/(2\pi) = 1$ or 6 kHz. The CSA_{H} value of the ^1H coupled to the ^{13}C or ^{27}Al nucleus is indicated in the figure captions, its asymmetry parameter is null, and the orientation of its principal axis systems with respect to the vector between its position and the ^{13}C or ^{27}Al nucleus is described by the Euler angles $(0, 30^\circ, 0)$.

The static magnetic field was fixed at $B_0 = 18.8$ T ($\nu_{0,^1\text{H}} = 800$ and $\nu_{0,^{13}\text{C}} = 201$ MHz) for all simulations, except for those of Fig. 7, which were carried out at 9.4 T. The use of high magnetic fields is beneficial for half-integer spin quadrupolar nuclei, notably because the line-widths of the central transition are inversely proportional to B_0 and hence, the spectral resolution is proportional to B_0^2 . In Figs. 2-6, 8 and S7, the MAS frequency was 22 or 24 kHz (indicated as $\nu_{\text{R}} \approx 23$ kHz), for R18_2^5 or SR4_1^2 schemes, respectively, to correctly sample the first maximum of the build-up curves with $|b_{\text{IS}}|/(2\pi) = 6$ kHz. Such MAS frequencies correspond to those typically used for rotor with 3.2 mm outer diameter. In Figs. S2 to S9 (except Fig. S7), the MAS frequency was $\nu_{\text{R}} = 60$ kHz for all recoupling sequences. This MAS frequency is accessible using rotor with an outer diameter of 1.3 mm and is generally required to achieve high resolution for the ^1H spectra without the use of ^1H - ^1H dipolar decoupling sequence [71]. In Fig. 7, the MAS frequency was 15 kHz to correspond to most DNP experiments.

We simulated the powder averaged signal of $\text{PRESTO-III-R18}_2^5$ and $D\text{-RINEPT-SR4}_1^2$ sequences, except in Fig. 7 where simulations were carried out for $\text{PRESTO-III-R16}_7^6$ using either single π -pulses or 270_090_{180} composite ones. The simulations carried out for PRESTO-II-R18_2^5 [not shown] confirm that this method is less robust to CSA_{H} than the PRESTO-III variant. The $\text{PRESTO-III-R18}_2^5$ and $D\text{-RINEPT-SR4}_1^2$ sequences are denoted PRESTO and RINEPT hereafter. The pulses, which do not belong to the recoupling blocks were simulated as ideal Dirac pulses, except in Fig. S1b, where ^{27}Al CT-selective long pulses had to be used. The pulses of the recoupling schemes were applied on resonance, except in Fig. 3 and S3, for which the ^1H resonance offset was varied. For spin systems containing ^{13}C , the transfer efficiencies of PRESTO and RINEPT were calculated as the ratios between the simulated signals of these experiments and that with a ^{13}C direct excitation with an ideal $\pi/2$ -pulse. For Fig. S1b, the transfer efficiency of $^1\text{H} \rightarrow ^{27}\text{Al}$ RINEPT experiment was calculated as the ratio of its simulated signal and that with a ^{27}Al direct excitation with a CT-selective $\pi/2$ -pulse. Note that in SIMPSON simulations, the signal is not proportional to the gyromagnetic ratio, and hence the calculation of the transfer efficiency does not require to be normalized by the ratio of the gyromagnetic ratios. The build-up curves shown in Figs. 2 and S2 were simulated using the shortest possible increments for $\tau/2 = \tau'/2$ delays, i.e. $2T_{\text{R}}/9$ for R18_2^5 , corresponding to the length of a $\pi_{50}\pi_{-50}$ block, and $T_{\text{R}}/2$ for SR4_1^2 , corresponding to the length of a $\pi_{90}\pi_{-90}$ block. The shortest possible increment for R16_7^6 and $\text{R16}_7^6\text{-C}$ schemes is $7T_{\text{R}}/8$.

III.2. Solid-state NMR experiments

For the experiments, all the rotors were fully packed. RINEPT experiments were acquired with SR4_1^2 recoupling scheme, whereas the PRESTO experiments were recorded with R18_2^5 and $\text{R16}_7^6\text{-C}$ at $\nu_{\text{R}} = 15$ kHz, R18_2^5 at $\nu_{\text{R}} = 20$ kHz and R16_3^2 , R18_4^1 and $\text{R16}_7^6\text{-C}$ at $\nu_{\text{R}} = 60$ kHz. $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT experiments were performed on an as-synthesized $\text{AlPO}_4\text{-14}$ sample with isopropylamine inserted into the pores [72], and the recycle delay was $\tau_{\text{RD}} = 1$ s. The ^{27}Al isotropic chemical shifts were referenced to 1 M solution $\text{Al}(\text{NO}_3)_3$, whereas the ^1H isotropic chemical shifts were referenced to tetramethylsilane using the resonance of adamantane (1.74 ppm) as a secondary reference. $^1\text{H} \rightarrow ^{17}\text{O}$ spectra were recorded on a fumed silica and the ^{17}O isotropic chemical shifts were referenced to water at 0 ppm.

III.2.a. Experiments at 9.4 T

Even if DNP-enhanced NMR experiments have been reported up to 21.1 T [73] and MAS frequency up to 40 kHz [74], DNP-enhanced PRESTO experiments have so far mainly been reported using 9.4 T and 3.2 mm rotors [23,48]. Therefore, we first recorded the $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT 2D experiments of $\text{AlPO}_4\text{-14}$ using a 9.4 T wide-bore magnet equipped with an Avance-II Bruker console. The experiments were recorded using a 3.2 mm HXY MAS probe used in the double resonance mode. The rf-fields of the pulses other than those used during the recoupling scheme were equal to 86 and 11 kHz on the ^1H and ^{27}Al channels, respectively.

III.2.b. Experiments at 18.8 T

As mentioned above, high B_0 field is beneficial for the detection of quadrupolar nuclei. Therefore, PRESTO and RINEPT experiments transferring the magnetization of protons to quadrupolar nuclei were also performed on an 18.8 T narrow-bore Bruker magnet with HX double-resonance MAS probes. $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT experiments were performed on $\text{AlPO}_4\text{-14}$ using rotors with outer diameter of 3.2 and 1.3 mm, whereas $^1\text{H} \rightarrow ^{17}\text{O}$ spectra were recorded with 3.2 mm diameter. With 3.2 mm rotors, experiments were recorded with an Avance III console, whereas we used an Avance IV one with 1.3 mm rotors.

For $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT experiments, the rf-field of the pulses other than those used during the recoupling scheme was $\nu_{1,1\text{H}} = 77$ and 208 kHz, and $\nu_{1,27\text{Al}} = 10$ and 14 kHz to achieve CT-selective excitation, at $\nu_R = 20$ and 60 kHz, respectively.

$^1\text{H} \rightarrow ^{17}\text{O}$ spectra were acquired at $\nu_R = 18$ kHz on a fumed silica with specific surface area of 350 m^2/g , for which the surface was ^{17}O enriched using a previously reported procedure [18]. The 1D ^{17}O direct excitation MAS spectra were acquired using a single-pulse and QCPMG (quadrupolar Carr-Purcell-Meiboom-Gill) sequences [75]. Except during the recoupling parts on the proton channel, the rf-fields were $\nu_{1,1\text{H}} = 100$ and $\nu_{1,17\text{O}} = 8$ kHz.

IV. Numerical simulations of the $^1\text{H} \rightarrow ^{13}\text{C}$ transfer

IV.1. Build-up curves at 18.8 T with single π -pulses

Fig. 2 shows the build-up curves of $^1\text{H} \rightarrow ^{13}\text{C}$ PRESTO and RINEPT transfers for an isolated $^{13}\text{C}\text{-}^1\text{H}$ spin pair at $\nu_R \approx 23$ kHz. For $\text{CSA}_\text{H} = 0$, the PRESTO sequence exhibits stronger oscillations than RINEPT and a higher maximal transfer efficiency (0.73 for the former instead of 0.52 for the latter). Such differences are consistent with the γ -encoding recoupling used for PRESTO and the non- γ -encoding employed in RINEPT. For both sequences, the optimal recoupling times are in agreement with those predicted from Average Hamiltonian theory (Eqs. 7 and 11). In addition, we can observe that the build-up curves of RINEPT are not affected by CSA_H , unlike those of PRESTO. This robustness of RINEPT to CSA_H stems from its commutation with $^{13}\text{C}\text{-}^1\text{H}$ dipolar terms recoupled by $\text{SR}4_1^2$, whereas those terms do not commute with $\text{R}18_2^5$ (see section II). Finally, the comparison of Figs. 2a and c proves a larger influence of the CSA_H on the PRESTO build-up curve in the case of small $^{13}\text{C}\text{-}^1\text{H}$ dipolar couplings. It is noted that similar build-up curves are obtained for $\nu_R = 60$ kHz (Fig. S2).

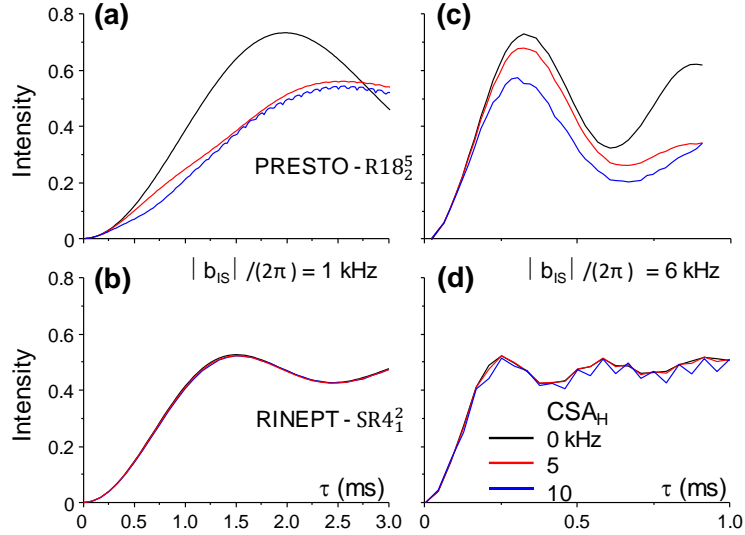


Fig.2. Simulated build-up curves of $^1\text{H} \rightarrow ^{13}\text{C}$ transfer at 18.8 T and $\nu_R \approx 23$ kHz of (a,c) PRESTO-R18 $_2^5$ with single π -pulses or (b,d) RINEPT for an isolated ^{13}C - ^1H spin-pair with $|b_{\text{IS}}|/(2\pi) =$ (a,b) 1 or (c,d) 6 kHz, and $\text{CSA}_{\text{H}} = 0, 5$ or 10 kHz.

IV.2. Robustness to offset at 18.8 T with single π -pulses

The robustness of RINEPT is higher than that of PRESTO, especially for weak I-S dipolar couplings (Fig.3). Similar results are obtained at $\nu_R = 60$ kHz (Fig.S3). This high robustness to offset of SR4 $_1^2$ stems from the super-cycling, which better eliminates the unwanted cross-terms involving offset in the higher-order terms of the Average Hamiltonian. In the case of RINEPT, as expected by the Average Hamiltonian theory, the efficiency of the transfer versus the offset weakly depends on CSA_{H} . Furthermore, the robustness to ^1H offset improves for increasing MAS frequency since the rf-fields of the recoupling sequences are proportional to the MAS frequency (compare Figs.3 and S3).

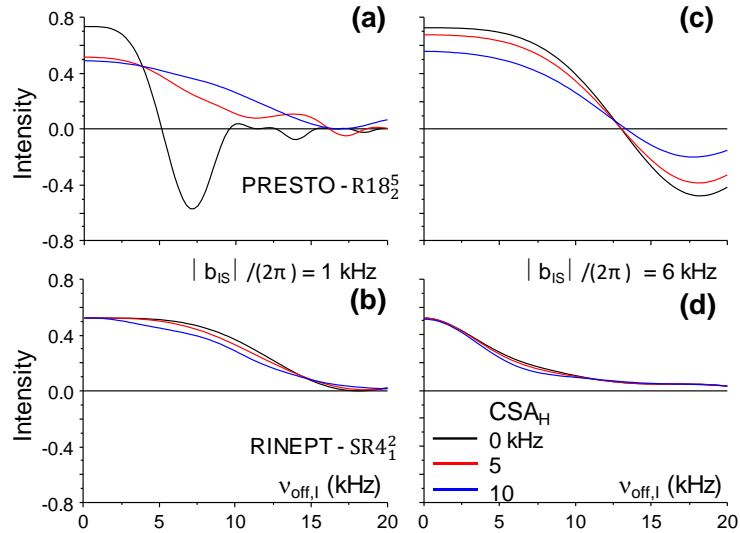


Fig.3. Simulated transfer efficiency versus the ^1H resonance offset, $\nu_{\text{off},l}$, at 18.8 T and $\nu_R \approx 23$ kHz for (a,c) PRESTO-R18 $_2^5$ with single π -pulses or (b,d) RINEPT, with $|b_{\text{IS}}|/(2\pi) =$ (a,b) 1 or (c,d) 6 kHz, and $\text{CSA}_{\text{H}} = 0, 5$ or 10 kHz. The τ value was set to its optimum value determined from Fig.2.

IV.3. Robustness to ^1H CSA at 18.8 T with single π -pulses

^1H CSA values can be as large as 30 ppm for protons involved in hydrogen bonds [76–78]. Such CSA corresponds to 24 kHz at 18.8 T. Figs. 4 and S4 display a comparison of the robustness to CSA_H of PRESTO and RINEPT transfers at $\nu_\text{R} \approx 23$ and 60 kHz, respectively. As already observed in Figs. 2 and 3 as well as S2 and S3, the RINEPT method exhibits higher robustness to CSA_H than PRESTO. This is consistent with the commutation between CSA_H and hetero-nuclear dipolar coupling terms recoupled by SR4_1^2 , whereas those terms do not commute with R18_2^5 . As expected, the robustness to CSA_H is improved at $\nu_\text{R} = 60$ kHz, as the rf-fields of the recoupling schemes are proportional to the MAS frequency.

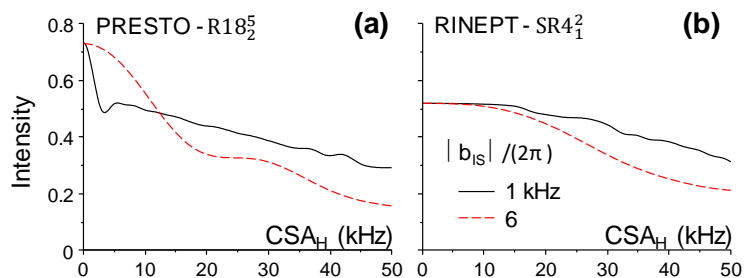


Fig. 4. Simulated transfer efficiency versus CSA_H at 18.8 T and $\nu_\text{R} \approx 23$ kHz for (a) PRESTO-R18₂⁵ with single π -pulses and (b) RINEPT with $|b_{\text{IS}}|/(2\pi) = 1$ or 6 kHz. The τ value was set to its optimum value determined from Fig. 2.

IV.4. Robustness to rf-field inhomogeneity at 18.8 T with single π -pulses

The rf-field produced by a solenoid coil depends on the position inside the rotor [79–83]. It is known that for rotor diameters of 1.3 and 3.2 mm the minimal rf-field at the ends of the rotor is approximately 25 % of its maximum value [79,81]. In Figs. 5 and S5, the simulated transfer efficiencies of PRESTO and RINEPT experiments are plotted against the ratio between the applied and theoretical rf-fields, $R_\text{rf} = \nu_\text{1l}/\nu_\text{1,th}$. For both schemes, the transfer efficiency is below 10 % for $R_\text{rf} \leq 0.25$ [not shown]. Furthermore, these simulations show that the RINEPT scheme is much more robust to rf-inhomogeneity than the PRESTO one with single π -pulses. This result is attributed to the use of the SR4_1^2 recoupling, which is constructed from the amplitude-modulated basic sequences $\text{R4}_1^{\pm 2}$, i.e. the phase shift between consecutive π -pulses is 180°. This amplitude modulation achieves a compensation for rf-field errors [84,85].

These simulations indicate that the robustness to rf-inhomogeneity of SR4_1^2 does not depend on the b_{IS} value, while that of PRESTO increases with this value. Actually, the difference $|\nu_\text{1l} - \nu_\text{1,th}|$ must be smaller than $|b_{\text{IS}}|/(2\pi)$ with R18_2^5 , as already observed for the γ -encoded R^3 recoupling [45]. Therefore, at $\nu_\text{R} = 60$ kHz, the R_rf relative interval yielding high transfer efficiency for PRESTO is smaller than at $\nu_\text{R} \approx 23$ kHz (compare Figs. 5a,b and S5a,b) since $|\nu_\text{1l} - \nu_\text{1,th}| \leq |b_{\text{IS}}|$ and $\nu_\text{1,th}$ is proportional to ν_R .

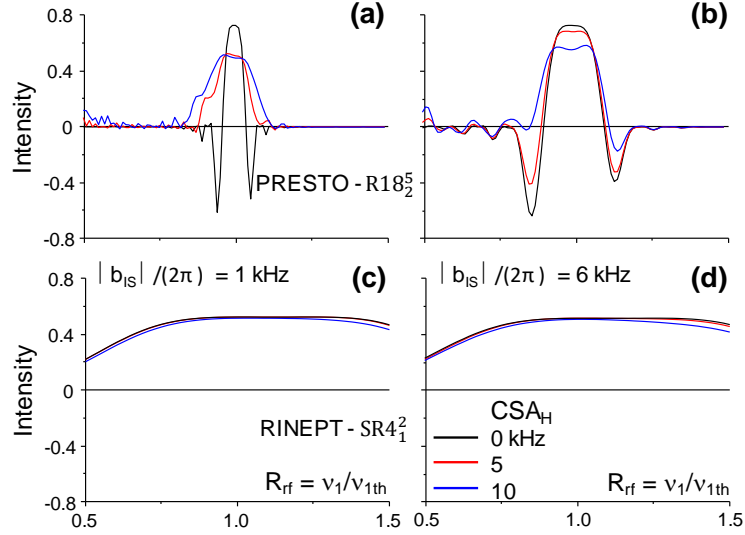


Fig.5. Simulated transfer efficiency at 18.8 T and $\nu_R \approx 23$ kHz versus the rf-inhomogeneity, $R_{rf} = \nu_1/\nu_{1th}$, for (a,b) PRESTO-R18₂⁵ with single π -pulses or (c,d) RINEPT, with $CSA_H = 0, 5, 10$ and $|b_{IS}|/(2\pi) = 1$ (a,c) or 6 (b,d) kHz. The τ value was set to its optimum value determined from Fig.2.

In Figs.5 and S5, no 1H - 1H dipolar coupling interaction was considered. Figs.6 and S6 display the simulated robustness to rf-inhomogeneity of PRESTO and RINEPT sequences at $\nu_R \approx 23$ and 60 kHz, respectively, for the ^{13}C - 1H_4 spin-system described in Section III.1, with $|b_{HH}|/(2\pi) = 0, 1$ or 7 kHz and $|b_{IS}|/(2\pi) = 1$ kHz. These simulations show that the robustness of PRESTO to rf-inhomogeneity does not depend on b_{HH} , whereas that of RINEPT decreases for increasing 1H - 1H dipolar interactions. The effect of these interactions on the RINEPT robustness does not depend on the MAS frequency (compare Figs.6b and S6b). Nevertheless, for the investigated spin systems with $|b_{HH}|/(2\pi)$ up to 7 kHz, RINEPT still exhibits higher robustness with respect to rf-inhomogeneity than PRESTO.

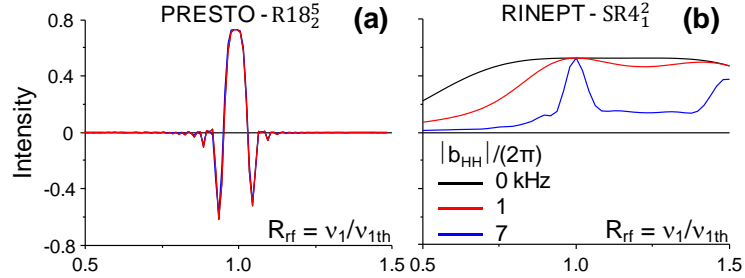


Fig.6. Simulated transfer efficiency at 18.8 T and $\nu_R \approx 23$ kHz versus the rf-inhomogeneity, $R_{rf} = \nu_1/\nu_{1th}$, for (a) PRESTO-R18₂⁵ with single π -pulses or (b) RINEPT sequences applied to ^{13}C - 1H_4 spin system with $|b_{HH}|/(2\pi) = 0, 1$ or 7 kHz, $|b_{IS}|/(2\pi) = 1$ kHz and $CSA_H = 0$. The τ value was set to its optimum value determined from Fig.2.

IV.5. Robustness to MAS frequency fluctuations at 18.8 T with single π -pulses

The transfer efficiencies of PRESTO and RINEPT versus the relative deviation, $R_{\nu_R} = (\nu_R - \nu_{R,th})/\nu_{R,th}$, of the actual MAS frequency from its theoretical value, $\nu_{R,th} \approx 23$ or 60 kHz, are plotted in Figs.S7 or S8, respectively. As expected, the sensitivity to MAS fluctuations is higher for smaller hetero-nuclear dipolar coupling constants, which require longer recoupling times. Moreover, the absolute line-width of the efficiency curve only depends on $|b_{IS}|$ and therefore the relative deviation (R_{ν_R}) is inversely proportional to the spinning speed (compare Figs.S7 and S8). For $|b_{IS}|/(2\pi) = 6$ kHz, PRESTO and RINEPT experiments exhibit similar robustness to MAS fluctuations. However, for $|b_{IS}|/(2\pi) = 1$ kHz, this robustness decreases for increasing CSA_H in the case of RINEPT, whereas it increases in the case of

PRESTO. Hence, in the case of small dipolar couplings between S-spin and protons subject to significant CSA_H , PRESTO is more robust to MAS fluctuations than RINEPT. However, even in that case, the simulations show that the stability of the MAS frequency achieved using the latest generation of MAS speed controllers is sufficient to avoid significant intensity losses for both PRESTO and RINEPT.

IV.6. Robustness to 1H - 1H dipolar interactions

Protons in solids, notably in organic and hybrid ones, are often coupled to several other protons. Therefore, the $^1H \rightarrow ^{13}C$ transfer efficiency was simulated for the ^{13}C - 1H_4 spin system (described in section III.1) versus the 1H - 1H dipolar coupling constant, $|b_{HH}|/(2\pi)$. The results are shown in Figs. 7, 8 and 9, for $\nu_R = 15, 23$ and 60 kHz, respectively. For the three MAS frequencies and for both PRESTO and RINEPT schemes, the effect of 1H - 1H dipolar couplings is larger for smaller $|b_{IS}|/(2\pi)$ values due to longer recoupling times. For the three MAS frequencies, PRESTO-R18₂⁵ and RINEPT exhibit similar robustness to 1H - 1H dipolar couplings, whereas SR4₁² recoupling employs an rf-field 2.25-fold smaller than that of R18₂⁵. As seen in Fig. 7, R16₁⁶ recoupling, which employs an rf-field 40% smaller than that of SR4₁² and 4-fold smaller than that of R18₂⁵, is much more sensitive to 1H - 1H dipolar couplings than these schemes. The replacement of single π -pulses by 270₀90₁₈₀ ones increases both the required rf-field and the robustness to 1H - 1H dipolar couplings. However, R16₁⁶-C recoupling is more sensitive to 1H - 1H dipolar couplings than R18₂⁵ and SR4₁² schemes, even if the rf-field of R16₁⁶-C is 28% higher than that of SR4₁². The comparison of Figs. 7, 8 and 9 also shows that the use of high MAS frequency improves the robustness to 1H - 1H dipolar interactions for both PRESTO and RINEPT experiments.

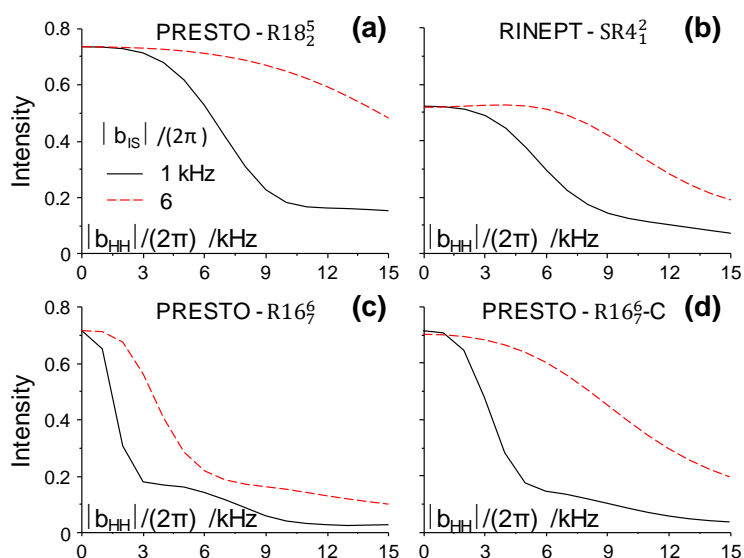


Fig. 7. Simulated on-resonance transfer efficiency at 9.4 T with $\nu_R = 15$ kHz, $CSA_H = 0$ and $|b_{IS}|/(2\pi) = 1$ or 6 kHz, versus $|b_{HH}|/(2\pi)$ in ^{13}C - 1H_4 spin system for (a,c,d) PRESTO with R18₂⁵ (a) and R16₁⁶ (c) with single π -pulses, or R16₁⁶-C with composite π -pulses (d), as well as (b) RINEPT. The τ value were set to 2, 1.46, 4.67 and 2.22 ms for subfigures a, b, c and d, respectively.

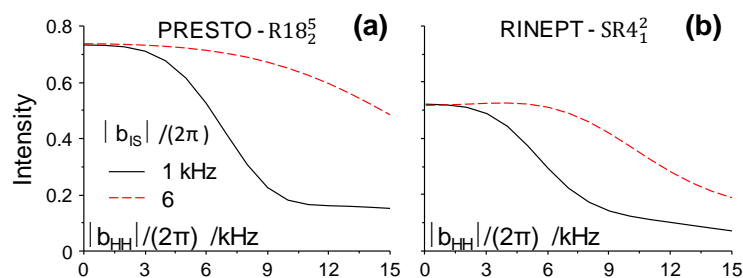


Fig.8. Simulated on-resonance transfer efficiency at 18.8 T and $\nu_R \approx 23$ kHz versus the $|b_{HH}|/(2\pi)$ constant in $^{13}\text{C}-^1\text{H}_4$ spin system for (a) PRESTO-R18 $_2^5$ with single π -pulses and (b) RINEPT schemes with $\text{CSA}_H = 0$ and $|b_{IS}|/(2\pi) = 1$ or 6 kHz. The τ value was set to its optimum value determined from Fig.2.

V. NMR experiments

V.1. Experiments on $\text{AlPO}_4\text{-14}$ at 9.4 T with $\nu_R = 15$ kHz

The crystal structure of $\text{AlPO}_4\text{-14}$ exhibits four crystallographically inequivalent Al sites: two AlO_4 , one AlO_5 and one AlO_6 , with $C_Q = 1.8, 4.1, 5.6$ and 2.6 MHz, respectively [86]. Even at 18.8 T, the two AlO_4 resonances overlap, and thus only their sum signal is given in Tables 1 to 3. After optimization of the τ delay and the rf-field of the recoupling scheme [not shown], we recorded three D -HETCOR 2D spectra with ^{27}Al detection, denoted $^{27}\text{Al}\{-^1\text{H}\}$ hereafter, of $\text{AlPO}_4\text{-14}$ using $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO-R18 $_2^5$ (Fig.9a), PRESTO-R16 $_7^6$ -C and RINEPT sequences. The ^1H dimension of the 2D spectra exhibits three resolved proton signals, NH_3^+ , CH and CH_3 . For the three 2D spectra, all ^1H - ^{27}Al cross-peaks were detected. As seen in Table 1, the signal-to-noise ratio (S/N) of the cross-peaks of PRESTO-R18 $_2^5$ is in average 50% higher than that of RINEPT. The higher sensitivity of PRESTO-R18 $_2^5$ stems from the γ -encoding of R18 $_2^5$ schemes, which results in higher transfer efficiency, whereas the SR4 $_1^2$ scheme is non- γ -encoded (Fig.2). Nevertheless, the S/N ratio of the RINEPT 2D spectrum is in average about 73% higher than that of PRESTO-R16 $_7^6$ -C. The lower sensitivity of PRESTO-R16 $_7^6$ -C must stem from its lower robustness to ^1H - ^1H dipolar couplings (Fig.7).

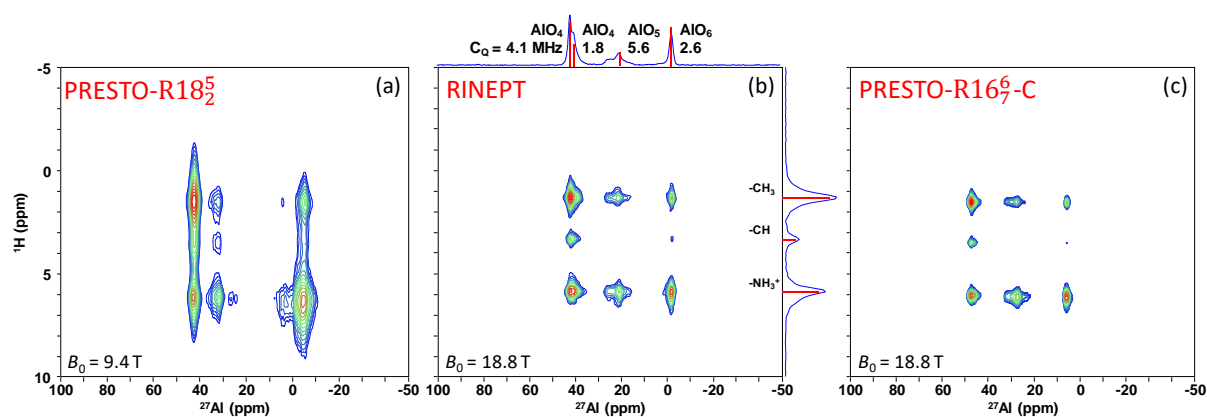


Fig.9. $^{27}\text{Al}\{-^1\text{H}\}$ 2D spectra of $\text{AlPO}_4\text{-14}$ acquired with $\tau = 800$ μs , using the following conditions: (a) $\nu_R = 15$ kHz, 9.4 T with PRESTO-R18 $_2^5$ and single π -pulses, (b) $\nu_R = 20$ kHz, 18.8 T with RINEPT, (c) $\nu_R = 60$ kHz, 18.8 T with PRESTO-C-R16 $_7^6$. The ^1H rf-field during the recoupling, the number of scans, the number of t_1 increments, and the total experimental time, were equal to $(\nu_1$ (kHz), NS, N_1 , T_{exp}) = (68, 128, 100, 3.5 h) for (a), (46, 4, 128, 10 min) for (b) and (140, 128, 128, 4.5 h) for (c). The assignment of ^1H and ^{27}Al signals is shown on the projections: The C_Q values are indicated on top of the ^{27}Al projection of the spectrum (b).

Table 1. S/N ratios at 9.4 T with AVANCE-II console and $\nu_R = 15$ kHz of the cross-peaks in $^{27}\text{Al}\{-^1\text{H}\}$ 2D HETCOR spectra of $\text{AlPO}_4\text{-14}$ acquired with RINEPT, PRESTO-R18 $_2^5$ with single π -pulse, and PRESTO-R16 $_7^6$ -C.

Sequence	PRESTO-R18 $_2^5$			RINEPT			PRESTO-R16 $_7^6$ -C		
$\delta_{\text{iso},^{27}\text{Al}}$ \ / \ $\delta_{\text{iso},^1\text{H}}$	42 Al ^{IV}	22 Al ^V	-2 Al ^{VI}	42 Al ^{IV}	22 Al ^V	-2 Al ^{VI}	42 Al ^{IV}	22 Al ^V	-2 Al ^{VI}
1.3 (CH $_3$)	68	37	65	44	25	34	27	16	17
3.3 (CH)	43	15	20	28	10	10	16	6	5
5.8 (NH $_3^+$)	95	17	32	81	16	25	52	9	15

V.2. Experiments on $\text{AlPO}_4\text{-14}$ at 18.8 T with single π -pulses and $\nu_R = 20$ kHz.

In order to test the influence of the B_0 field on the sensitivity, $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO-R18 $_2^5$ and RINEPT experiments were carried out on $\text{AlPO}_4\text{-14}$ at $B_0 = 18.8$ T with $\nu_R = 20$ kHz. Figs.10b,e show the build-up curves of the four ^{27}Al signals obtained with $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT. The experimental build-up curves significantly differ from the simulated ones of Fig.2. In particular, it is noted that the build-up curve of PRESTO exhibits smaller oscillations than the simulated curves. Furthermore, the experimental optimal recoupling time is $\tau \approx 800$ μs for PRESTO and RINEPT, whereas according to average Hamiltonian theory and numerical simulations for an isolated spin-pair, the optimal recoupling time of PRESTO-R18 $_2^5$ should be 32 % longer than that of RINEPT (Eqs.7 and 11). The discrepancy between simulations and experiments may be attributed to the presence of several protons in the sample instead of isolated spin-pairs (Fig.9), as well as the coherent and incoherent losses during the τ delays.

The $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT signal intensity as function of the rf-field of the recoupling scheme is shown in Figs.10a,d. The maximal signal intensity is achieved for nutation frequencies close to the theoretical ones: $\nu_1 \approx 4.5\nu_R$ for PRESTO and $2\nu_R$ for RINEPT. Furthermore, in agreement with simulations (Fig.5), the RINEPT-SR4 $_1^2$ recoupling is more robust to rf-inhomogeneity than PRESTO with single π -pulses. The intervals of rf-field, for which the signal intensity is larger than half of its maximal value, are equal to 15 and 40 kHz for PRESTO and RINEPT, respectively, which correspond to 17 and 100 %, in relative value.

Furthermore, as seen in Figs.10c,f the PRESTO sequence, which employs γ -encoded R18 $_2^5$ recoupling, is slightly less sensitive to the MAS frequency than RINEPT, which uses the non- γ -encoded SR4 $_1^2$ scheme. However, modern speed controllers can achieve stability much better than ± 20 Hz. In such interval of MAS frequency, both PRESTO and RINEPT sequences are insensitive to MAS frequency fluctuations.

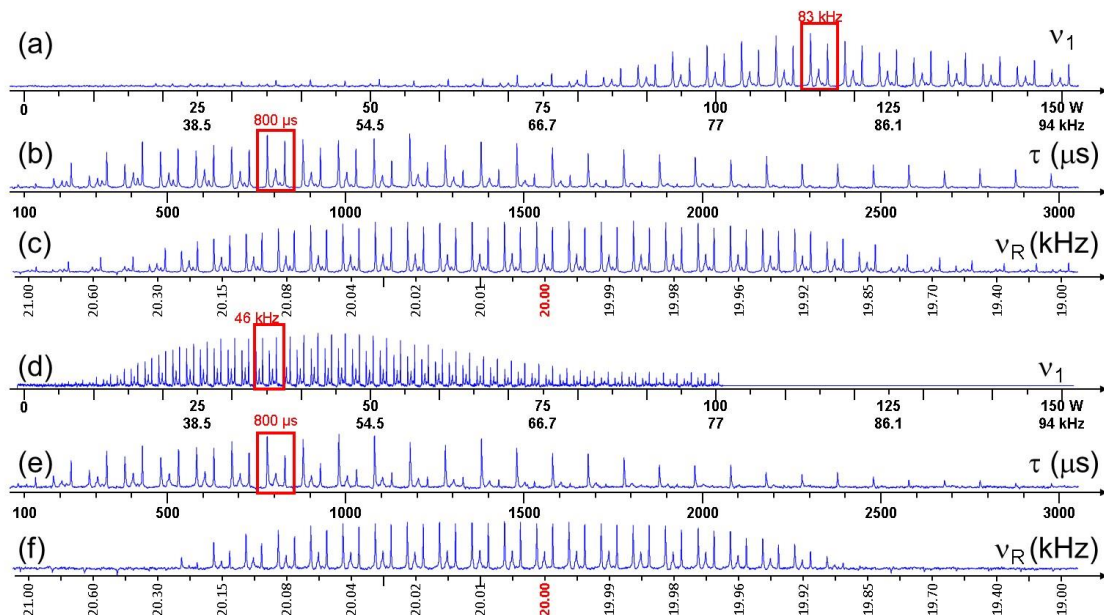


Fig. 10. $^1\text{H} \rightarrow ^{27}\text{Al}$ (a-c) PRESTO-R18₂⁵ with single π -pulses and (d-f) RINEPT spectra of AlPO₄-14 at $B_0 = 18.8$ T versus ν_1 (a,d), τ (b,e) and ν_R (c,f) using $\nu_R = 20$ kHz in (a,b,d,e), $\tau = 800$ μs in (a,c,d,f), $\nu_1 = 83$ (b,c) or 46 kHz (e,f). Each spectrum was recorded using $NS = 32$.

Fig. 9b shows the $^{27}\text{Al}\{-^1\text{H}\}$ RINEPT 2D spectrum of AlPO₄-14 at $B_0 = 18.8$ T with $\nu_R = 20$ kHz. As expected, it exhibits a higher resolution along the ^{27}Al dimension than in Fig. 9a, since for half-integer quadrupolar nuclei, the resolution is proportional to B_0^2 . Furthermore, the ^{27}Al isotropic shifts are distinct between 9.4 and 18.8 T since the quadrupolar induced shifts are inverse-proportional to B_0 . For both $^{27}\text{Al}\{-^1\text{H}\}$ RINEPT and PRESTO-R18₂⁵ 2D spectra, all $^{27}\text{Al}\{-^1\text{H}\}$ cross-peaks were detected, even if the (AlO₅, CH) cross-peak exhibits a small intensity and is not visible in Fig. 9b. As seen in Table 2, the cross-peaks of the PRESTO-R18₂⁵ spectrum exhibit a S/N ratio in average 36% higher than for the RINEPT spectrum acquired within an identical experimental time. The sensitivity gain for PRESTO-R18₂⁵ method with respect to RINEPT scheme decreases with increasing B_0 field. Such decrease stems notably from the larger ^1H offset and CSA at high field.

Table 2. S/N ratios at 18.8 T with AVANCE-III console and $\nu_R = 20$ kHz of the cross-peaks in $^{27}\text{Al}\{-^1\text{H}\}$ 2D HETCOR spectra of AlPO₄-14 acquired with RINEPT and PRESTO-R18₂⁵ with single π -pulses.

Sequence	PRESTO-R18 ₂ ⁵			RINEPT		
$\delta_{\text{iso},^{27}\text{Al}}$	42	22	-2	42	22	-2
$\delta_{\text{iso},^1\text{H}}$	Al ^{IV}	Al ^{IV}	Al ^{VI}	Al ^{IV}	Al ^V	Al ^{VI}
1.3 (CH ₃)	83	14	35	61	11	20
3.3 (CH)	23	4	7	18	3	4
5.8 (NH ₃ ⁺)	66	24	72	65	21	53

V.3. Experiments on AlPO₄-14 at 18.8 T with $\nu_R = 60$ kHz.

It can be desirable to transfer the magnetization of protons to half-integer spin quadrupolar nuclei at fast MAS, which improves the ^1H resolution by averaging out the $^1\text{H}\{-^1\text{H}\}$ dipolar couplings [71]. Fast MAS also enhances by a factor of 3 to 4 the spectral resolution of half-integer spin quadrupolar nuclei subject to large quadrupole interactions by separating the spinning sidebands from the center-band [87]. As the rf-field requirement of R18₂⁵, $\nu_1 \approx 4.5\nu_R$ with single π -pulses, is incompatible with the rf-specifications of most 1.3 mm MAS probes, PRESTO experiments were carried out using R16₃², R18₄¹ and R16₇⁶-C recoupling schemes, which only require $\nu_1/\nu_R \approx 2.66$, 2.25 and 2.28, respectively.

We acquired the build-up curves of RINEPT and PRESTO experiments at $\nu_R = 60$ kHz [not shown] and found optimal recoupling times of $\tau \approx 800$ μ s. This is the same value as that obtained for the experiments performed at $\nu_R = 20$ kHz. Fig. 11 shows the $^1\text{H} \rightarrow ^{27}\text{Al}$ PRESTO and RINEPT signals of $\text{AlPO}_4\text{-14}$ versus ν_1 . The full widths at half maximum of NMR signal intensity as function of the rf-field are similar (ca. 20 kHz) for R16_3^2 and R18_4^1 schemes at $\nu_R = 60$ kHz and for R18_5^2 at $\nu_R = 20$ kHz, whereas they are two-fold broader at $\nu_R = 60$ than at 20 kHz for RINEPT (compare Figs. 10a,d and 11). Hence, the tolerated relative deviation of the rf-field is about $\pm 6\%$ for R18_4^1 and R16_3^2 and $\pm 20\%$ for RINEPT at $\nu_R = 60$ kHz, instead of $\pm 12\%$ for R18_5^2 and $\pm 30\%$ for RINEPT at $\nu_R = 20$ kHz. These experimental results are consistent with the simulations shown in Figs. 5 and S5 and they indicate that the PRESTO sequence with single π -pulses is more sensitive to the rf-field homogeneity at high MAS frequency than RINEPT since the rf-field inhomogeneity corresponds to a relative variation of the rf-field amplitude in the sample space. On the contrary, PRESTO with composite π -pulses is much more robust to rf-inhomogeneity than with single π -pulses, as observed when comparing Fig. 11d with Figs. 11b and c. Furthermore, for both PRESTO with single π -pulses and RINEPT experiments, the maximal signal intensity is obtained for ν_1 values higher than the theoretical ones: 161, 147 and 177 instead of 120, 135, and 160 kHz, for SR4_1^2 , R18_4^1 and R16_3^2 recoupling, respectively. In fact, owing to the rf-field inhomogeneity in the sample space, higher signal can be obtained when the field at the center of the rotor exceeds the theoretical value.

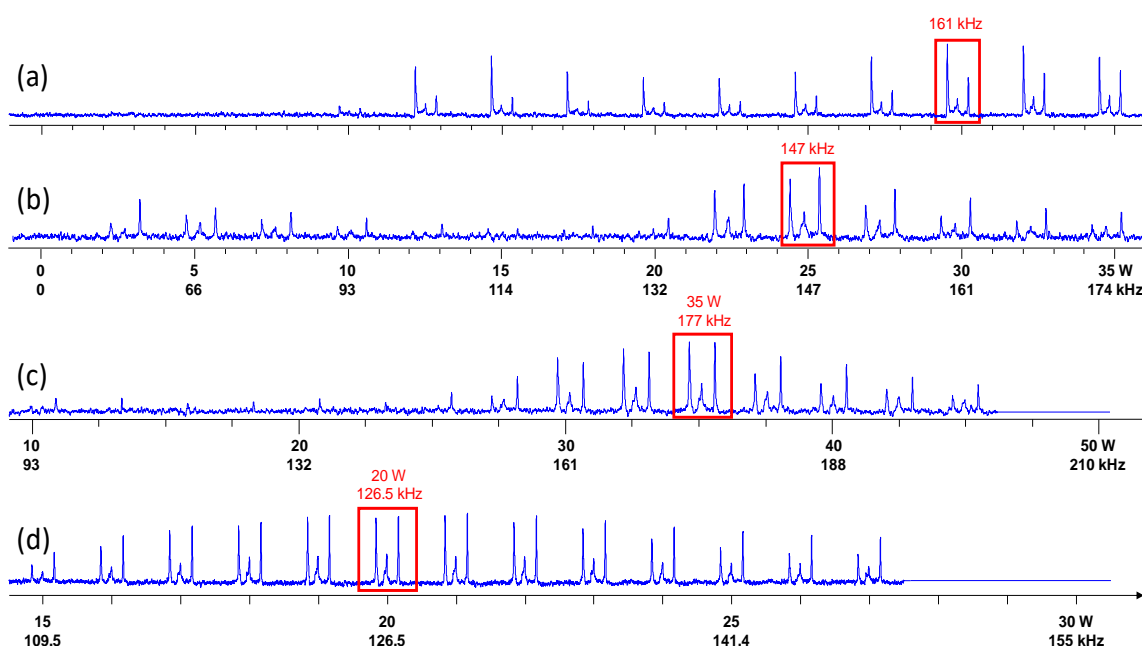


Fig. 11. $^1\text{H} \rightarrow ^{27}\text{Al}$ spectra of $\text{AlPO}_4\text{-14}$ versus ν_1 at 18.8 T with RINEPT (a), and PRESTO with either single (b,c) or composite π -pulses (d), with $\nu_R = 60$ kHz, $\tau = 800$ μ s and $\text{NS} = 32$. With PRESTO, the recoupling scheme is (b) R18_4^1 , (c) R16_3^2 or (d) $\text{R16}_3^2\text{-C}$. It must be noted that the horizontal scales of (c) and (d) differ from that of (a) and (b).

Table 3 compares the S/N ratios of the cross-peaks in $^{27}\text{Al}\{-^1\text{H}\}$ 2D spectra recorded at $B_0 = 18.8$ T with $\nu_R = 60$ kHz using PRESTO- R16_3^2 with single π -pulses, PRESTO- $\text{R16}_3^2\text{-C}$ and RINEPT. The PRESTO- R16_3^2 variant was chosen since it exhibits a slightly higher transfer efficiency than PRESTO- R18_4^1 . Fig. 9c displays the $^{27}\text{Al}\{-^1\text{H}\}$ PRESTO- $\text{R16}_3^2\text{-C}$ 2D spectrum of $\text{AlPO}_4\text{-14}$. As expected, the linewidth of the cross-peaks along the ^1H dimension is about three times lower at $\nu_R = 60$ than at 20 kHz, the B_0 field being constant (compare Figs. 9b and c). Even if the acquisition times of the 2D spectra are 42 times longer at $\nu_R = 60$ than at 20 kHz, their S/N are much smaller. Such sensitivity decrease stems from the smaller sample volume in 1.3 mm rotor with respect to 3.2 mm one. Table 3 indicates that RINEPT is in average

1.8 and 1.3-fold more sensitive than PRESTO-R16₃² and PRESTO-R16₇⁶-C, respectively, at $B_0 = 18.8$ T with $\nu_R = 60$ kHz. Under such condition, RINEPT exhibits the highest sensitivity since it is much more robust to rf-inhomogeneity, especially at high MAS frequencies (see Figs. **S5**, **S6** and **11**). Furthermore, we have shown that our 1.3 mm HX probe suffers from poor rf-field homogeneity [81]. Interestingly, PRESTO-R16₇⁶-C method is more sensitive than PRESTO-R16₃² at 18.8 T with $\nu_R = 60$ kHz, whereas it is less sensitive than PRESTO-R18₂⁵ at 9.4 T with $\nu_R = 15$ kHz. Such difference stems from (i) the higher robustness to rf-inhomogeneity of PRESTO-R16₇⁶-C owing to the use of 270₀90₁₈₀ π -pulses, (ii) the decrease of non-averaged ¹H-¹H interactions at ultra-fast MAS, and (iii) the increased robustness to these interactions of PRESTO-R16₇⁶-C for higher MAS frequency, and hence higher rf-field of the pulses.

Table 3. S/N ratios at 18.8 T with AVANCE-IV console and $\nu_R = 60$ kHz of the cross-peaks in ²⁷Al-{¹H} 2D HETCOR spectra of AlPO₄-14 with RINEPT, PRESTO-R16₃² with single π -pulses, and PRESTO-R16₇⁶-C.

Sequence	PRESTO-R16 ₃ ²			RINEPT			PRESTO-R16 ₇ ⁶ -C		
$\delta_{\text{iso},27\text{Al}}$	42	22	-2	42	22	-2	42	22	-2
$\delta_{\text{iso},1\text{H}}$	Al ^{IV}	Al ^V	Al ^{VI}	Al ^{IV}	Al ^V	Al ^{VI}	Al ^{IV}	Al ^V	Al ^{VI}
1.3 (CH ₃)	56	10	21	121	14	39	111	18	29
3.3 (CH)	12	2	3	28	3	7	21	3	4
5.8 (NH ₃ ⁺)	37	30	75	85	36	111	63	34	70

V.4. Experiments on ¹⁷O labeled fumed silica at 18.8 T with $\nu_R = 18$ kHz

1D MAS spectra of ¹⁷O-labelled fumed silica were recorded at $B_0 = 18.8$ T with $\nu_R = 18$ kHz and they are shown in Fig. **12**. The direct excitation spectrum is shown in Fig. **12a**. The de-shielded resonance is assigned to ¹⁷O nuclei in siloxane bridges, whereas the shielded one is assigned to ¹⁷O nuclei of silanol groups. As seen in Fig. **12b**, the use of QCPMG detection improves the sensitivity for the ¹⁷O siloxane signal. However, the silanol signal is then absent owing to its short T_2' constant time since the dipolar coupling between ¹H and ¹⁷O nuclei leads to a rapid decay of the maximum of the echo signals during the QCPMG scheme. Conversely, the siloxane ¹⁷O nuclei, which are not bonded to protons, exhibit longer T_2' value, hence allowing the acquisition of 16 echoes.

No signal was detected with ¹H \rightarrow ¹⁷O CPMAS experiments for this sample. Such lack of signal illustrates the difficulty to optimize the CPMAS experiment when the S/N ratio is low, as it is the case for this sample. Conversely, as seen in Figs. **12c** and **d**, signals were detected for ¹H \rightarrow ¹⁷O RINEPT and PRESTO-R18₂⁵ with single π -pulses. Both spectra exhibit signals for ¹⁷O siloxane and silanol nuclei; the last signal being more intense than that of siloxane, whereas it is the reverse for the direct excitation of ¹⁷O spectra (compare Figs. **12a** with **c,d**). Such variation in signal intensity stems from the more efficient ¹H \rightarrow ¹⁷O magnetization transfer for silanol than for siloxane since the ¹H-¹⁷O distance is shorter for the former group than for the latter one. In addition, the S/N ratio of the ¹H \rightarrow ¹⁷O RINEPT spectrum of fumed silica is 50 % higher than that of PRESTO-R18₂⁵, whereas for AlPO₄-14, more efficient ¹H \rightarrow ²⁷Al transfers were achieved at $\nu_R = 20$ kHz using PRESTO, instead of RINEPT. RINEPT exhibits comparable robustness to rf-field inhomogeneity for both samples since (i) the strength of the ¹H-¹H dipolar interactions determines the robustness to rf-field inhomogeneity of RINEPT (see Fig. **6b**) and (ii) the protons of fumed silica at $\nu_R = 18$ kHz and AlPO₄-14 at $\nu_R = 20$ kHz [not shown] exhibit NMR signals of comparable widths and hence, are subject to comparable ¹H-¹H dipolar interactions. Furthermore, the shorter optimal τ delay for PRESTO in the case of fumed silica with respect to AlPO₄-14 (222 and 800 μ s) indicates a larger ¹H-S dipolar coupling for the former sample, which should increase the robustness to rf-inhomogeneity of PRESTO. Therefore, the inversion of the relative efficiencies of RINEPT and PRESTO in fumed silica with respect to AlPO₄-14 does not stem from a change in the robustness to rf inhomogeneity but from the higher robustness to CSA_H of RINEPT with respect to PRESTO (Fig. **4**) since the ¹H spectrum of the sample is dominated by a resonance at 2.9 ppm, typical of hydrogen-bonded

silanols [18], which can be subject to CSA_H as large as 30 ppm, i.e. 24 kHz at 18.8 T [77]. Conversely the 1H CSAs in $AlPO_4-14$ are expected to be smaller than 18 and 7 ppm for NH_3^+ and aliphatic protons, respectively [78].

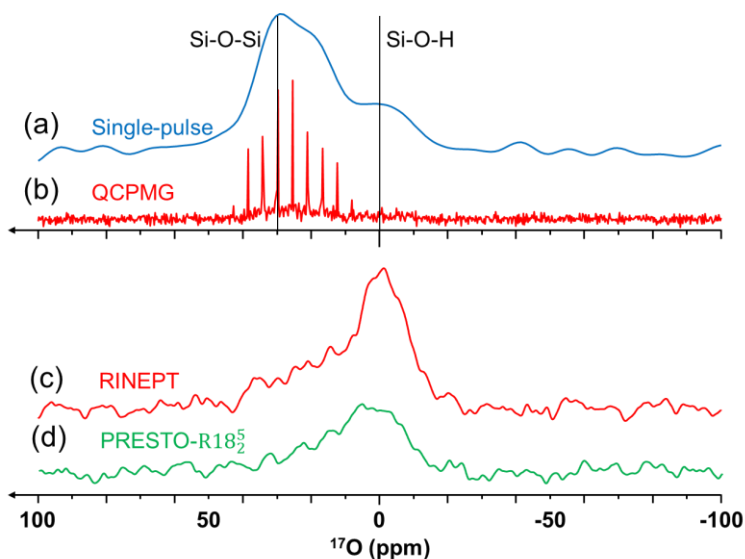


Fig. 12. ^{17}O MAS 1D spectra of labelled fumed silica at 18.8 T with $\nu_R = 18$ kHz. Direct excitation recorded with (a) single-pulse and (b) QCPMG recycling with 16 echoes. $^1H \rightarrow ^{17}O$ (c) RINEPT and (d) PRESTO-R18 $_2^5$ with single π -pulses, $\tau = 222$ μs and $\nu_1 = 36$ and 81 kHz for RINEPT and PRESTO-R18 $_2^5$, respectively. The assignment of ^{17}O resonances is displayed on the top. Spectra (a) and (b) were recorded with $NS = 4,000$ and $\tau_{RD} = 0.5$ s, i.e. $T_{exp} = 34$ min, whereas the spectra (c) and (d) were recorded with $NS = 40,000$ and $\tau_{RD} = 1$ s, i.e. $T_{exp} = 11$ h 7 min.

VI. Conclusion

We have compared the performances of PRESTO and RINEPT experiments to transfer the magnetization of protons to half-integer spin quadrupolar isotopes. These two methods use different types of schemes: PRESTO employs γ -encoded RN_n^γ recouplings, whereas RINEPT uses the $SR4_1^2$ scheme, which is not. The reported simulations and experiments indicate that these techniques complement each other. Owing to its γ -encoding, PRESTO yields the highest transfer efficiency at low MAS frequency and low B_0 field, especially in the case of limited 1H offset and CSA. Conversely, RINEPT benefits from the highest transfer efficiency at high MAS frequency and high B_0 field, owing to its higher robustness to rf-field inhomogeneity, offset and CSA. In particular, RINEPT is beneficial for the transfer of polarization from protons subject to large CSA, such as those involved in hydrogen bonds, and notably will be useful for DNP experiments at high B_0 field, in which the DNP-enhanced magnetization of protons involved in hydrogen bonds has to be transferred to quadrupolar nuclei [23]. At high MAS frequency, the transfer efficiency of PRESTO can be enhanced by the use of composite 270_090_{180} π -pulses within the R16 $_2^5$ -based recoupling scheme. These composite pulses improve the robustness of PRESTO to rf-field inhomogeneity. Nevertheless, the transfer efficiency of this PRESTO-R16 $_2^5$ -C technique remains smaller than that of RINEPT at high MAS frequency and high B_0 field. Furthermore, the robustness to 1H - 1H dipolar coupling of this PRESTO-R16 $_2^5$ -C sequence is not sufficient at low MAS frequencies. PRESTO sequences employing composite pulses suitable for low MAS frequencies, and hence DNP, are currently under investigation and will be presented elsewhere.

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